



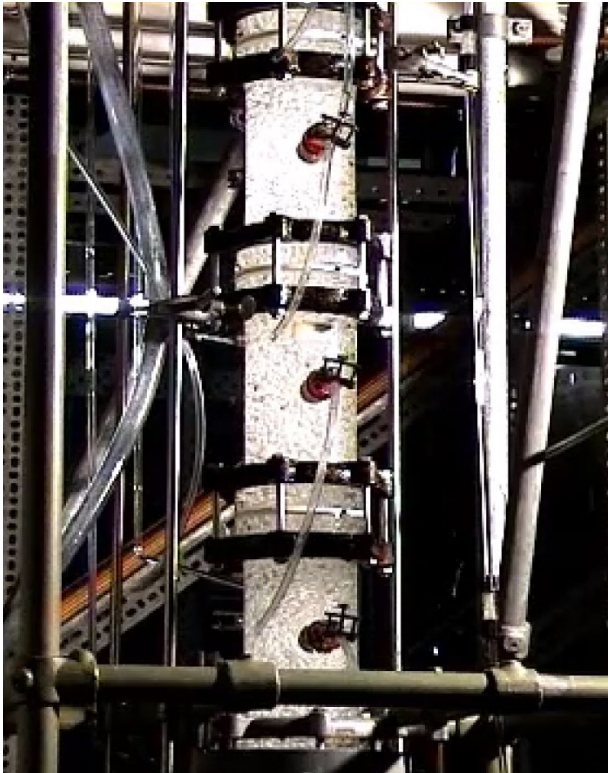
Sharp-interface modeling of mass transfer in multicomponent two-phase fluid systems

Dieter Bothe

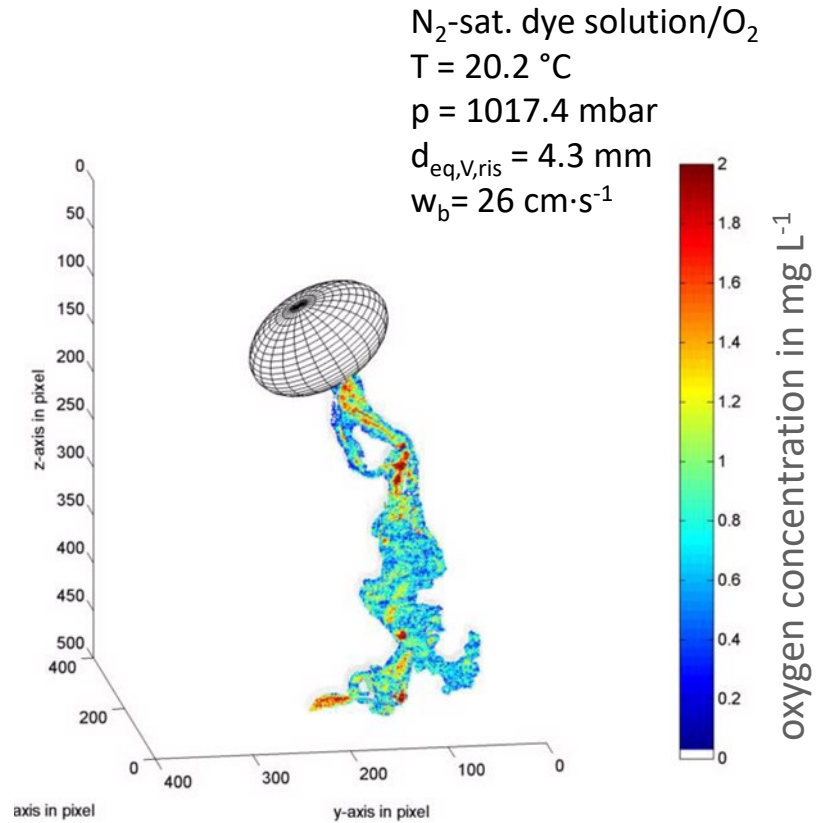


Mathematical
Modeling and Analysis

Motivation: Reactive Mass Transfer



Chemical Engineering Lab, U Paderborn



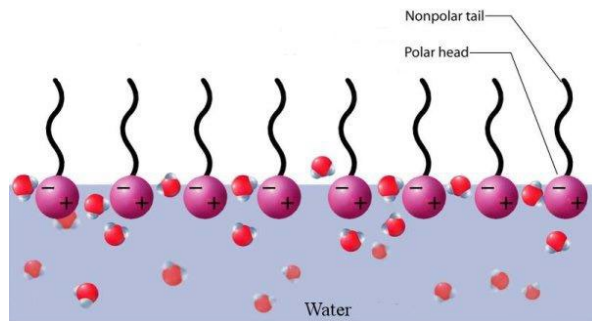
Time Resolved Scanning Laser Induced Fluorescence

Michael Schlüter, IMS, TU Hamburg-Harburg

Rüttinger, Spille, Hoffmann, Schlüter,
ChemBioEng reviews 5.4 (2018): 253-269

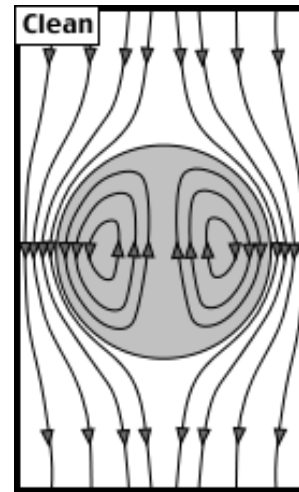
Surface Active Agents

surfactant = surface active agent

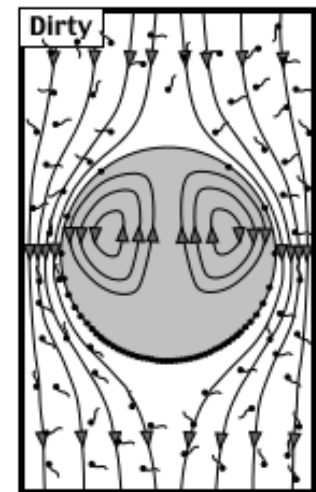


Interfacial concentration c_i^Σ

<https://commons.wikimedia.org/wiki/File:Surfactant.jpg>

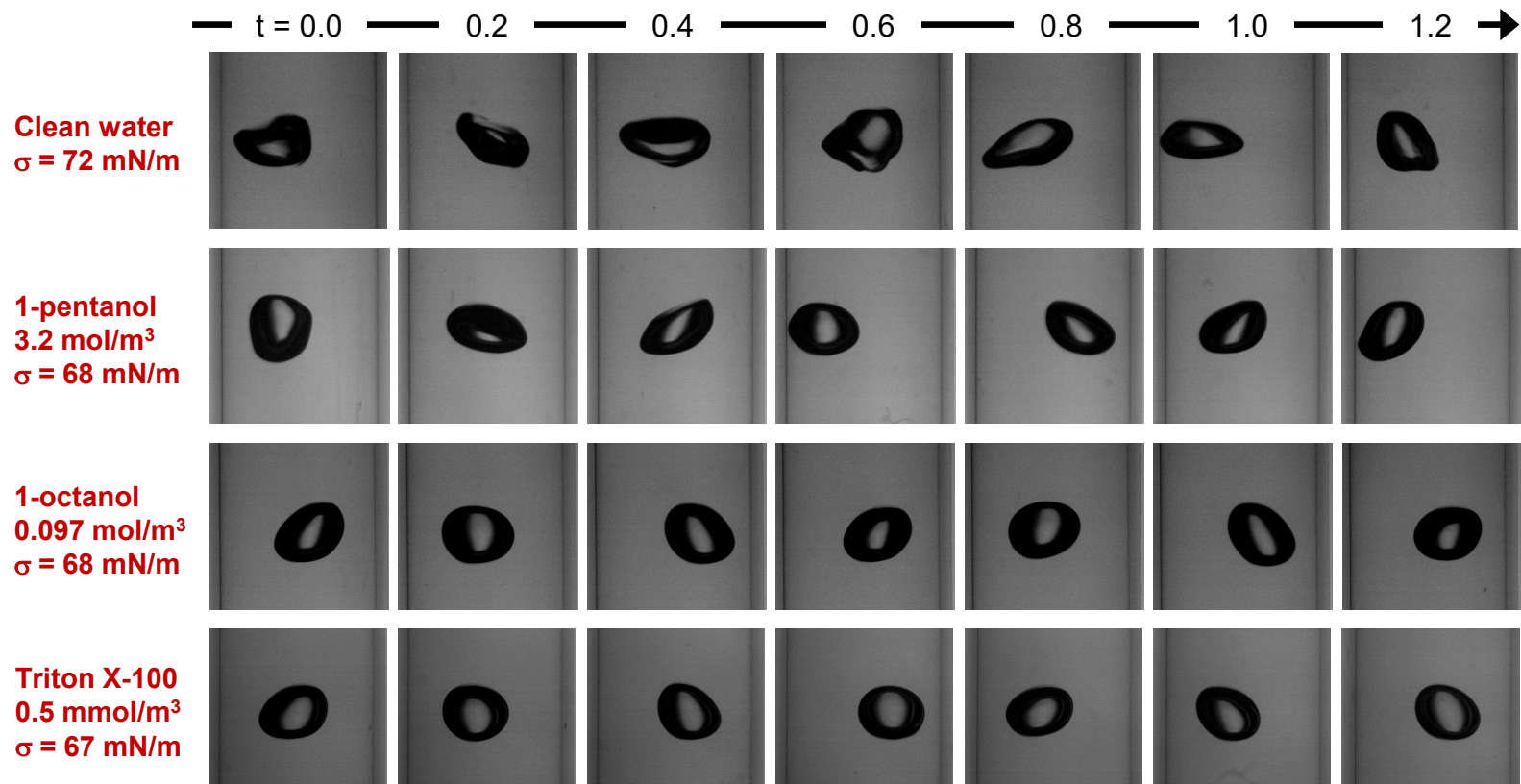


<http://www.bubbleology.com>



- **Surface equation of state:** $\sigma = \sigma(T, c_1^\Sigma, \dots, c_N^\Sigma)$
- **Surface tension gradients induce Marangoni stress**
- **Surface coverage hinders mass transfer**

Dissolving CO₂ Bubbles in Surfactant Solutions



A. Tomiyama et al.,
Kobe University

Experimental Findings by Akio Tomiyama et al.

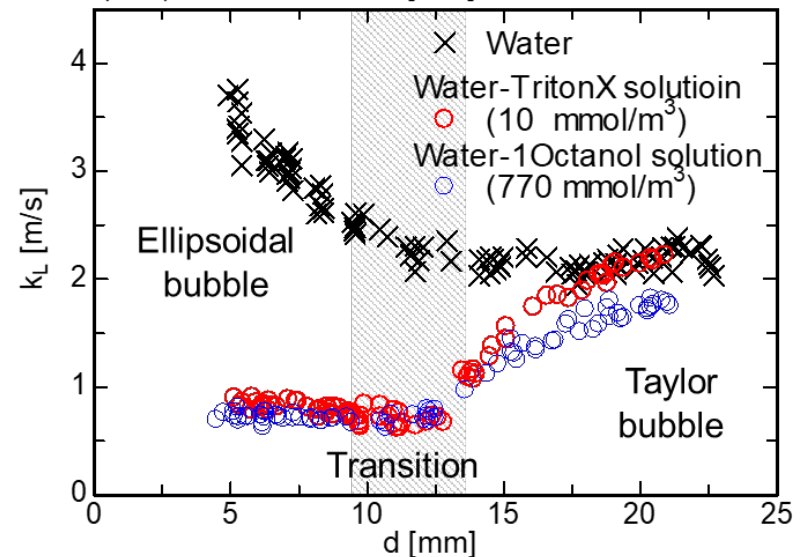
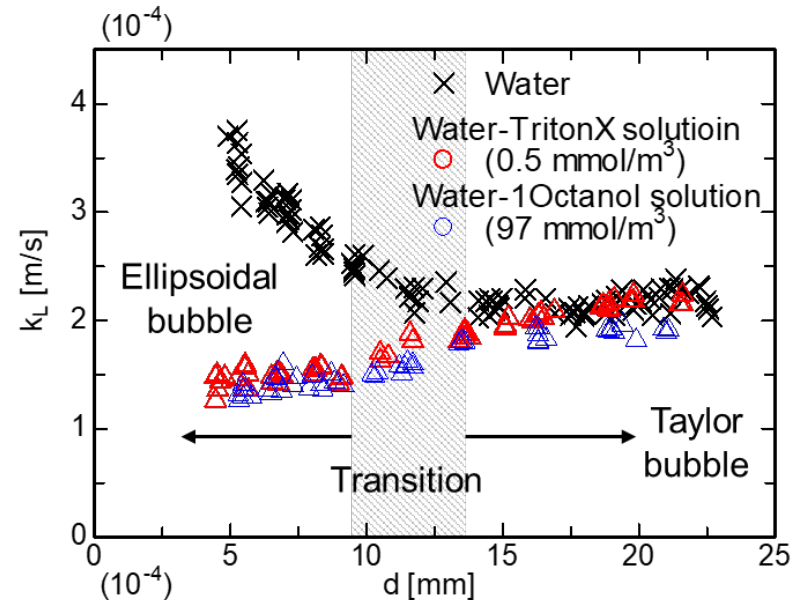
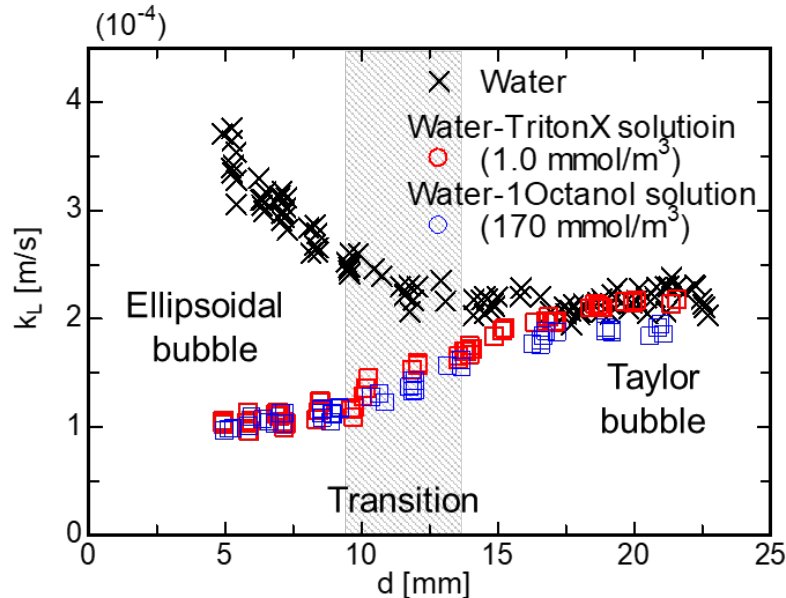
Triton X-100

C_{sol} [mmol/m ³]	σ [mN/m]
0.5	68
1.0	65
10	52

1-Octanol

C_{sol} [mmol/m ³]	σ [mN/m]
97	68
170	65
770	52

(C_{sol} : Surfactant concentration, σ : Surface tension)



Effects of surface active agents on mass transfer

- **experimental findings**

- surfactant coverage leads to local mass transfer hindrance
- mass transfer reduction depends on the *reduction of surface tension* rather than on the surfactant concentration (type of surfactant seems irrelevant)

$$\dot{m}_i^{\text{contam}} = \alpha(\sigma) \dot{m}_i^{\text{clean}}$$

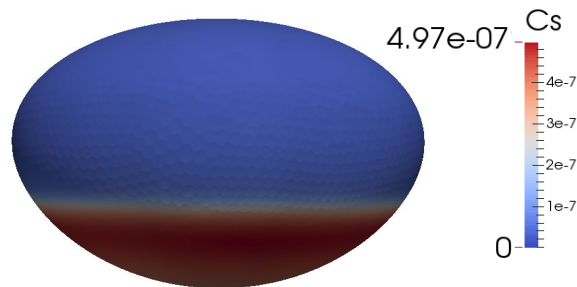
**Thermodynamically consistent continuum model
of hindered mass transfer?**

J. Aoki, et al., Tomiyama, Chem. Eng. Techn. 38 (2015)

Y. Hori, et al., Tomiyama, Int. J. Heat Mass Trans. 136 (2019)

Effects of surface active agents on mass transfer

- *hindrance* of mass transfer due to *surface coverage*



Surfactant surface concentration
 $[C_s] = \text{mol/m}^2$, $t = 0.08 \text{ s}$.

surface coverage ratio

$$Se = \frac{\Gamma_{eq}}{\Gamma_{\infty}} = \frac{\beta C_{sol}}{\alpha + \beta C_{sol}}$$

k_L^0	k_L	k_L^1
$Se = 0$	$0 < Se < 1$	$Se = 1$

Sardeing model : $Sh = (1 - Se)Sh_{clean} + SeSh_{dirty}$

$$Sh_{clean} = 0.42[\Delta\rho g d^3 / \rho_L v_L^2]^{1/3} Sc^{1/2}$$

$$Sh_{dirty} = 0.54(\alpha / \beta)^{0.0837} (g d^3 / v_L^2)^{1/3} Sc^{1/3}$$

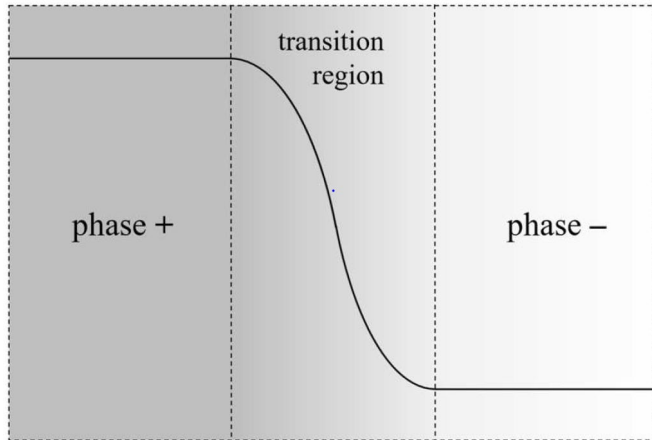
R. Sardeing et al., Chem. Eng. Sci. 61 (2006)

No quantitative match with experimental data (A. Tomiyama, Kobe Univ.)

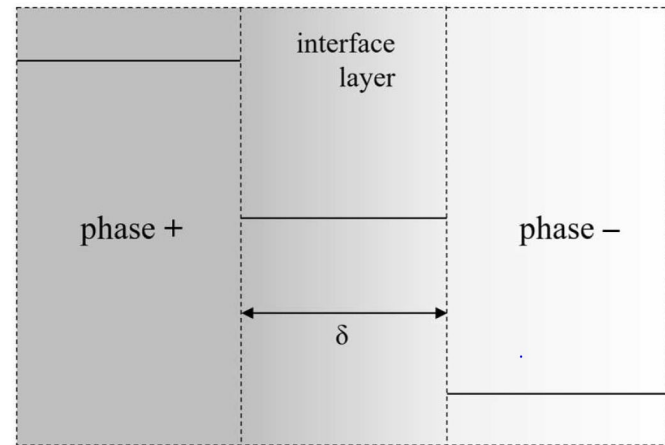
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- I – Motivation
 - II – Sharp-Interface Framework**
 - III – Interfacial Entropy Principle
 - IV – Mass Transfer Closure with
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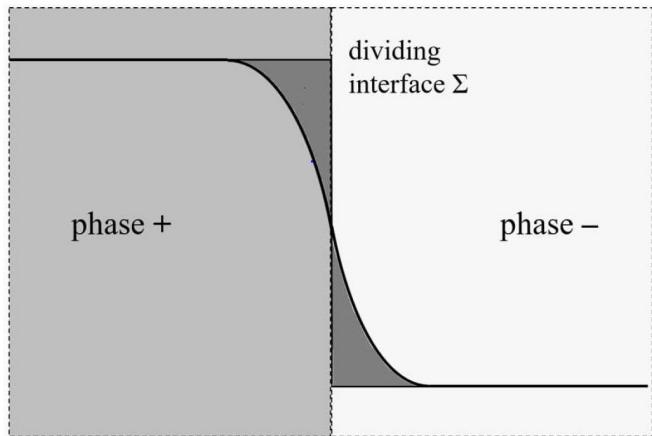
Various Interface Conceptions



Diffuse interface



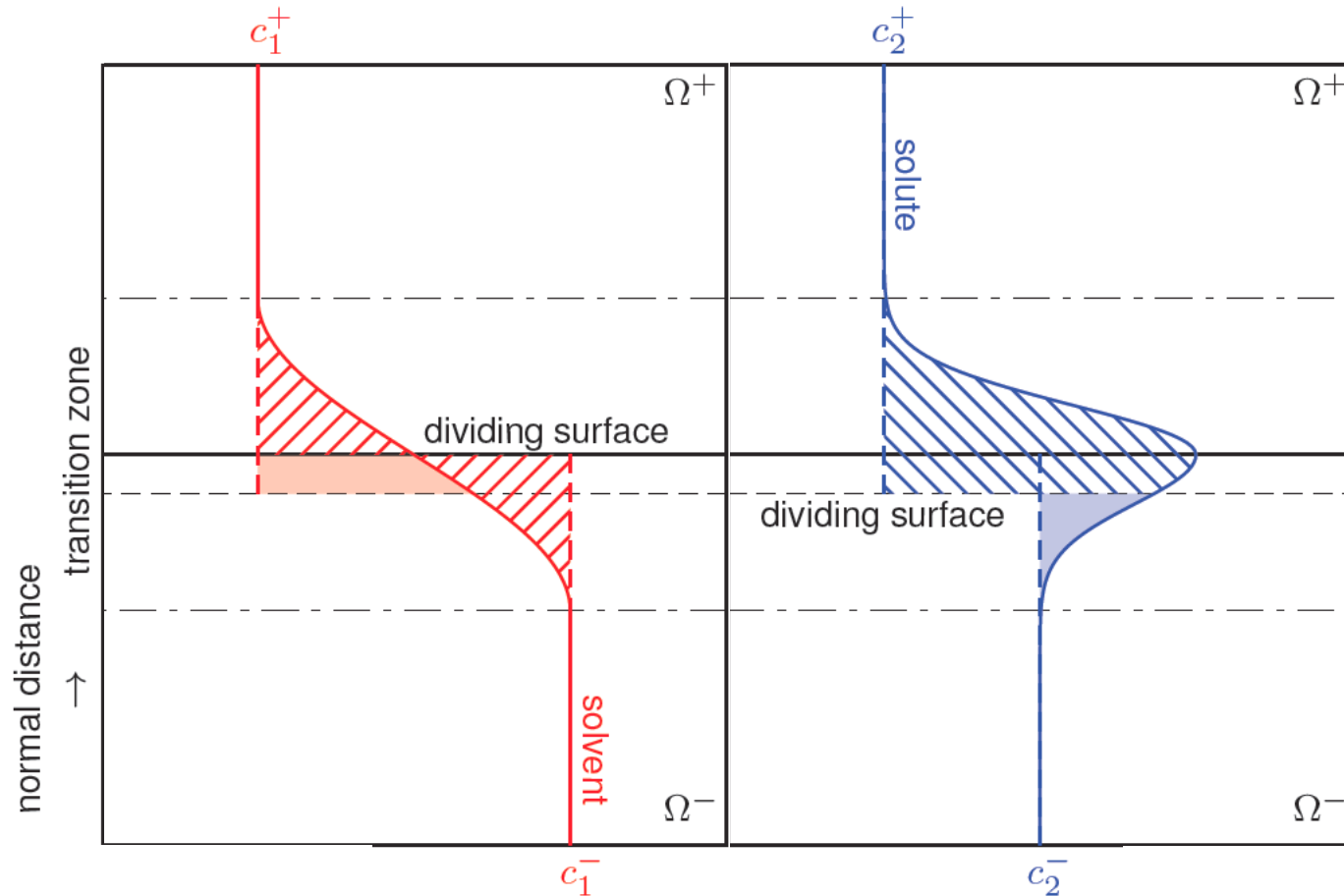
Guggenheim's interface layer



Gibbs' dividing interface

Choosing the Interface Position

Gibbs' construction of a dividing interface

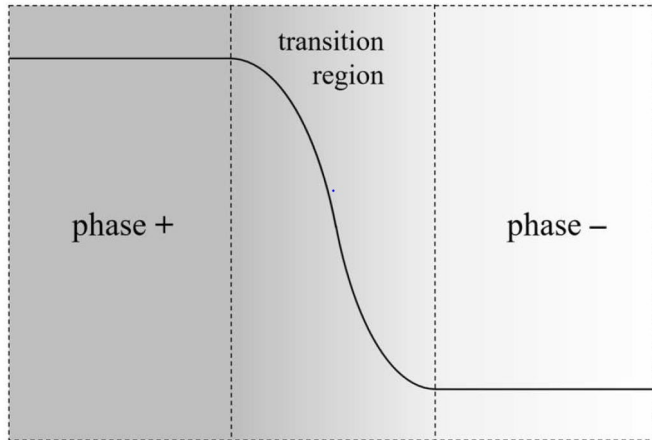


- replace steep gradient by discontinuous field
- introduce excess quantities to conserve (partial) mass

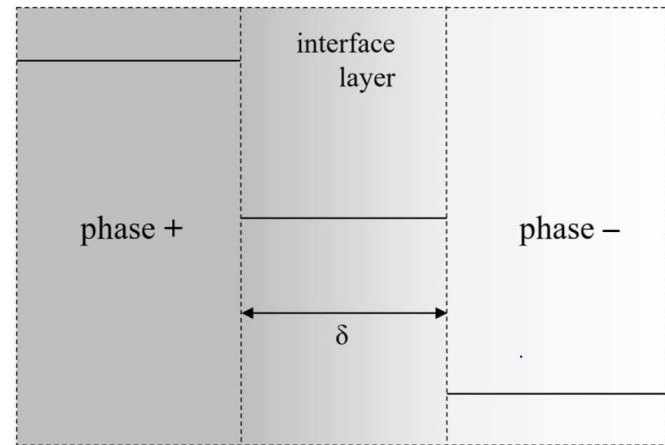
Choosing the Interface Concept

- **Gibbs' construction tricky for multicomponent systems**
 - excess concentrations can become negative
 - components appear/disappear on Σ depending on position of Σ
 - modified construction of van den Tempel / Lucassen-Reynders: positive excess, but interface can be shifted far away
- **Is a transfer species like CO_2 present on the interface ?**
 - answer depends on the modeling approach !
 - in diffuse interface models, any species is present in the transition layer
 - unspecified in the Gibbs construction
 - model by Guggenheim keeps the interface layer

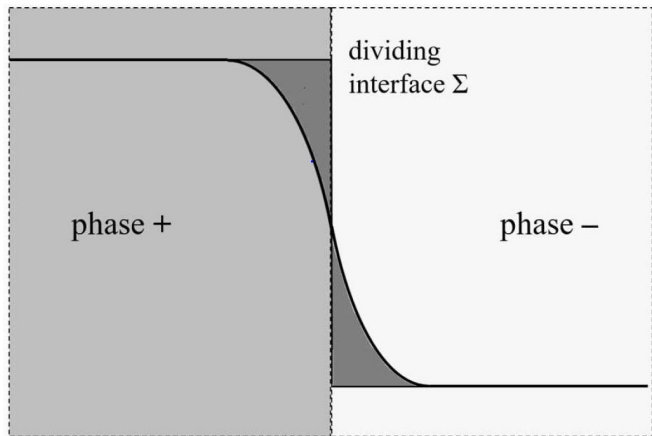
Various Interface Conceptions



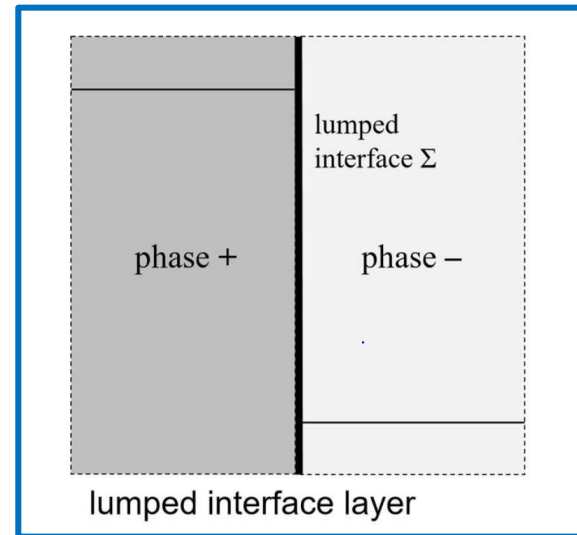
Diffuse interface



Guggenheim's interface layer



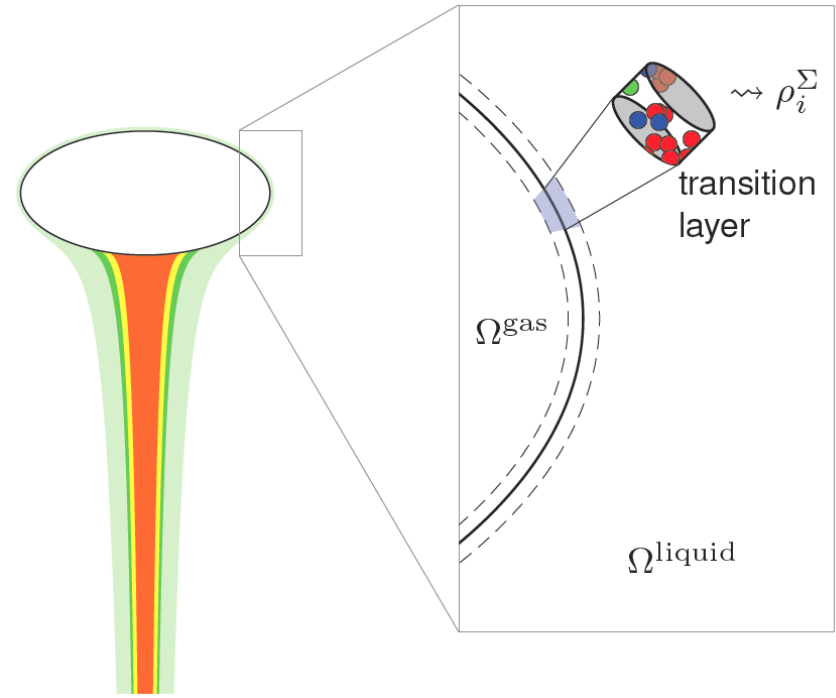
Gibbs' dividing interface



lumped interface layer

Effects of surface active agents on mass transfer

- sharp-interface continuum thermodynamics of multi-component fluid systems with ***non-vanishing interfacial mass densities*** for all constituents !



Aim: Thermodynamically consistent theory of local mass transfer resistance

D. Bothe: Sharp-interface continuum thermodynamics of multicomponent fluid systems with interfacial mass. Int. Journal of Engineering Science **179**, 103731 (2022).

D. Bothe: Multi-velocity sharp-interface continuum thermodynamics of fluid systems with adsorption. arXiv:2502.00906 (2025).

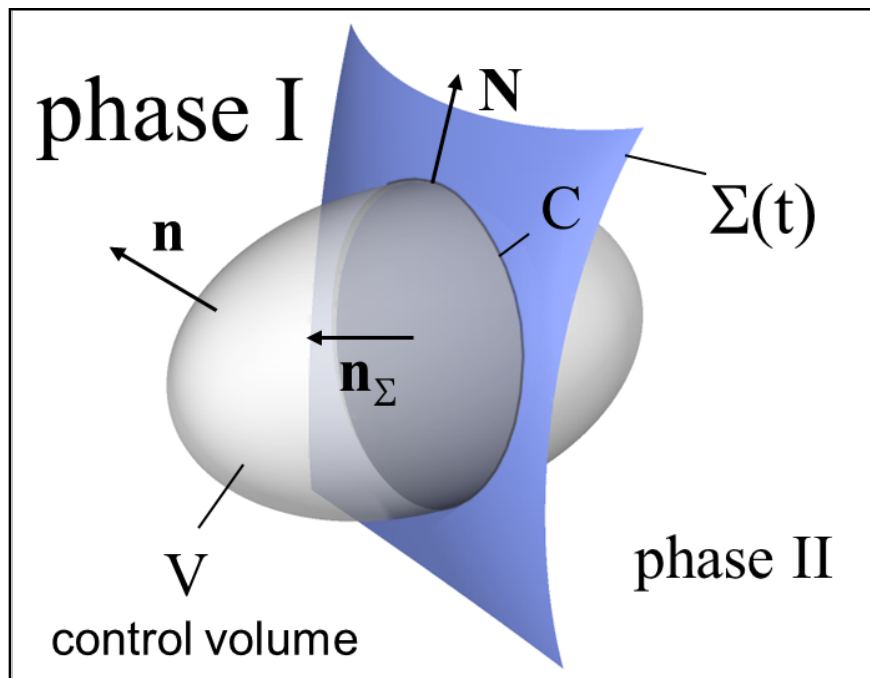
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Continuum Physical Modeling

Two-phase integral balance equations

$$\frac{d}{dt} \left(\int_V \phi \, dx + \int_{\Sigma_V} \phi^\Sigma \, do \right) = - \int_{\partial V} \mathbf{j} \cdot \mathbf{n} \, do + \int_V f \, dx - \int_{\partial \Sigma_V} \mathbf{j}^\Sigma \cdot \mathbf{N} \, ds + \int_{\Sigma_V} f^\Sigma \, do$$



- $\Omega^+(t), \Omega^-(t)$ bulk phases
- $\mathbf{n}_\Sigma$ unit interface normal (into Ω^-)
- $\mathbf{n}$ outer unit normal to V
- $\Sigma(t)$ interface, $\Sigma_V(t) := \Sigma(t) \cap V$
- $\partial \Sigma_V$ boundary curve of Σ_V
- $\mathbf{N}$ outer unit normal to $\partial \Sigma_V$, $\mathbf{N} \perp \mathbf{n}_\Sigma$

Balance equations for mass, momentum, energy

partial mass balance

$$\partial_t \rho_i + \operatorname{div} (\rho_i \mathbf{v} + \mathbf{j}_i) = M_i r_i$$

$$\partial_t^\Sigma \rho_i^\Sigma + \operatorname{div}_\Sigma (\rho_i^\Sigma \mathbf{v}^\Sigma + \mathbf{j}_i^\Sigma) + \llbracket \rho_i (\mathbf{v} - \mathbf{v}^\Sigma) + \mathbf{j}_i \rrbracket \cdot \mathbf{n}_\Sigma = M_i r_i^\Sigma$$

momentum balance

$$\partial_t (\rho \mathbf{v}) + \operatorname{div} (\rho \mathbf{v} \otimes \mathbf{v} - \mathbf{S}) = \rho \mathbf{b},$$

$$\partial_t^\Sigma (\rho^\Sigma \mathbf{v}^\Sigma) + \nabla_\Sigma \cdot (\rho^\Sigma \mathbf{v}^\Sigma \otimes \mathbf{v}^\Sigma - \mathbf{S}^\Sigma) + \llbracket \rho \mathbf{v} \otimes (\mathbf{v} - \mathbf{v}^\Sigma) - \mathbf{S} \rrbracket \cdot \mathbf{n}_\Sigma = \rho^\Sigma \mathbf{b}^\Sigma$$

internal energy balance

$$\partial_t (\rho e) + \nabla \cdot (\rho e \mathbf{v} + \mathbf{q}) = \nabla \mathbf{v} : \mathbf{S} + \sum_{i=1}^N \mathbf{j}_i \cdot \mathbf{b}_i$$

$$\begin{aligned} \partial_t^\Sigma (\rho^\Sigma e^\Sigma) + \nabla_\Sigma \cdot (\rho^\Sigma e^\Sigma \mathbf{v}^\Sigma + \mathbf{q}^\Sigma) + \llbracket \dot{m} \left(e + \frac{(\mathbf{v} - \mathbf{v}^\Sigma)^2}{2} \right) \rrbracket \\ - \llbracket (\mathbf{v} - \mathbf{v}^\Sigma) \cdot (\mathbf{S} \mathbf{n}_\Sigma) \rrbracket + \llbracket \mathbf{q} \cdot \mathbf{n}_\Sigma \rrbracket = \nabla_\Sigma \mathbf{v}^\Sigma : \mathbf{S}^\Sigma + \sum_{i=1}^N \mathbf{j}_i^\Sigma \cdot \mathbf{b}_i^\Sigma \end{aligned}$$

Entropy Balance for Bulk-Interface

integral entropy balance

$$\frac{d}{dt} \left(\int_{V(t)} \rho s \, dx + \int_{A(t)} \rho^\Sigma s^\Sigma \, do \right) = - \int_{\partial V(t)} \mathbf{\Phi} \cdot \mathbf{n} \, do + \int_{V(t)} \zeta \, dx$$
$$- \int_{\partial A(t)} \mathbf{\Phi}^\Sigma \cdot \mathbf{N} \, ds + \int_{A(t)} \zeta^\Sigma \, do$$

$\zeta \geq 0 \quad \zeta^\Sigma \geq 0$

local entropy balance

$$\partial_t(\rho s) + \nabla \cdot (\rho s \mathbf{v} + \mathbf{\Phi}) = \zeta \quad \text{in } \Omega \setminus \Sigma$$
$$\partial_t^\Sigma(\rho^\Sigma s^\Sigma) + \nabla_\Sigma \cdot (\rho^\Sigma s^\Sigma \mathbf{v}^\Sigma + \mathbf{\Phi}^\Sigma) + \llbracket \rho s(\mathbf{v} - \mathbf{v}^\Sigma) + \mathbf{\Phi} \rrbracket \cdot \mathbf{n}_\Sigma = \zeta^\Sigma \quad \text{on } \Sigma$$

$$\llbracket \phi \rrbracket(t, \mathbf{x}) := \lim_{h \rightarrow 0+} (\phi(t, \mathbf{x} + h \mathbf{n}_\Sigma) - \phi(t, \mathbf{x} - h \mathbf{n}_\Sigma))$$

Interfacial Entropy Production

Reduced interfacial entropy production (zero interfacial viscosities):

$$\begin{aligned}\zeta^\Sigma = & \mathbf{q}^\Sigma \cdot \nabla_\Sigma \frac{1}{T^\Sigma} - \sum_{i=1}^N \mathbf{j}_i^\Sigma \cdot \left(\nabla_\Sigma \frac{\mu_i^\Sigma}{T^\Sigma} - \frac{\mathbf{b}_i^\Sigma}{T^\Sigma} \right) - \frac{1}{T} \sum_{a=1}^{N_R^\Sigma} R_a^\Sigma \mathcal{A}_a^\Sigma \\ & - \frac{1}{T^\Sigma} (\mathbf{v}^+ - \mathbf{v}^\Sigma)_\parallel \cdot (\mathbf{S}^+ \mathbf{n}^+)_\parallel - \frac{1}{T^\Sigma} (\mathbf{v}^- - \mathbf{v}^\Sigma)_\parallel \cdot (\mathbf{S}^- \mathbf{n}^-)_\parallel \\ & + \left(\frac{1}{T^\Sigma} - \frac{1}{T^+} \right) \left(\dot{m}^{+, \Sigma} \left(e^+ + \frac{p^+}{\rho^+} \right) + \mathbf{q}^+ \cdot \mathbf{n}^+ \right) \\ & + \left(\frac{1}{T^\Sigma} - \frac{1}{T^-} \right) \left(\dot{m}^{-, \Sigma} \left(e^- + \frac{p^-}{\rho^-} \right) + \mathbf{q}^- \cdot \mathbf{n}^- \right) \\ & + \sum_{i=1}^N \dot{m}_i^{+, \Sigma} \left(\frac{\mu_i^+}{T^+} - \frac{\mu_i^\Sigma}{T^\Sigma} + \frac{1}{T^\Sigma} \left(\frac{(\mathbf{v}^+ - \mathbf{v}^\Sigma)^2}{2} - \mathbf{n}^+ \frac{\mathbf{S}^{+, \text{irr}}}{\rho^+} \mathbf{n}^+ \right) \right) \\ & + \sum_{i=1}^N \dot{m}_i^{-, \Sigma} \left(\frac{\mu_i^-}{T^-} - \frac{\mu_i^\Sigma}{T^\Sigma} + \frac{1}{T^\Sigma} \left(\frac{(\mathbf{v}^- - \mathbf{v}^\Sigma)^2}{2} - \mathbf{n}^- \frac{\mathbf{S}^{-, \text{irr}}}{\rho^-} \mathbf{n}^- \right) \right)\end{aligned}$$

Interfacial Entropy Production - Literature

Kovac, J.: Non-equilibrium thermodynamics of interfacial systems, *Physica A* **86**, 1-24 (1977).

Bedeaux, D.: Nonequilibrium thermodynamics and statistical physics of surfaces, pp. 47-109 in *Advance in Chemical Physics* **64** (I. Prigogine, S.A. Rice, eds), Jon Wiley & Sons 1986.

Slattery, J.C.: *Interfacial Transport Phenomena*. Springer, New York 1990.

Sagis, L.M.C.: Dynamic behavior of interfaces: Modeling with nonequilibrium thermodynamics, *Adv. Colloid & Interf. Sci.* **206**, 328-343 (2014).

Dreyer, W., Gohlke, C., Müller, R.: Bulk-surface electrothermodynamics and applications to electrochemistry, *Entropy* **20**, 939,1-44 (2018).

Bothe, D.: Sharp-interface continuum thermodynamics of multicomponent fluid systems with interfacial mass*, *Int. J. Eng. Sci.* **179**, 103731 (2022).

*based on conference proceedings: D.B. IBW7 (2015), D.B. RIMS (2016)

Interfacial Entropy Production

Reduced interfacial entropy production (zero interfacial viscosities):

$$\begin{aligned}
 \zeta^\Sigma = & \mathbf{q}^\Sigma \cdot \nabla_\Sigma \frac{1}{T^\Sigma} - \sum_{i=1}^N \mathbf{j}_i^\Sigma \cdot \left(\nabla_\Sigma \frac{\mu_i^\Sigma}{T^\Sigma} - \frac{\mathbf{b}_i^\Sigma}{T^\Sigma} \right) - \frac{1}{T} \sum_{a=1}^{N_R^\Sigma} R_a^\Sigma \mathcal{A}_a^\Sigma \\
 & - \frac{1}{T^\Sigma} (\mathbf{v}^+ - \mathbf{v}^\Sigma)_\parallel \cdot (\mathbf{S}^+ \mathbf{n}^+)_\parallel - \frac{1}{T^\Sigma} (\mathbf{v}^- - \mathbf{v}^\Sigma)_\parallel \cdot (\mathbf{S}^- \mathbf{n}^-)_\parallel \\
 & + \left(\frac{1}{T^\Sigma} - \frac{1}{T^+} \right) \left(\dot{m}^{+, \Sigma} \left(e^+ + \frac{p^+}{\rho^+} \right) + \mathbf{q}^+ \cdot \mathbf{n}^+ \right) \\
 & + \left(\frac{1}{T^\Sigma} - \frac{1}{T^-} \right) \left(\dot{m}^{-, \Sigma} \left(e^- + \frac{p^-}{\rho^-} \right) + \mathbf{q}^- \cdot \mathbf{n}^- \right) \\
 & + \sum_{i=1}^N \dot{m}_i^{+, \Sigma} \left(\frac{\mu_i^+}{T^+} - \frac{\mu_i^\Sigma}{T^\Sigma} + \frac{1}{T^\Sigma} \left(\frac{(\mathbf{v}^+ - \mathbf{v}^\Sigma)^2}{2} - \mathbf{n}^+ \frac{\mathbf{S}^{+, \text{irr}}}{\rho^+} \mathbf{n}^+ \right) \right) \\
 & + \sum_{i=1}^N \dot{m}_i^{-, \Sigma} \left(\frac{\mu_i^-}{T^-} - \frac{\mu_i^\Sigma}{T^\Sigma} + \frac{1}{T^\Sigma} \left(\frac{(\mathbf{v}^- - \mathbf{v}^\Sigma)^2}{2} - \mathbf{n}^- \frac{\mathbf{S}^{-, \text{irr}}}{\rho^-} \mathbf{n}^- \right) \right)
 \end{aligned}$$

Class I

are these terms grouped correctly?

Interfacial Entropy Production – Class II

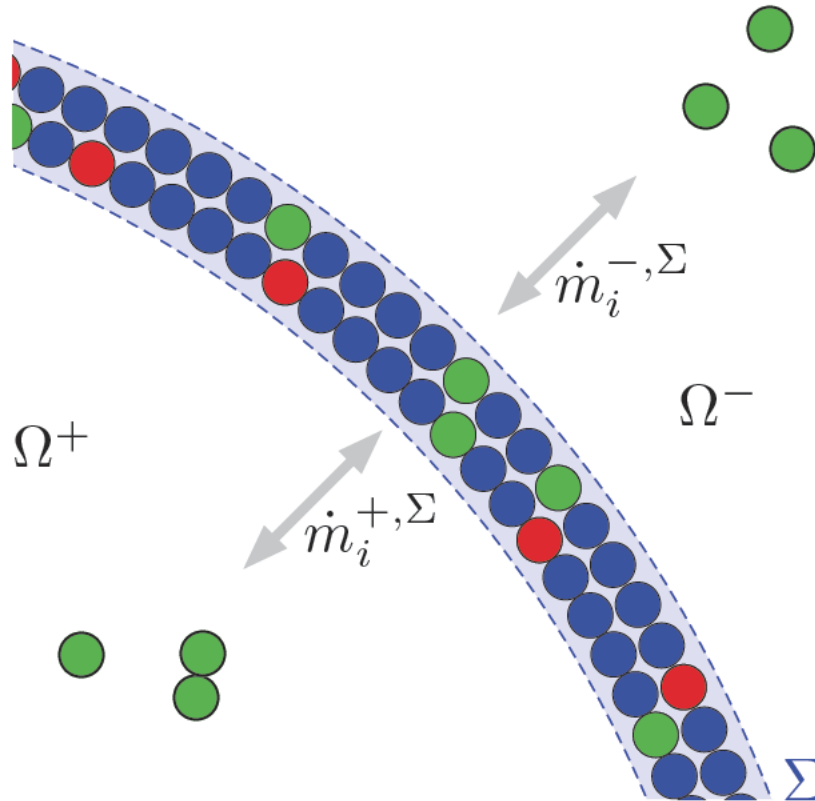
$$\begin{aligned}
 \zeta^\Sigma = & \frac{1}{T^\Sigma} \sum_{i=1}^N \nabla_\Sigma \mathbf{S}_i^{\Sigma, \circ} : \mathbf{D}_i^{\Sigma, \circ} + \frac{1}{T^\Sigma} \sum_{i=1}^N \pi_i^\Sigma \nabla_\Sigma \cdot \mathbf{v}_i^\Sigma + \mathbf{q}^\Sigma \cdot \nabla_\Sigma \frac{1}{T^\Sigma} - \frac{1}{T^\Sigma} \sum_a R_a^\Sigma \mathcal{A}_a^\Sigma \\
 & - \sum_{i=1}^N \mathbf{u}_i^\Sigma \cdot \left(\rho_i^\Sigma \nabla_\Sigma \frac{\mu_i^\Sigma}{T^\Sigma} + \frac{1}{T^\Sigma} (\mathbf{f}_i^\Sigma - r_i^\Sigma (\mathbf{v}_i^\Sigma - \frac{\mathbf{u}_i^\Sigma}{2}) - \nabla_\Sigma p_i^\Sigma) \right) \quad \text{Class II} \\
 & - \frac{1}{T^\Sigma} \sum_{i=1}^N (\mathbf{v}_i^+ - \mathbf{v}_i^\Sigma)_\parallel \cdot (\mathbf{S}_i^{+, \text{irr}} \mathbf{n}^+)_\parallel - \frac{1}{T^\Sigma} \sum_{i=1}^N (\mathbf{v}_i^- - \mathbf{v}_i^\Sigma)_\parallel \cdot (\mathbf{S}_i^{-, \text{irr}} \mathbf{n}^-)_\parallel \\
 & - \left(\frac{1}{T^+} - \frac{1}{T^\Sigma} \right) (\dot{m}^{+, \Sigma} h^+ + \mathbf{q}^+ \cdot \mathbf{n}^+) - \left(\frac{1}{T^-} - \frac{1}{T^\Sigma} \right) (\dot{m}^{-, \Sigma} h^- + \mathbf{q}^- \cdot \mathbf{n}^-) \\
 & + \sum_{i=1}^N \dot{m}_i^{+, \Sigma} \left(\frac{\mu_i^+}{T^+} - \frac{\mu_i^\Sigma}{T^\Sigma} + \frac{1}{T^\Sigma} \left(\frac{(\mathbf{v}_i^+ - \mathbf{v}_i^\Sigma)^2}{2} - \mathbf{n}^+ \cdot \frac{\mathbf{S}_i^{+, \text{irr}}}{\rho_i^+} \cdot \mathbf{n}^+ \right) \right) \\
 & + \sum_{i=1}^N \dot{m}_i^{-, \Sigma} \left(\frac{\mu_i^-}{T^-} - \frac{\mu_i^\Sigma}{T^\Sigma} + \frac{1}{T^\Sigma} \left(\frac{(\mathbf{v}_i^- - \mathbf{v}_i^\Sigma)^2}{2} - \mathbf{n}^- \cdot \frac{\mathbf{S}_i^{-, \text{irr}}}{\rho_i^-} \cdot \mathbf{n}^- \right) \right), \quad \Rightarrow \text{correctly grouped !}
 \end{aligned}$$

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-

One-sided Mass Transfer Processes in Series

- mass transfer as a sequence of „ad- and desorption“ processes



- transfer species pass through the force field of the other (adsorbed) constituents

Mass Transfer Entropy Production

Reduced interfacial entropy production (zero interfacial viscosities):

$$\begin{aligned}\zeta^\Sigma = & \mathbf{q}^\Sigma \cdot \nabla_\Sigma \frac{1}{T^\Sigma} - \sum_{i=1}^N \mathbf{j}_i^\Sigma \cdot \left(\nabla_\Sigma \frac{\mu_i^\Sigma}{T^\Sigma} - \frac{\mathbf{b}_i^\Sigma}{T^\Sigma} \right) - \frac{1}{T} \sum_{a=1}^{N_R^\Sigma} R_a^\Sigma \mathcal{A}_a^\Sigma \\ & - \frac{1}{T^\Sigma} (\mathbf{v}^+ - \mathbf{v}^\Sigma)_\parallel \cdot (\mathbf{S}^+ \mathbf{n}^+)_\parallel - \frac{1}{T^\Sigma} (\mathbf{v}^- - \mathbf{v}^\Sigma)_\parallel \cdot (\mathbf{S}^- \mathbf{n}^-)_\parallel \\ & + \left(\frac{1}{T^\Sigma} - \frac{1}{T^+} \right) \left(\dot{m}^{+, \Sigma} \left(e^+ + \frac{p^+}{\rho^+} \right) + \mathbf{q}^+ \cdot \mathbf{n}^+ \right) \\ & + \left(\frac{1}{T^\Sigma} - \frac{1}{T^-} \right) \left(\dot{m}^{-, \Sigma} \left(e^- + \frac{p^-}{\rho^-} \right) + \mathbf{q}^- \cdot \mathbf{n}^- \right) \\ & + \sum_{i=1}^N \dot{m}_i^{+, \Sigma} \left(\frac{\mu_i^+}{T^+} - \frac{\mu_i^\Sigma}{T^\Sigma} + \frac{1}{T^\Sigma} \left(\frac{(\mathbf{v}^+ - \mathbf{v}^\Sigma)^2}{2} - \mathbf{n}^+ \frac{\mathbf{S}^{+, \text{irr}}}{\rho^+} \mathbf{n}^+ \right) \right) \\ & + \sum_{i=1}^N \dot{m}_i^{-, \Sigma} \left(\frac{\mu_i^-}{T^-} - \frac{\mu_i^\Sigma}{T^\Sigma} + \frac{1}{T^\Sigma} \left(\frac{(\mathbf{v}^- - \mathbf{v}^\Sigma)^2}{2} - \mathbf{n}^- \frac{\mathbf{S}^{-, \text{irr}}}{\rho^-} \mathbf{n}^- \right) \right)\end{aligned}$$

Mass Transfer Closure

Interfacial mass transfer: decompose into ad- and desorption

$$\dot{m}_i^{+, \Sigma} = s_i^{\text{ad}, +} - s_i^{\text{de}, +}, \quad \dot{m}_i^{-, \Sigma} = s_i^{\text{ad}, -} - s_i^{\text{de}, -}$$

Mass transfer entropy production (simplified):

$$\zeta_{\text{TRANS}}^{\pm} = \frac{1}{T} \sum_{i=1}^N (s_i^{\text{ad}, \pm} - s_i^{\text{de}, \pm})(\mu_i^{\pm} - \mu_i^{\Sigma})$$

Closure of sorption kinetics: analogous to chemical reactions!

$$\ln \frac{s_i^{\text{ad}, \pm}}{s_i^{\text{de}, \pm}} = \frac{a_i^{\pm}}{RT} (\mu_i^{\pm} - \mu_i^{\Sigma}) \quad \text{with} \quad a_i^{\pm} \geq 0$$

Desorption modeled explicitly, adsorption follows: ($a_i^{\pm}$ set to 1)

$$s_i^{\text{de}, \pm} = k_i^{\text{de}, \pm} x_i^{\Sigma}, \quad s_i^{\text{ad}, \pm} = k_i^{\text{de}, \pm} x_i^{\Sigma} \exp\left(\frac{\mu_i^{\pm} - \mu_i^{\Sigma}}{RT}\right)$$

Mass Transfer Closure

Simple mixture assumption:

$$\begin{aligned}\mu_i^\pm(T, p, x_1, \dots, x_{N-1}) &= g_i^\pm(T, p) + RT \ln x_i^\pm \\ \mu_i^\Sigma(T, p, x_1^\Sigma, \dots, x_{N-1}^\Sigma) &= g_i^\Sigma(T, p^\Sigma) + RT \ln x_i^\Sigma\end{aligned}$$

Resulting (ad-)sorption kinetics:

$$s_i^{\text{de}, \pm} = k_i^{\text{de}, \pm} x_i^\Sigma, \quad s_i^{\text{ad}, \pm} = k_i^{\text{de}, \pm} \exp\left(\frac{g_i^\pm - g_i^\Sigma}{RT}\right) x_i^\pm =: k_i^{\text{ad}, \pm} x_i^\pm$$

approximation: $[\dot{m}_i] = 0 \quad \Leftrightarrow \quad \dot{m}_i^{+, \Sigma} + \dot{m}_i^{-, \Sigma} = 0$

Resulting mass transfer rates:

$$\dot{m}_i = \underbrace{\frac{k_i^{\text{de}, +} k_i^{\text{de}, -}}{k_i^{\text{de}, +} + k_i^{\text{de}, -}}}_{k_{\text{trans}}} \underbrace{\exp\left(-\frac{g_i^\Sigma}{RT}\right)}_{\text{influence of surface tension}} \left(\exp\left(\frac{g_i^+}{RT}\right) x_i^+ - \exp\left(\frac{g_i^-}{RT}\right) x_i^- \right)$$

Mass Transfer Hindrance Model

- combination of several* interface free energy models

$$\mu_k^{\Sigma, \text{mol}} = a_k(T) + \alpha_k RT \ln\left(1 + \frac{\pi^\Sigma}{K^\Sigma}\right) + \beta_k \pi^\Sigma + \sum_{i=1}^N \cancel{f_i(c_j^\Sigma)} c_i^\Sigma + RT \ln \chi_k^\Sigma$$

χ_k^Σ is one of c_k^Σ / c_S^Σ , $c_k^\Sigma / c_k^{\Sigma, \infty}$ or $x_k^\Sigma = c_k^\Sigma / c^\Sigma$

final result:

$$\dot{m}_i^{\text{contam}} = \alpha(\sigma) \dot{m}_i^{\text{clean}}$$

matches experimental findings by A. Tomiyama

up to non-linear corrections for non-ideal interface mixtures

* lattice-based, area-based, gas-type, explicit solvent contribution, excluded area

Mass Transfer Hindrance Model

- combination of several* interface free energy models

$$\mu_k^{\Sigma, \text{mol}} = a_k(T) + \alpha_k RT \ln\left(1 + \frac{\pi^{\Sigma}}{K^{\Sigma}}\right) - \beta_k \pi^{\Sigma} + \sum_{i=1}^N f_i(c_j^{\Sigma}) c_i^{\Sigma} + RT \ln \chi_k^{\Sigma}$$

main effect of
surface pressure = - σ

final result:

Langmuir
energy
barrier

$$\dot{m}_i^{\text{contam}} = \alpha(\sigma) \dot{m}_i^{\text{clean}}$$

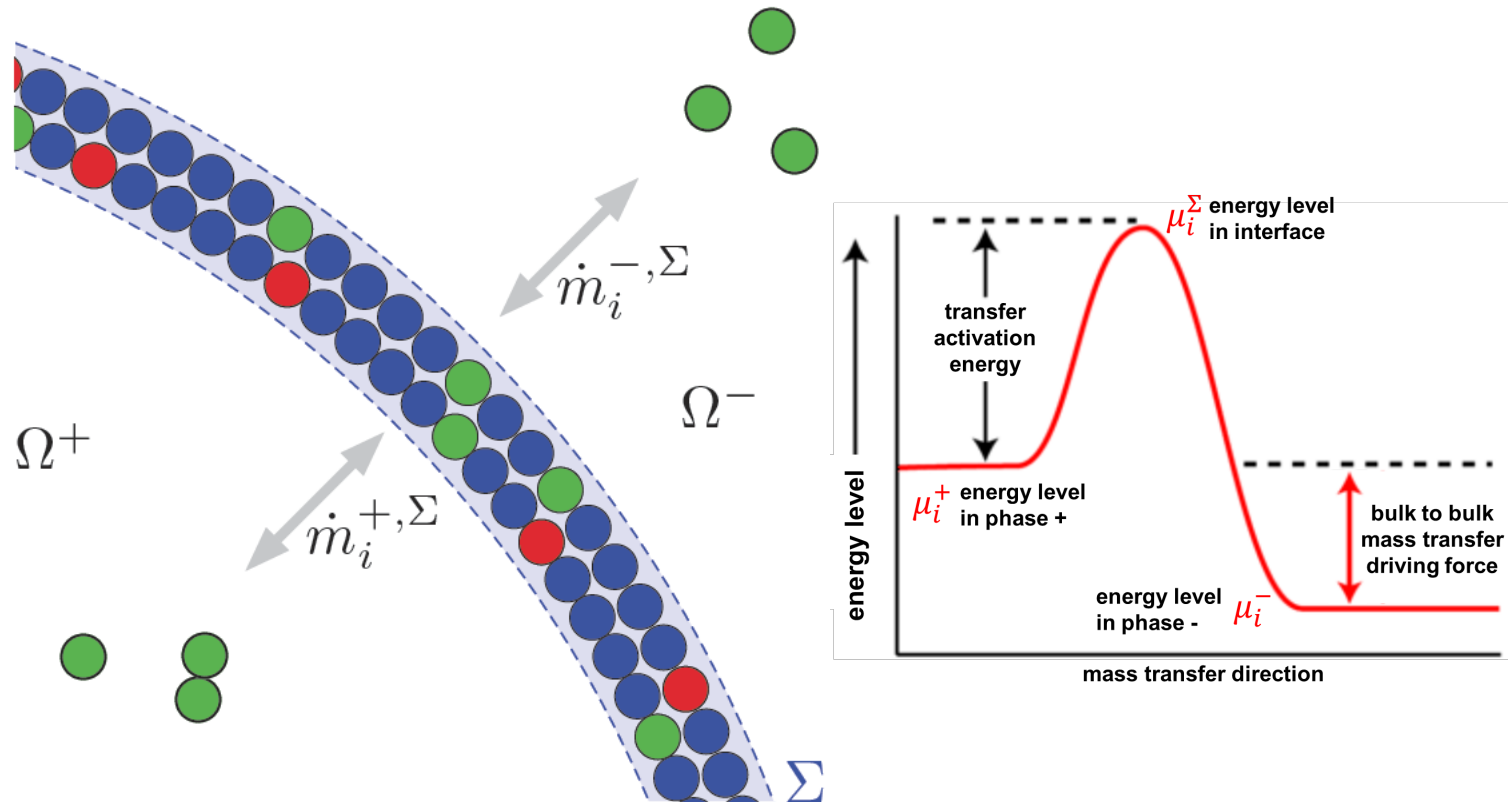
$$\alpha(\sigma) = \exp\left(-\frac{\sigma_0 - \sigma}{c_{\Sigma}^{\infty} RT}\right) \quad \text{Boltzmann factor}$$

Since $\mu_i^{\Sigma} = \mu_i^{\Sigma}(T^{\Sigma}, \sigma, x_k^{\Sigma})$ depends on surface tension:

surfactant $k \rightarrow$ changes $\sigma \rightarrow$ changes all $\mu_i^{\Sigma} \rightarrow$ changes $\dot{m}_i^{\pm}$

One-sided Mass Transfer Processes in Series

- mass transfer as a sequence of „ad- and desorption“ processes



- transfer species pass through the force field of the other (adsorbed) constituents

Thank You for Your
kind attention !