

Implicitly constituted materials with fading memory

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Incompressible simple fluid

Truesdell and Noll (1965):

$$\mathbb{T} = -p\mathbb{1} + \mathfrak{F}_{s=0}^{+\infty}(\mathbb{C}_t(t-s))$$

Differential type models

General form, Rivlin and Ericksen (1955):

$$\mathbb{T} = -p\mathbb{I} + \mathfrak{f}(\mathbb{A}_1, \mathbb{A}_2, \mathbb{A}_3, \dots)$$

where

$$\mathbb{A}_1 = 2\mathbb{D}$$

$$\mathbb{A}_n = \frac{d\mathbb{A}_{n-1}}{dt} + \mathbb{A}_{n-1}\mathbb{L} + \mathbb{L}^\top\mathbb{A}_{n-1}$$

Coleman and Noll (1960): These models can be understood as successive approximations of the history functional

$$\mathbb{T} = -p\mathbb{I} + \mathfrak{F}_{s=0}^{+\infty}(\mathbb{C}_t(t-s))$$

with “fading memory”.

Pressure dependent viscosity

$$\mathbb{T} = -p\mathbb{I} + 2\mu(p)\mathbb{D}$$

Stress dependent viscosity

$$\mathbb{T} = -p\mathbb{I} + \mu(\mathbb{T})\mathbb{D}$$

Seely (1964):

$$\mu(\mathbb{T}) = \mu_{\infty} + (\mu_0 - \mu_{\infty}) e^{-\frac{|\mathbb{T}_{\delta}|}{\tau_0}}$$

Blatter (1995):

$$\mu(\mathbb{T}) = \frac{A}{\left(|\mathbb{T}_{\delta}|^2 + \tau_0^2\right)^{\frac{n-1}{2}}}$$

Matsuhisa and Bird (1965):

$$\mu(\mathbb{T}) = \frac{\mu_0}{1 + \alpha |\mathbb{T}_{\delta}|^{n-1}}$$

Implicit constitutive relation

Rajagopal (2003, 2006); Rajagopal and Srinivasa (2008):

$$f(\mathbb{T}, \mathbb{D}) = 0$$

Rate type models

$$\mathbb{T} = -\pi \mathbb{I} + \mathbb{S}$$

Oldroyd (1958):

$$\begin{aligned} \mathbb{S} + \lambda_1 \overset{\nabla}{\mathbb{S}} + \frac{\lambda_3}{2} (\mathbb{D}\mathbb{S} + \mathbb{S}\mathbb{D}) + \frac{\lambda_5}{2} (\text{Tr } \mathbb{S}) \mathbb{D} + \frac{\lambda_6}{2} (\mathbb{S} : \mathbb{D}) \mathbb{I} \\ = -\mu \left(\mathbb{D} + \lambda_2 \overset{\nabla}{\mathbb{D}} + \lambda_4 \mathbb{D}^2 + \frac{\lambda_7}{2} (\mathbb{D} : \mathbb{D}) \mathbb{I} \right) \end{aligned}$$

Phan Thien (1978):

$$\begin{aligned} \mathbb{Y}\mathbb{S} + \lambda \overset{\nabla}{\mathbb{S}} + \frac{\lambda \xi}{2} (\mathbb{D}\mathbb{S} + \mathbb{S}\mathbb{D}) = -\mu \mathbb{D} \\ \mathbb{Y} = e^{-\varepsilon \frac{\lambda}{\mu} \text{Tr } \mathbb{S}} \end{aligned}$$

Notation:

$$\overset{\nabla}{\mathbb{b}^b} =_{\text{def}} \frac{d\mathbb{b}^b}{dt} - [\nabla \mathbf{v}] \mathbb{b}^b - \mathbb{b}^b [\nabla \mathbf{v}]^T$$

Materials with fading memory

Implicit algebraic relation:

$$f(\mathbb{T}, \mathbb{D}) = 0$$

Implicit relation between the histories:

$$\mathfrak{H}_{s=0}^{+\infty} (\mathbb{T}(t-s), \mathbb{C}_t(t-s)) = 0$$

Questions:

- ▶ Are rate type and differential type models (and other known models) special instances or **approximations** of the material with fading memory?
- ▶ Is something like the celebrated **retardation theorem** by Coleman and Noll (1960) available for implicit type materials with fading memory?

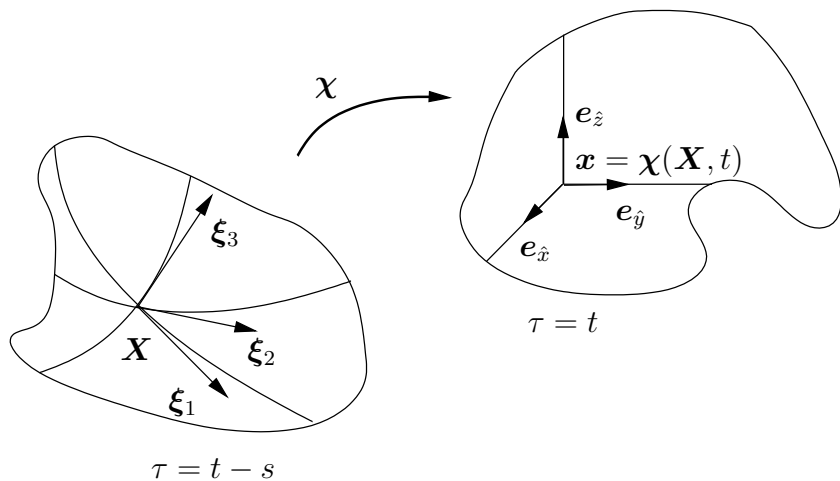
Independent variables

[...] the properties of a material element may depend upon the previous rheological states through which that element has passed, but *not in any way on the states of neighbouring elements and not on the motion of the element as a whole in the space.*

[...] only those tensor quantities need to be considered which have a significance for the material element independent of its motion as a whole in space.

Oldroyd (1950)

Convected coordinate system



Plan

- ▶ Formulate the constitutive relation in the *convected coordinate system*, Hencky (1925), Oldroyd (1950).
- ▶ Formulate the constitutive relation in an implicit form, Rajagopal (2003).
- ▶ Expand the functional using an analogue of the *retardation theorem*, Coleman and Noll (1960).
- ▶ Use *representation theorems* for isotropic linear and bilinear functions, Truesdell and Noll (1965).
- ▶ Transform the constitutive relation to a *fixed-in-space coordinate system*, Oldroyd (1950).

Constitutive assumptions

General relation:

$$\mathfrak{H}_{s=0}^{+\infty} (\mathfrak{D}(\xi, t - s), \mathfrak{T}(\xi, t - s)) = \mathbb{0},$$

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Special form of the general constitutive relation:

$$\mathfrak{T} = -\pi \mathbb{1} + \mathfrak{O}$$

$$\mathbb{0} = \mathfrak{G}_{s=0}^{+\infty} (\mathfrak{D}(\xi, t - s) - \mathbb{1}, \mathfrak{O}(\xi, t - s))$$

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Constitutive assumption:

$$\|\mathfrak{O}(\xi, t-s)\|_S = O(\|\mathfrak{D}(\xi, t-s) - \mathbb{1}\|_T) \text{ as } \|\mathfrak{D}(\xi, t-s) - \mathbb{1}\|_T \rightarrow 0+,$$

Norm

[...] the deformations that occurred in the distant past should have less influence in determining the present stress than those that occurred in the recent past.

Truesdell and Noll (1965)

$$\|\vartheta\|_{L_h^2} =_{\text{def}} \left(\int_{s=0}^{+\infty} |\vartheta(s)|^2 h(s) \, ds \right)^{\frac{1}{2}}$$
$$\left\| [\vartheta, \sigma]^\top \right\|_{L_h^2 \times L_h^2} =_{\text{def}} \left(\|\vartheta\|_{L_h^2}^2 + \|\sigma\|_{L_h^2}^2 \right)^{\frac{1}{2}}$$

Taylor series for the functional

$$\pi = -\pi \mathbb{1} + \sigma$$

$$0 = \mathfrak{G}_{s=0}^{+\infty} (\vartheta(\xi, t-s) - \mathbb{1}, \sigma(\xi, t-s))$$

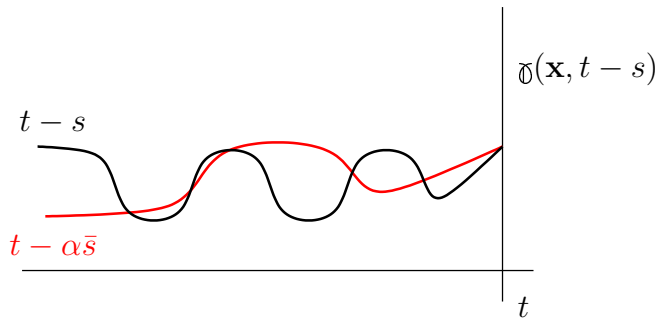
$$\mathfrak{G}_{s=0}^{+\infty} \left(\begin{bmatrix} \vartheta(\xi, t-s) - \mathbb{1} \\ \sigma(\xi, t-s) \end{bmatrix} \right) = A+B+C+o \left(\left\| \begin{bmatrix} \vartheta(\xi, t-s) - \mathbb{1} \\ \sigma(\xi, t-s) \end{bmatrix} \right\|_{L_h^2 \times L_h^2}^2 \right)$$

$$A = \mathfrak{G}_{s=0}^{+\infty} \left(\begin{bmatrix} 0 \\ 0 \end{bmatrix} \right)$$

$$B = \delta \mathfrak{G}_{s=0}^{+\infty} \left(\begin{bmatrix} 0 \\ 0 \end{bmatrix} \right) \begin{bmatrix} \vartheta(\xi, t-s) - \mathbb{1} \\ \sigma(\xi, t-s) \end{bmatrix}$$

$$C = [\vartheta(\xi, t-s) - \mathbb{1}, \sigma(\xi, t-s)]^\top \delta^2 \mathfrak{G}_{s=0}^{+\infty} \left(\begin{bmatrix} 0 \\ 0 \end{bmatrix} \right) \begin{bmatrix} \vartheta(\xi, t-s) - \mathbb{1} \\ \sigma(\xi, t-s) \end{bmatrix}$$

Slow history



Taylor series for the metric tensor–stress tensor history

Informal expansion:

$$\begin{aligned} & \begin{bmatrix} \bar{\vartheta}(\xi, t - \alpha \bar{s}) - \mathbb{1} \\ \bar{\sigma}(\xi, t - \alpha \bar{s}) \end{bmatrix} \\ &= \begin{bmatrix} \bar{\vartheta}(\xi, t) - \mathbb{1} \\ \bar{\sigma}(\xi, t) \end{bmatrix} - \alpha \bar{s} \begin{bmatrix} \frac{d\bar{\vartheta}}{dt}(\xi, t) \\ \frac{d\bar{\sigma}}{dt}(\xi, t) \end{bmatrix} + \frac{1}{2} \alpha^2 \bar{s}^2 \begin{bmatrix} \frac{d^2\bar{\vartheta}}{dt^2}(\xi, t) \\ \frac{d^2\bar{\sigma}}{dt^2}(\xi, t) \end{bmatrix} + o(\alpha^2) \\ &= \underbrace{\begin{bmatrix} 0 \\ \bar{\sigma}(\xi, t) \end{bmatrix}}_{\mathbf{g}^{(0)}} - \alpha \underbrace{\begin{bmatrix} \frac{d\bar{\vartheta}}{dt}(\xi, t) \\ \frac{d\bar{\sigma}}{dt}(\xi, t) \end{bmatrix}}_{\mathbf{g}^{(1)}} \bar{s} + \frac{1}{2} \alpha^2 \underbrace{\begin{bmatrix} \frac{d^2\bar{\vartheta}}{dt^2}(\xi, t) \\ \frac{d^2\bar{\sigma}}{dt^2}(\xi, t) \end{bmatrix}}_{\mathbf{g}^{(2)}} \bar{s}^2 + o(\alpha^2) \end{aligned}$$

Rigorous result:

$$\lim_{\alpha \rightarrow 0^+} \frac{1}{\alpha^2} \left\| \begin{bmatrix} \bar{\vartheta}(\xi, t - \alpha \bar{s}) - \mathbb{1} \\ \bar{\sigma}(\xi, t - \alpha \bar{s}) \end{bmatrix} - \left(\mathbf{g}^{(0)} - \alpha \mathbf{g}^{(1)} \bar{s} + \frac{1}{2} \alpha^2 \mathbf{g}^{(2)} \bar{s}^2 \right) \right\|_{L_h^2 \times L_h^2} = 0$$

Approximation for slow histories

Use constitutive assumption

$$\|\sigma(\xi, t - s)\|_S = O(\|\bar{\sigma}(\xi, t - s) - \bar{\sigma}\|_\Gamma) \text{ as } \|\bar{\sigma}(\xi, t - s) - \bar{\sigma}\|_\Gamma \rightarrow 0+,$$

and substitute to the Taylor formula for the functional.

First order:

$$\mathfrak{G}_{s=0}^{+\infty} \left(\begin{bmatrix} \bar{\sigma}(\xi, t - \alpha \bar{s}) - \bar{\sigma} \\ \sigma(\xi, t - \alpha \bar{s}) - \sigma(\xi, t) \end{bmatrix} \right) = f_0 \left(\begin{smallmatrix} \textcircled{0} \\ \mathbf{g} \end{smallmatrix} \right) + f_1 \left(\begin{smallmatrix} \textcircled{1} \\ \mathbf{g} \end{smallmatrix} \right) + o(\alpha)$$

Second order:

$$\begin{aligned} \mathfrak{G}_{s=0}^{+\infty} \left(\begin{bmatrix} \bar{\sigma}(\xi, t - \alpha \bar{s}) - \bar{\sigma} \\ \sigma(\xi, t - \alpha \bar{s}) - \sigma(\xi, t) \end{bmatrix} \right) &= f_0 \left(\begin{smallmatrix} \textcircled{0} \\ \mathbf{g} \end{smallmatrix} \right) + f_1 \left(\begin{smallmatrix} \textcircled{1} \\ \mathbf{g} \end{smallmatrix} \right) + f_2 \left(\begin{smallmatrix} \textcircled{2} \\ \mathbf{g} \end{smallmatrix} \right) \\ &+ g_{00} \left(\begin{smallmatrix} \textcircled{0} & \textcircled{0} \\ \mathbf{g} & \mathbf{g} \end{smallmatrix} \right) + g_{10} \left(\begin{smallmatrix} \textcircled{1} & \textcircled{0} \\ \mathbf{g} & \mathbf{g} \end{smallmatrix} \right) + g_{01} \left(\begin{smallmatrix} \textcircled{0} & \textcircled{1} \\ \mathbf{g} & \mathbf{g} \end{smallmatrix} \right) + g_{11} \left(\begin{smallmatrix} \textcircled{1} & \textcircled{1} \\ \mathbf{g} & \mathbf{g} \end{smallmatrix} \right) + o(\alpha^2) \end{aligned}$$

Linear and bilinear tensor functions

Representation for isotropic linear functions:

$$\mathfrak{h}(\mathbb{A}) = a_1 (\text{Tr } \mathbb{A}) \mathbb{I} + a_2 \mathbb{A}$$

Representation for isotropic bilinear functions:

$$\begin{aligned} \mathfrak{h}(\mathbb{A}, \mathbb{B}) = & (c_1 \text{Tr } \mathbb{A} \text{Tr } \mathbb{B} + c_2 \text{Tr}(\mathbb{A}\mathbb{B})) \mathbb{I} \\ & + c_3 (\text{Tr } \mathbb{A}) \mathbb{B} + c_4 (\text{Tr } \mathbb{B}) \mathbb{A} + c_5 (\mathbb{A}\mathbb{B} + \mathbb{B}\mathbb{A}) \end{aligned}$$

Time derivatives with respect to fixed-in-space coordinate system

Oldroyd (1950):

$$\frac{\partial b^k_i}{\partial t} =_{\text{def}} \left(\frac{\partial b^k_i}{\partial t} + v^m b^k_{i|m} \right) - v^k|_m b^m_i + b^k_m v^m|_i,$$

$$\frac{\partial b_{ki}}{\partial t} =_{\text{def}} \left(\frac{\partial b_{ki}}{\partial t} + v^m b_{ki|m} \right) + v^m|_k b_{mi} + b_{km} v^m|_i,$$

$$\frac{\partial b^{ki}}{\partial t} =_{\text{def}} \left(\frac{\partial b^{ki}}{\partial t} + v^m b^{ki|m} \right) - v^k|_m b^{mi} - b^{km} v^i|_m,$$

Time derivatives with respect to fixed-in-space coordinate system

$$\frac{\partial \mathbf{b}}{\partial t} = \frac{d\mathbf{b}}{dt} - [\nabla \mathbf{v}] \mathbf{b} + \mathbf{b} [\nabla \mathbf{v}],$$

$$\frac{\partial \mathbf{b}^\#}{\partial t} = \frac{d\mathbf{b}^\#}{dt} + [\nabla \mathbf{v}]^\top \mathbf{b}^\# + \mathbf{b}^\# [\nabla \mathbf{v}],$$

$$\frac{\partial \mathbf{b}^b}{\partial t} = \frac{d\mathbf{b}^b}{dt} - [\nabla \mathbf{v}] \mathbf{b}^b - \mathbf{b}^b [\nabla \mathbf{v}]^\top,$$

Identification of the derivatives

$$\sigma(\xi, t) \mapsto \mathbb{S}(\mathbf{x}, t)$$

$$\frac{d\sigma}{dt}(\xi, t) \mapsto \mathbb{D}(\mathbf{x}, t)$$

$$\frac{d^2\sigma}{dt^2}(\xi, t) \mapsto \overset{\nabla}{\mathbb{D}}(\mathbf{x}, t)$$

$$\frac{d\sigma}{dt}(\xi, t) \mapsto \overset{\nabla}{\mathbb{S}}(\mathbf{x}, t)$$

$$\frac{d^2\sigma}{dt^2}(\xi, t) \mapsto \overset{\nabla\nabla}{\mathbb{S}}(\mathbf{x}, t)$$

Approximation formulae

First order:

$$\mathfrak{G}_{s=0}^{+\infty} \left(\begin{bmatrix} \mathfrak{D}(\xi, t - \alpha \bar{s}) - \mathbb{1} \\ \sigma(\xi, t - \alpha \bar{s}) - \sigma(\xi, t) \end{bmatrix} \right) \\ \mapsto b_0 (\text{Tr } \mathbb{S}) \mathbb{1} + b_1 \mathbb{S} + 2b_3 \mathbb{D} + b_4 \left(\text{Tr } \overset{\vee}{\mathbb{S}} \right) \mathbb{1} + b_5 \overset{\vee}{\mathbb{S}} + o(\alpha)$$

Approximation formulae

Second order:

$$\begin{aligned}
 & \mathfrak{G}_{s=0}^{+\infty} \left(\left[\begin{array}{c} \mathfrak{D}(\xi, t - \alpha \bar{s}) - \mathbb{1} \\ \sigma(\xi, t - \alpha \bar{s}) - \sigma(\xi, t) \end{array} \right] \right) \\
 & \mapsto \left[b_0 (\text{Tr } \mathbb{S}) + b_4 \left(\text{Tr } \overset{\vee}{\mathbb{S}} \right) + b_6 \left(\text{Tr } \overset{\vee\vee}{\mathbb{S}} \right) + (b_{15} - 2b_8) \text{Tr} (\mathbb{D})^2 \right. \\
 & + \left(b_{10} (\text{Tr } \mathbb{S})^2 + b_{11} \text{Tr} (\mathbb{S})^2 \right) + \left(b_{18} \left(\text{Tr } \overset{\vee}{\mathbb{S}} \right)^2 + b_{19} \text{Tr} \left(\overset{\vee}{\mathbb{S}} \right)^2 \right) \\
 & + b_{23} \text{Tr} \left(\mathbb{D} \overset{\vee}{\mathbb{S}} \right) + \left(b_{27} \text{Tr } \mathbb{S} \text{Tr } \overset{\vee}{\mathbb{S}} + b_{28} \text{Tr} \left(\mathbb{S} \overset{\vee}{\mathbb{S}} \right) \right) + b_{33} \text{Tr} (\mathbb{S} \mathbb{D}) \Big] \mathbb{1} \\
 & + \left[b_1 + b_{12} (\text{Tr } \mathbb{S}) + b_{30} \left(\text{Tr } \overset{\vee}{\mathbb{S}} \right) \right] \mathbb{S} + \left[b_3 + b_{25} \left(\text{Tr } \overset{\vee}{\mathbb{S}} \right) + b_{34} (\text{Tr } \mathbb{S}) \right] \mathbb{D} \\
 & \quad + b_{13} (\mathbb{S})^2 + b_{17} (\mathbb{D})^2 + b_{36} (\mathbb{S} \mathbb{D} + \mathbb{D} \mathbb{S}) \\
 & \quad + b_9 \overset{\vee}{\mathbb{D}} + \left[b_5 + b_{20} \left(\text{Tr } \overset{\vee}{\mathbb{S}} \right) + b_{29} (\text{Tr } \mathbb{S}) \right] \overset{\vee}{\mathbb{S}} + b_{21} \left(\overset{\vee}{\mathbb{S}} \right)^2 \\
 & \quad + b_{26} \left(\mathbb{D} \overset{\vee}{\mathbb{S}} + \overset{\vee}{\mathbb{S}} \mathbb{D} \right) + b_{31} \left(\mathbb{S} \overset{\vee}{\mathbb{S}} + \overset{\vee}{\mathbb{S}} \mathbb{S} \right) + b_7 \overset{\vee\vee}{\mathbb{S}} + o(\alpha^2)
 \end{aligned}$$

Rate type models

$$\mathbb{T} = -\pi \mathbb{I} + \mathbb{S}$$

Oldroyd (1958):

$$\begin{aligned} \mathbb{S} + \lambda_1 \overset{\nabla}{\mathbb{S}} + \frac{\lambda_3}{2} (\mathbb{D}\mathbb{S} + \mathbb{S}\mathbb{D}) + \frac{\lambda_5}{2} (\text{Tr } \mathbb{S}) \mathbb{D} + \frac{\lambda_6}{2} (\mathbb{S} : \mathbb{D}) \mathbb{I} \\ = -\mu \left(\mathbb{D} + \lambda_2 \overset{\nabla}{\mathbb{D}} + \lambda_4 \mathbb{D}^2 + \frac{\lambda_7}{2} (\mathbb{D} : \mathbb{D}) \mathbb{I} \right) \end{aligned}$$

Phan Thien (1978):

$$\begin{aligned} \mathbb{Y}\mathbb{S} + \lambda \overset{\nabla}{\mathbb{S}} + \frac{\lambda \xi}{2} (\mathbb{D}\mathbb{S} + \mathbb{S}\mathbb{D}) = -\mu \mathbb{D} \\ \mathbb{Y} = e^{-\varepsilon \frac{\lambda}{\mu} \text{Tr } \mathbb{S}} \end{aligned}$$

Notation:

$$\overset{\nabla}{\mathbb{b}^b} =_{\text{def}} \frac{d\mathbb{b}^b}{dt} - [\nabla \mathbf{v}] \mathbb{b}^b - \mathbb{b}^b [\nabla \mathbf{v}]^T$$

Conclusion

- ▶ Implicit constitutive relations provide a general concept that can accommodate **both differential and rate type models**.
Generalization of the concept of simple fluid.
- ▶ Rate type models for viscoelastic materials can be seen as special instances of a general material with fading memory.
(After a **rigorous approximation procedure**.)
- ▶ This is only a proof of concept—there are better ways how to derive implicit type constitutive relations that are **consistent with the laws of thermodynamics**.

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