

Functional Grading of Rubber-Elastic Materials: From Chemistry to Mechanics

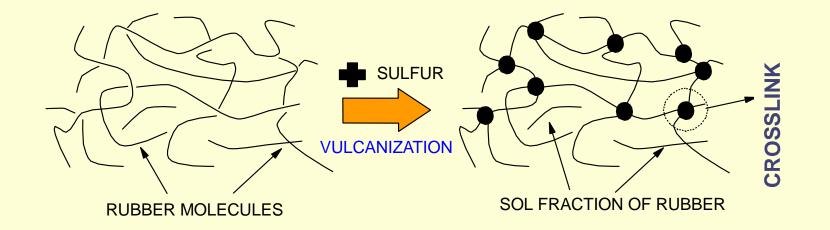
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Outline of the Presentation

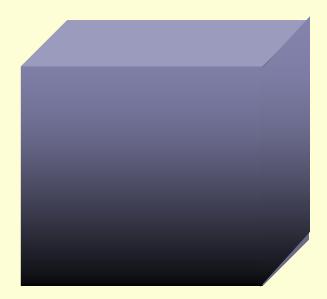
- Overview of rubbers and basic concepts
- Experimental results: potential of grading
- A theoretical study on the effects of
 - material non-homogeneity
 - temperature gradients
 - functional grading
- How to manufacture FGREMs?
 - modeling the structure development
- Conclusions

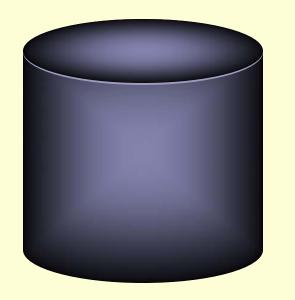
Overview of Rubber-like Materials

- Rubbers widely used in industry and daily life
- Applications: tires, bearings, bushings, sealants, etc.
- Relatively soft, light-weight, and flexible. Undergo large strains, energy dissipation, sub-ambient T_a.
- 3-D network structure of long, flexible polymeric chains:

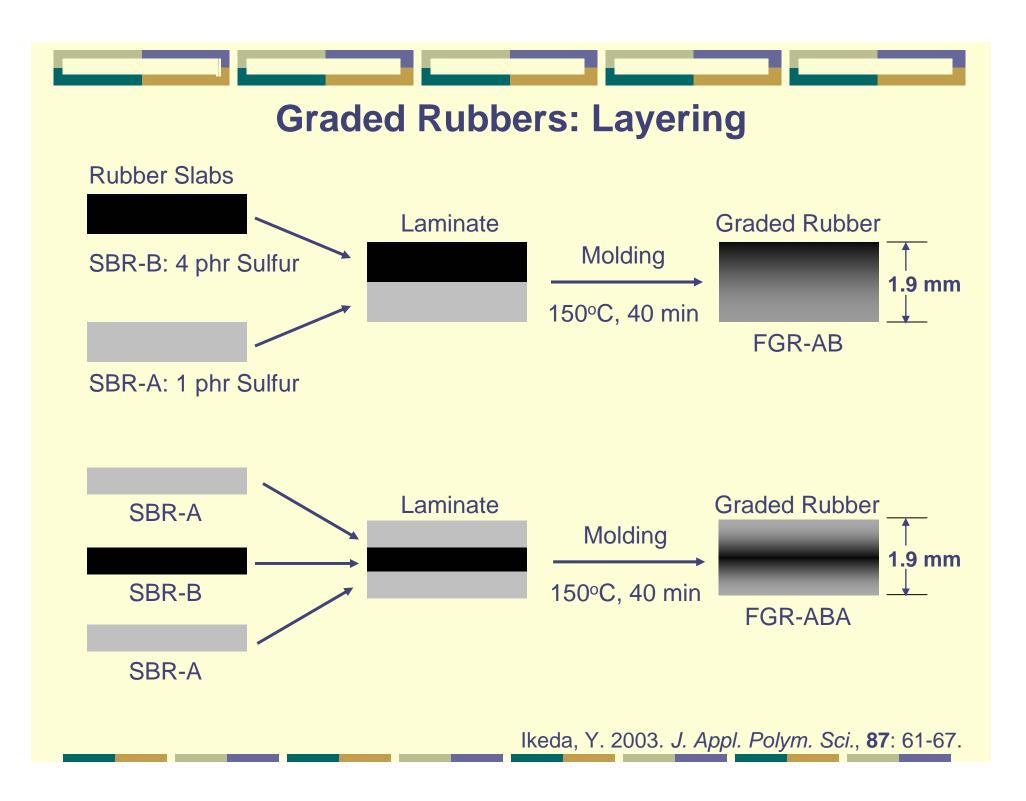


Functional Grading Concept





Tailoring the material non-homogeneity P(X) for reducing stress-strain localizations and temperature gradients controlling the dynamic properties, surface properties, and the permeability to fluids

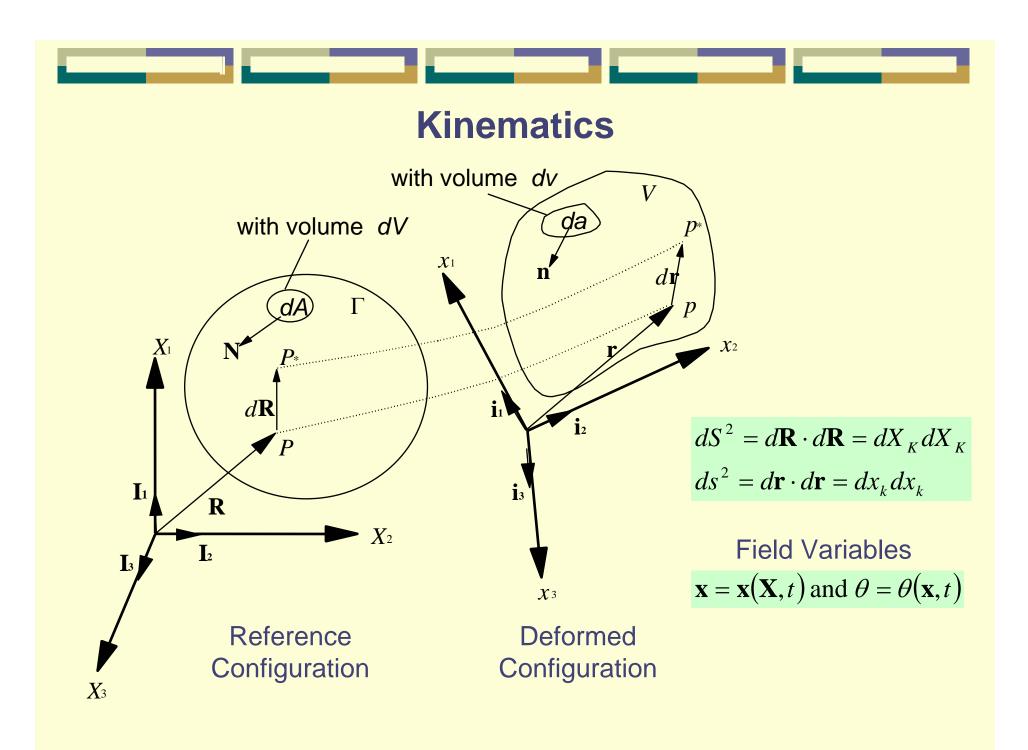


Milestones in Rubber Thermoelasticity
 Entropic origin for the stress (Anthony et al., 1942; Meyer and van der Wyk, 1946).

- Large, multi-axial, isothermal deformations (Treloar, 1944; Rivlin and Saunders, 1951)
- Universal controllable deformations of isotropic elastic materials (Ericksen, 1954, 1955).
- Universal deformations within the context of finite thermoelasticity (Petroski and Carlson, 1968,1970).
- Generalized thermoelastic models: Chadwick (1974),
 Chadwick and Creasy (1984), Ogden (1992).
- Many BVPs solved in the last decade (e.g., Horgan, Rajagopal, Saccomandi, Wineman). Non-homogeneity???

Consequences of Material Non-Homogeneity

- Material homogeneity is simply assumed in the rheological characterization of rubbers.
- What do standard tests really measure?
- Three fundamental issues to be addressed:
 - isothermal/non-isothermal behavior of non-homogeneous rubbers (forward problem)
 - tailoring the non-homogeneity
 - characterization of non-homogeneous rubbers (inverse problem)



Deformation Measures

The deformation gradient F

 $\mathbf{F} = \nabla \mathbf{x} \equiv \partial \mathbf{x} / \partial \mathbf{X}$

The left (right) Cauchy-Green deformation tensor B (C) is obtained via polar decomposition of F:

$$\mathbf{F} = \mathbf{R}\mathbf{U} = \mathbf{V}\mathbf{R} \Rightarrow \mathbf{B} \equiv \mathbf{V}^2 = \mathbf{F}\mathbf{F}^T$$
 and $\mathbf{C} \equiv \mathbf{U}^2 = \mathbf{F}^T\mathbf{F}$,

Material isotropy requires that the response functions depend on the principal invariants of **B**, i.e., *I*₁, *I*₂, and *I*₃:

$$I_1 = \text{tr} \mathbf{B}, \quad I_2 = \frac{1}{2} \left[(\text{tr} \mathbf{B})^2 - \text{tr} \mathbf{B}^2 \right], \quad I_3 = \frac{1}{6} \left[(\text{tr} \mathbf{B})^3 - 3 \text{tr} \mathbf{B} \text{tr} \mathbf{B}^2 + 2 \text{tr} \mathbf{B}^3 \right]$$

Finite Thermoelasticity

Thermo-mechanical behavior of non-homogeneous, isotropic

non-linearly elastic materials can be described by

$$\psi = \psi(I_1, I_2, I_3, \theta, \mathbf{X}), \quad \eta = -\partial \psi / \partial \theta, \quad \varepsilon = \psi + \theta \eta$$

The Cauchy stress T and the heat flux vector q are given by

$$\mathbf{T} = \frac{2\rho_0(\mathbf{X})}{\sqrt{I_3}} \left[\left(I_2 \frac{\partial \psi}{\partial I_2} + I_3 \frac{\partial \psi}{\partial I_3} \right) \mathbf{I} + \frac{\partial \psi}{\partial I_1} \mathbf{B} - I_3 \frac{\partial \psi}{\partial I_2} \mathbf{B}^{-1} \right]$$

$$\mathbf{q} = \mathbf{K} \nabla \theta = \left(\Gamma_0 \mathbf{I} + \Gamma_1 \mathbf{B} + \Gamma_{-1} \mathbf{B}^{-1} \right) \nabla \theta, \quad \Gamma_i = \Gamma_i \left(I_1, I_2, I_3, I_4, I_5, I_6, \theta, \mathbf{X} \right)$$

where

$$i \equiv 0, 1, \text{ or} - 1, \quad I_4 = \nabla \theta \cdot \nabla \theta, \quad I_5 = \nabla \theta \cdot \mathbf{B} \nabla \theta, \quad I_6 = \nabla \theta \cdot \mathbf{B}^{-1} \nabla \theta$$

 ψ , η , and ε : specific Helmholtz free energy, entropy, and internal energy Γ_i : thermo-mechanical response functions ρ_0 and θ : a reference density and temperature

Entropic Finite Thermoelasticity

A constrained thermoelastic model:

$$\psi = \psi(I_1, I_2, I_3, \theta, \mathbf{X}) = \frac{\theta}{\rho_0(\mathbf{X})\theta_0} W(I_1, I_2, \mathbf{X}) - \int_{\theta_0}^{\theta} \left(\frac{\theta}{\overline{\theta}} - 1\right) c(\mathbf{X}, \overline{\theta}) d\overline{\theta} - \frac{\pi}{\rho_0(\mathbf{X})} \left[\sqrt{I_3} - \frac{1}{\phi(\theta)}\right]$$

The Cauchy stress T and the heat flux vector q are given by

$$\mathbf{T} = -p\mathbf{I} + 2\phi(\theta)\frac{\theta}{\theta_0} \left[\frac{\partial W}{\partial I_1}\mathbf{B} - \frac{1}{\phi^2(\theta)}\frac{\partial W}{\partial I_2}\mathbf{B}^{-1}\right]$$
$$\mathbf{q} = \left[K_0 - K_1 + K_2(\theta - \theta_0)\right]\mathbf{I} + K_1\mathbf{B}$$

Clausius-Duhem Inequalities:

$$\sqrt{I_3}\mathbf{q}^T \nabla \theta \ge 0 \Longrightarrow \Gamma_0 I_4 + \Gamma_1 I_5 + \Gamma_{-1} I_6 \ge 0$$

W: isothermal strain energy function; Constraint: $\rho/\rho_0 = \phi(\theta) = 1/\sqrt{I_3}$

Local Balance Equations

Balance of Mass:

$$\frac{\mathrm{D}\,\rho}{\mathrm{D}\,t} + \rho\nabla\cdot\mathbf{v} = 0 \quad \text{or} \quad \rho(\mathbf{x},t) = \frac{\rho_{\mathrm{o}}(\mathbf{X})}{\det\mathbf{F}} = \frac{\rho_{\mathrm{o}}(\mathbf{X})}{\sqrt{\mathrm{I}_{3}}}$$

Balance of Linear Momentum:

$$\rho \frac{\mathrm{D} \mathbf{v}}{\mathrm{D} t} - \nabla \cdot \mathbf{T} - \rho \mathbf{b} = \mathbf{0}$$

Balance of Thermal Energy:

$$\rho \frac{\mathrm{D}\varepsilon}{\mathrm{D}t} - \mathbf{T} : \mathbf{d} - \nabla \cdot \mathbf{q} - \rho s = 0$$

Isothermal Deformation of Homogeneous Rubbers

Neo-Hookean model: Mooney model: Yeoh model:

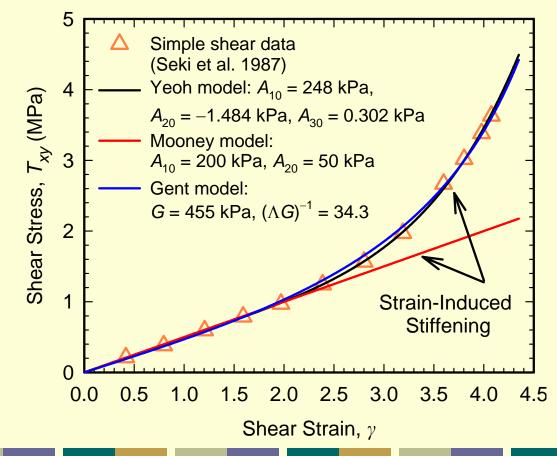
Gent model:

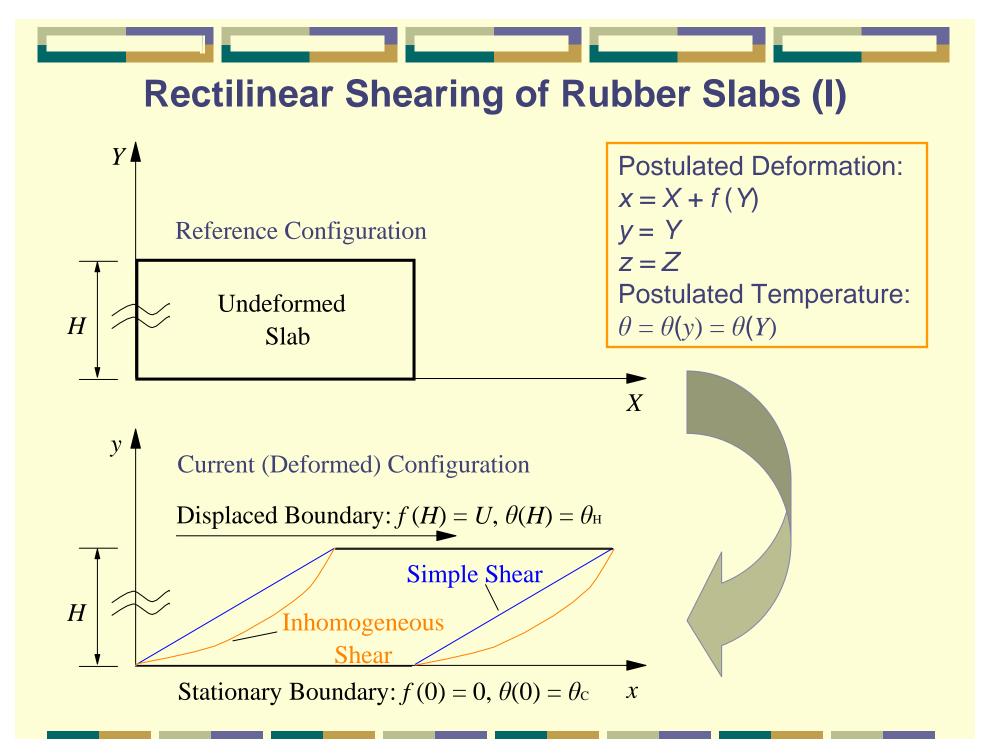
$$W(I_{1}) = A_{10}(I_{1} - 3)$$

$$W(I_{1}, I_{2}) = A_{10}(I_{1} - 3) + A_{01}(I_{2} - 3)$$

$$W(I_{1}) = A_{10}(I_{1} - 3) + A_{20}(I_{1} - 3)^{2} + A_{30}(I_{1} - 3)^{3}$$

$$W(I_{1}) = -\mu J_{m} \ln[1 - (I_{1} - 3)/J_{m}]/2$$





Rectilinear Shear (II): Deformation Measures

The deformation measures are implicitly given by

$$\mathbf{F} = \frac{\partial \mathbf{x}}{\partial \mathbf{X}} = \begin{bmatrix} 1 & f' & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{bmatrix}, \quad \mathbf{B} = \mathbf{F}\mathbf{F}^{\mathrm{T}} = \begin{bmatrix} 1 + (f')^{2} & f' & 0 \\ f' & 1 & 0 \\ 0 & 0 & 1 \end{bmatrix}$$

$$I_1 = tr B = 3 + (f')^2, \quad I_2 = \frac{1}{2} [(tr B)^2 - tr B^2] = 3 + (f')^2, \quad I_3 = det B = 1$$

Due to det F = 1, incompressibility constraint is satisfied. The strain inhomogeneity is determined from:

$$\left\| \nabla \mathbf{F} \right\| \coloneqq \sqrt{\frac{\partial F_{iJ}}{\partial X_M}} \frac{\partial F_{iJ}}{\partial X_M} = \left| f'' \right| \qquad i \equiv x, y, z \quad \text{and} \quad J, M \equiv X, Y, Z$$

f and f': shearing displacement and shear strain

|f''|: absolute value of the shear strain gradient

Rectilinear Shear (III): Field Equations

The Cauchy's equations of motion can be recorded as

$$\nabla \cdot \mathbf{T} = \frac{\partial T_{ij}}{\partial x_j} = \frac{\partial T_{ij}}{\partial X_M} \frac{\partial X_M}{\partial x_j} = \mathbf{0} \quad i, j \equiv x, y, z \quad \text{and} \quad M \equiv X, Y, Z$$

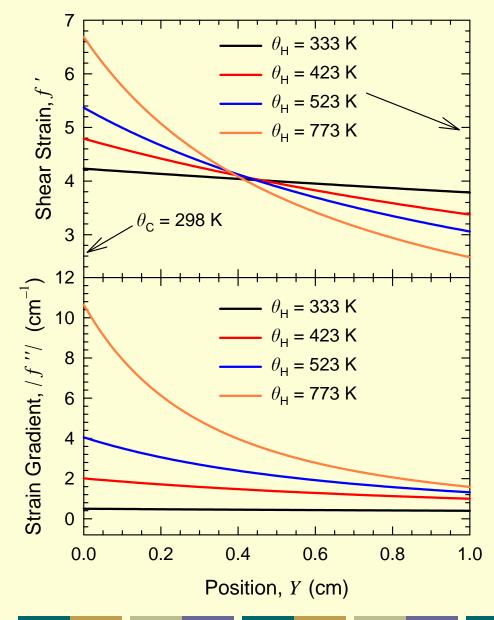
$$\frac{dT_{xy}}{dY} = \frac{d}{dY} \left[\frac{2\theta}{\theta_{o}} \left(\frac{\partial W}{\partial I_{1}} + \frac{\partial W}{\partial I_{2}} \right) f' \right] = 0, \quad f(0) = 0 \quad \text{and} \quad f(H) = U \text{ or } T_{xy}(H) = S$$

The local energy balance equation with Fourier's law:

$$\nabla \cdot \mathbf{q} = \frac{d}{dy} \left(k \frac{d\theta}{dy} \right) = k\theta'' = 0, \quad \theta(0) = \theta_{\rm C} \quad \text{and} \quad \theta(H) = \theta_{\rm H}$$

$$\theta = \theta(Y) = \theta_{\rm C} + \frac{\theta_{\rm H} - \theta_{\rm C}}{H}Y$$

Rectilinear Shear (IV): Neo-Hookean Model



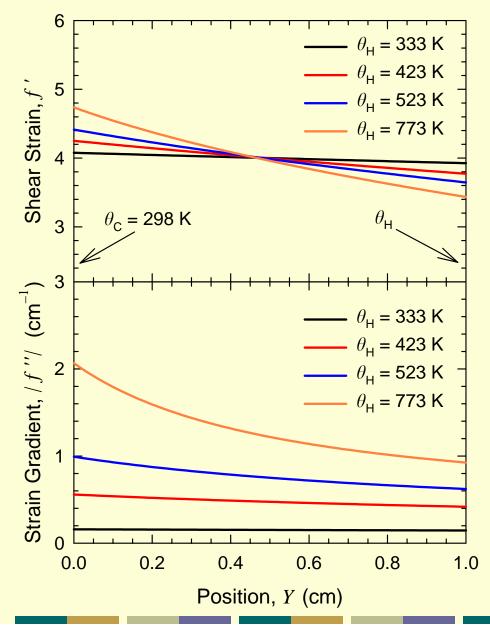
Non-isothermal Deformation of a Homogeneous Slab: $W(I_1) = A_{10}(I_1 - 3)$

An increase in temperature gradient enhances the strain inhomogeneity.

The strain inhomogeneity is more pronounced near the colder boundary.

H = 1 cm, *U* = 4 cm, $\theta_0 = 296$ K, $\theta_C = 298$ K, $A_{10} = 250$ kPa, θ_H varied

Rectilinear Shear (V): Yeoh Model

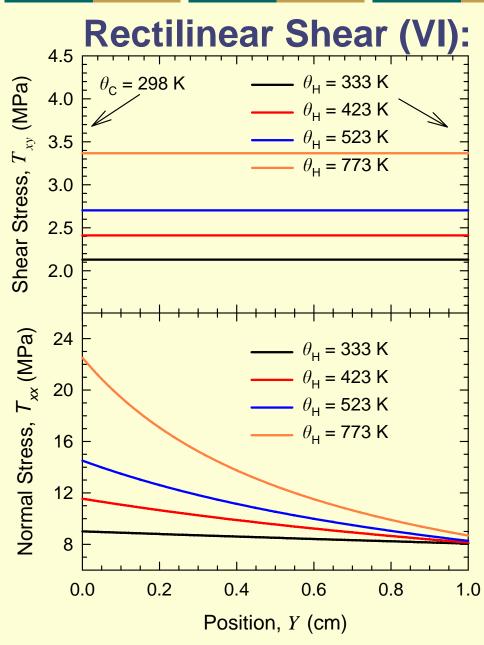


Non-isothermal Deformation of a Homogeneous Slab:

$$W(I_1) = A_{10}(I_1 - 3) + A_{20}(I_1 - 3)^2 + A_{30}(I_1 - 3)^3$$

The strain-induced stiffening homogenizes the strain field.

H = 1 cm, U = 4 cm, $\theta_0 = 296 \text{ K}, \theta_C = 298 \text{ K},$ $A_{10} = 250 \text{ kPa}, A_{10} = -1.484 \text{ kPa},$ $A_{30} = 0.302 \text{ kPa}, \theta_H \text{ varied}$



Rectilinear Shear (VI): Neo-Hookean Model

Non-isothermal Deformation of a Homogeneous Slab:

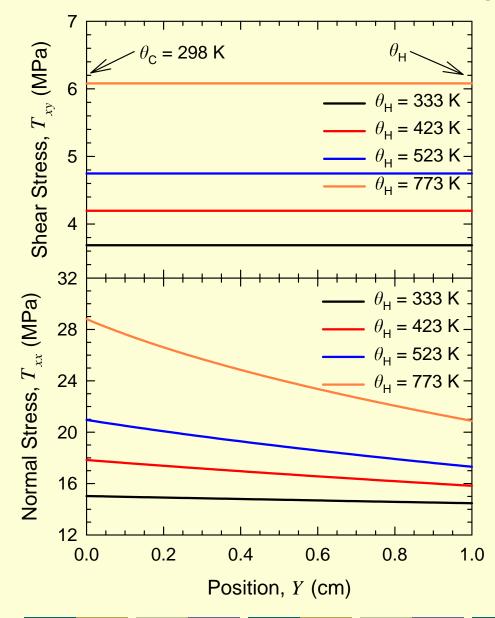
 $W(I_1) = A_{10}(I_1 - 3)$

Stress increase due to temperature-induced stiffening

Colder boundary subjected to a higher normal stress!

H = 1 cm, U = 4 cm, $\theta_0 = 296 \text{ K}, \ \theta_C = 298 \text{ K},$ $A_{10} = 250$ kPa, $\theta_{\rm H}$ varied

Rectilinear Shear (VII): Yeoh Model

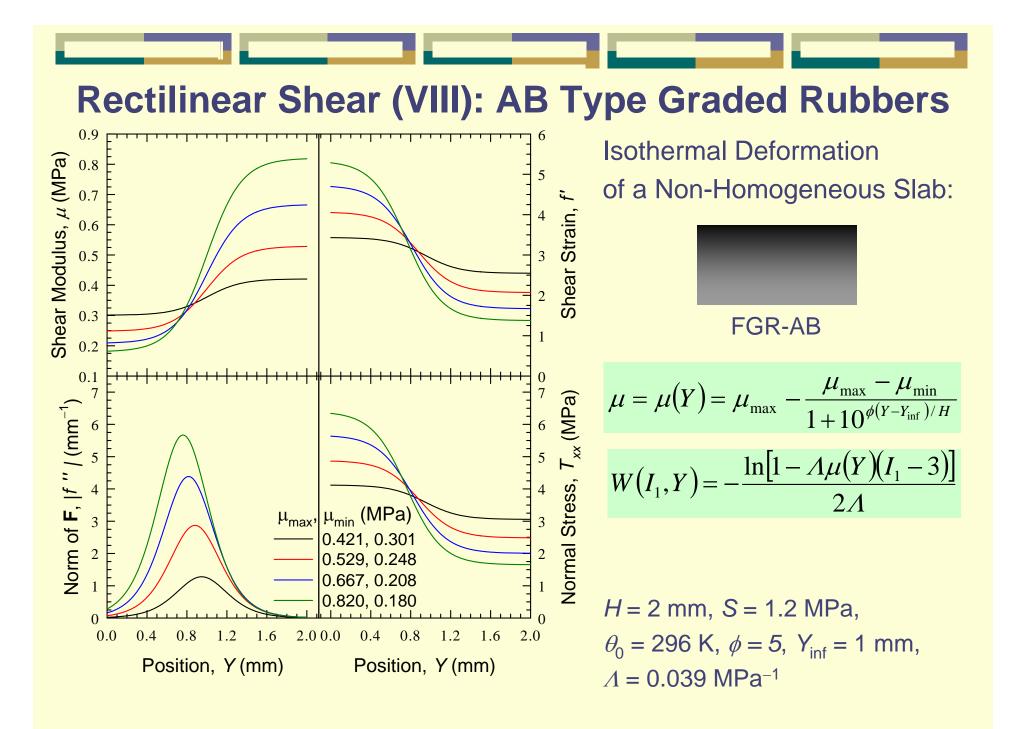


Non-isothermal Deformation of a Homogeneous Slab:

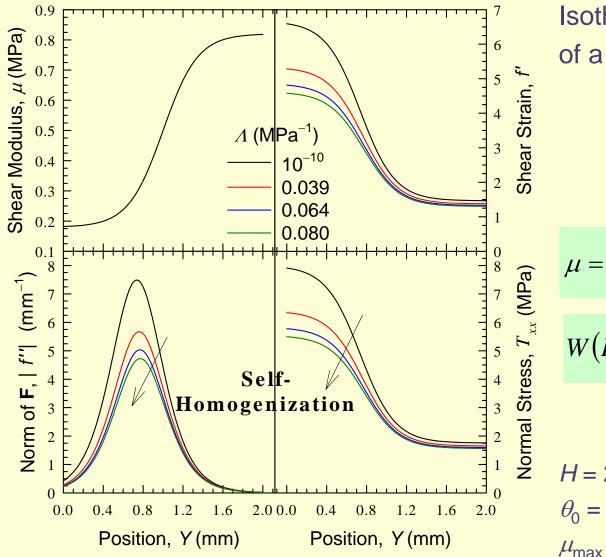
$$W(I_1) = A_{10}(I_1 - 3) + A_{20}(I_1 - 3)^2 + A_{30}(I_1 - 3)^3$$

The stresses are higher and less inhomogeneous in the Yeoh slab than in the Neo-Hookean slab.

$$\begin{split} H &= 1 \text{ cm}, \ U &= 4 \text{ cm}, \\ \theta_0 &= 296 \text{ K}, \ \theta_{\text{C}} = 298 \text{ K}, \\ A_{10} &= 250 \text{ kPa}, \ A_{10} &= -1.484 \text{ kPa}, \\ A_{30} &= 0.302 \text{ kPa}, \ \theta_{\text{H}} \text{ varied} \end{split}$$



Rectilinear Shear (IX): AB Type Graded Rubbers



Isothermal Deformation of a Non-Homogeneous Slab:

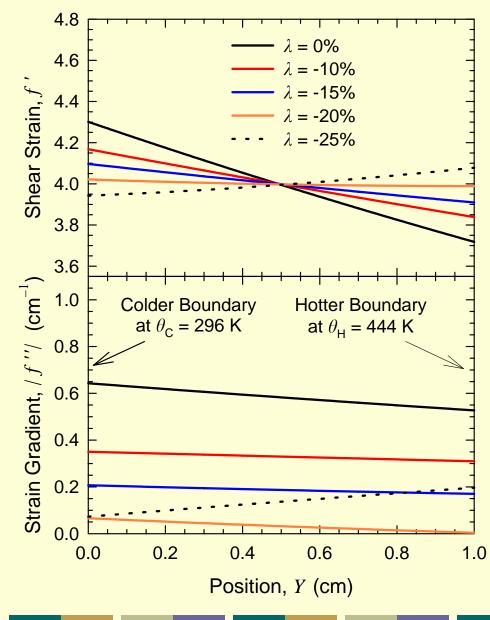


FGR-AB

$$\mu = \mu(Y) = \mu_{\max} - \frac{\mu_{\max} - \mu_{\min}}{1 + 10^{\phi(Y - Y_{\inf})/H}}$$
$$W(I_1, Y) = -\frac{\ln[1 - \Lambda\mu(Y)(I_1 - 3)]}{2\Lambda}$$

H = 2 mm, *S* = 1.2 MPa, $\theta_0 = 296$ K, $\phi = 5$, $Y_{inf} = 1$ mm, $\mu_{max} = 0.82$ MPa, $\mu_{min} = 0.18$ MPa

Rectilinear Shear (X): Functional Grading?



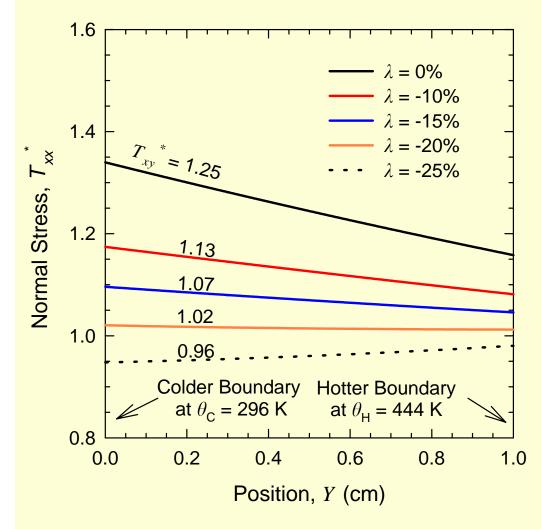
Non-isothermal Deformation of a Graded Slab:

$$\mu(Y) = \mu_{\rm B} \left[1 + \frac{\lambda (1+\kappa)Y}{100(\kappa H+Y)} \right]$$
$$W(I_1, Y) = -\frac{\ln[1 - \Lambda \mu(Y)(I_1 - 3)]}{2\Lambda}$$

An optimum grading exists!

H = 1 cm, *U* = 4 cm, $\theta_0 = \theta_C = 296$ K, $\theta_H = 444$ K, $\kappa = 5.0$, $\mu_B = 455$ kPa, $\Lambda = 0.0641$ MPa⁻¹

Rectilinear Shear (XI): Functional Grading?

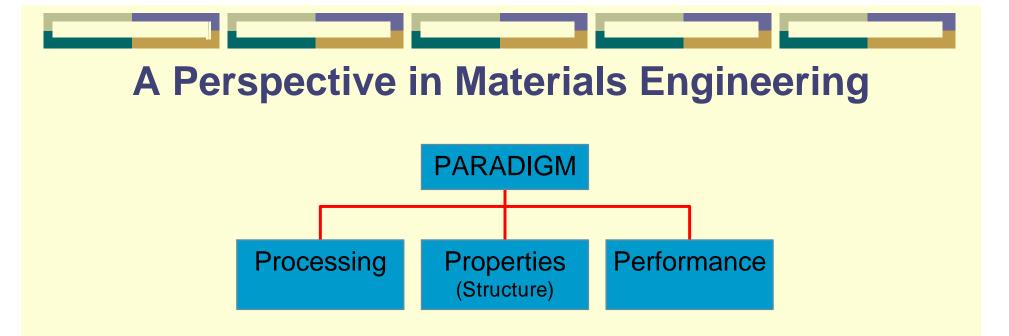


Non-isothermal Deformation of a Graded Slab:

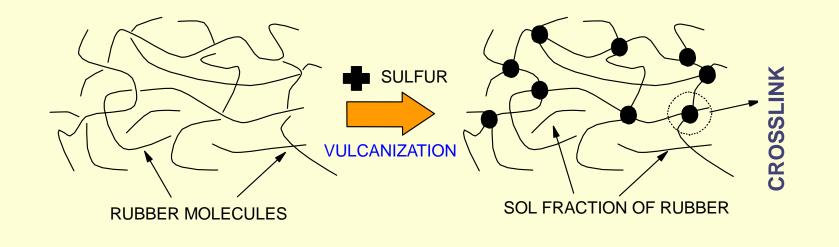
$$\mu(Y) = \mu_{\rm B} \left[1 + \frac{\lambda (1+\kappa)Y}{100(\kappa H + Y)} \right]$$
$$W(I_1, Y) = -\frac{\ln[1 - A\mu(Y)(I_1 - 3)]}{2\Lambda}$$

Grading leads to a reduction in stresses.

H = 1 cm, *U* = 4 cm, $\theta_0 = \theta_C = 296$ K, $\theta_H = 444$ K, $\kappa = 5.0$, $\mu_B = 455$ kPa, $\Lambda = 0.0641$ MPa⁻¹

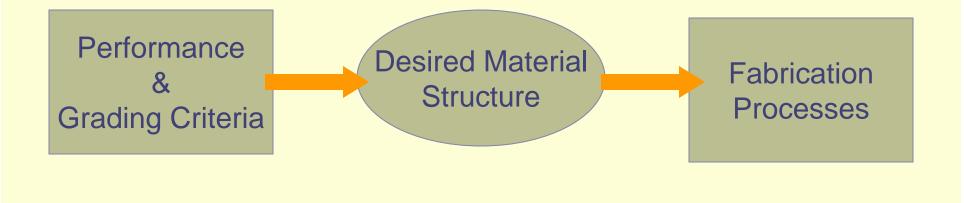


How to generate the spatial variation of the crosslink density V(X)?



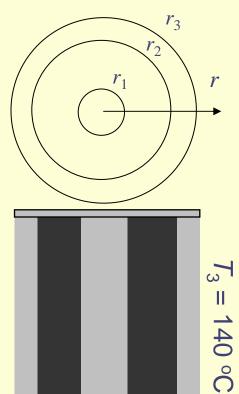
How to Produce FGREMs?

- Construction-based methods: layers with different contents of
 - curing agents (Ikeda, 2003)
 - secondary phase (fillers, recycled rubber powder)
 - rubber blends
- Transport-based methods:
 - Thermally-induced structure (crosslink density)



Molding of a Rubber Tube (I)

Geometry



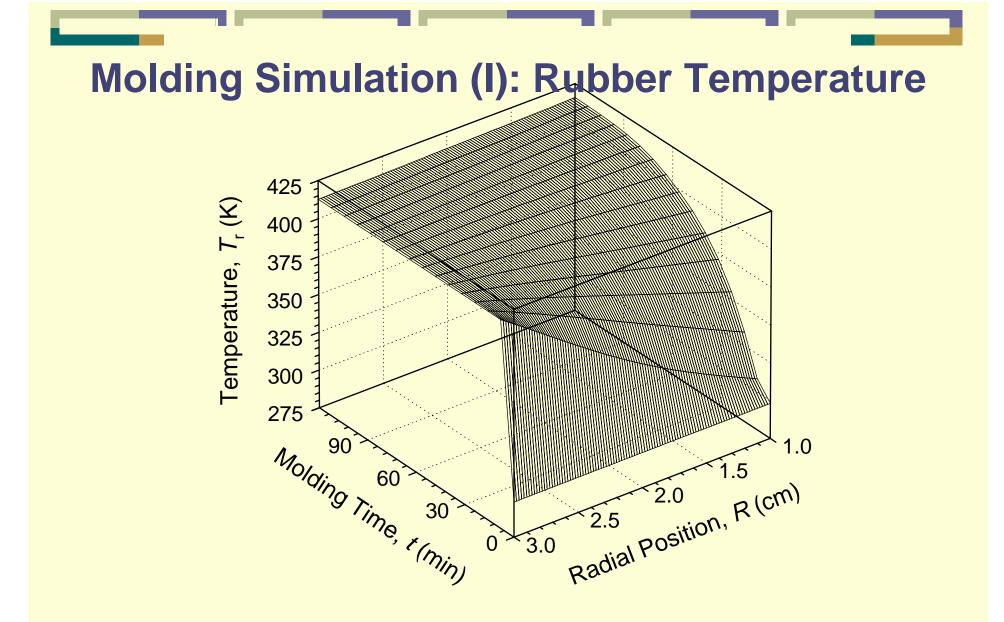
Rubber: black, Gray: Metal mold

Rubber with curing agents is heated by the steam-jacketed metal mold.

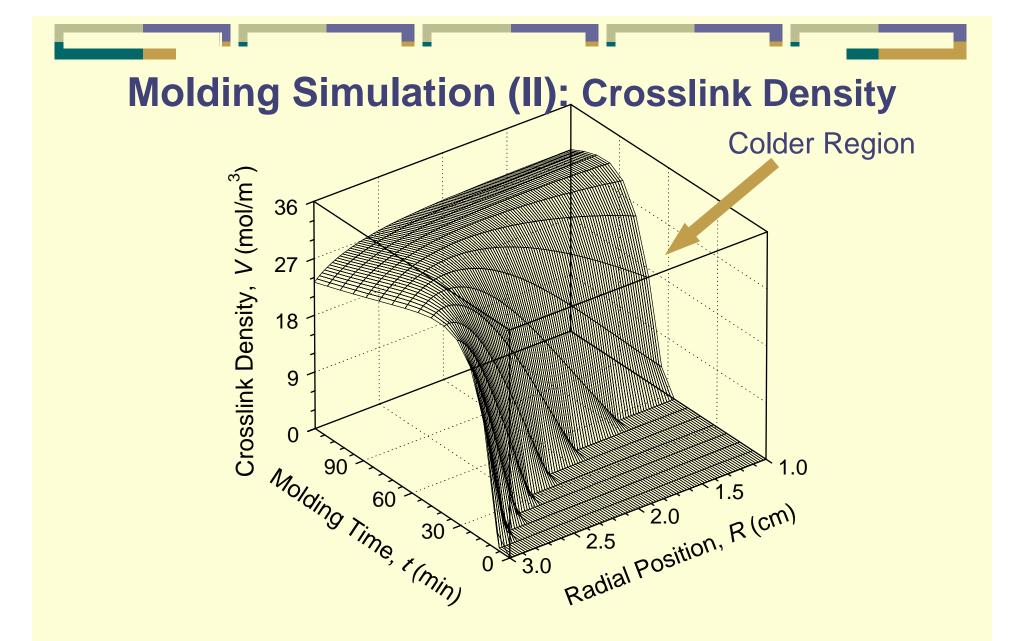
The spatio-temporal development of temperature and crosslink density was simulated.

• Long tube, insulated ends, constant physical properties, except $\lambda_r = \lambda_0 - \lambda_1 (T_r - 273.15)$

(E. Bilgili, AIChE Annu. Meet., 2003)



The rubber temperature rises fast near the outer mold (about 413 K), which is heated by steam. It monotonically increases and evolves to a steady state.



The evolution of crosslink density *V* is non-monotonic due to reversion. Transport-based grading of rubbers is possible via control of the process.

Conclusions

- Rubbers can be graded in a variety of ways.
- Functional grading can reduce stress-strain localizations.
- Sophisticated process models enable us to generate a desired structure.
- An attempt to connect processing-properties-performance aspects of rubbers has been made.

Publications

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E. Bilgili, "A Parametric Study of the Circumferential Shearing of Rubber Tubes: Beyond Isothermality and Material Homogeneity," *Kaut. Gummi Kunstst.*, 56, 671–676 (2003).

E. Bilgili, "Controlling the Stress–Strain Inhomogeneities in Axially Sheared and Radially Heated Hollow Rubber Tubes via Functional Grading," *Mech. Res. Commun.* 30, 257–266 (2003).

E. Bilgili, B. Bernstein, H. Arastoopour, "Effect of Material Non-homogeneity on the Inhomogeneous Shearing Deformation of a Gent Slab Subjected to a Temperature Gradient," *Int. J. Non-Linear Mech.* 38, 1351–1368 (2003).

E. Bilgili, "Computer Simulation as a Tool to Investigate the Shearing Deformation of Non-Homogeneous Elastomers," *J. Elastom. Plast.* 34, 239–264 (2002).

E. Bilgili, B. Bernstein, H. Arastoopour, "Influence of Material Non-Homogeneity on the Shearing Response of a Neo-Hookean Slab," *Rubber Chem. Technol.* 75, 347–363 (2002).

E. Bilgili, B. Bernstein, H. Arastoopour, "Inhomogeneous Shearing Deformation of a Rubber-like Slab within the Context of Finite Thermoelasticity with Entropic Origin for the Stress," *Int. J. Non-Linear Mech.* 36, 887–900 (2001).

Acknowledgments

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