

# Mechanical oscillators described by a system of differential-algebraic equations

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# Problem

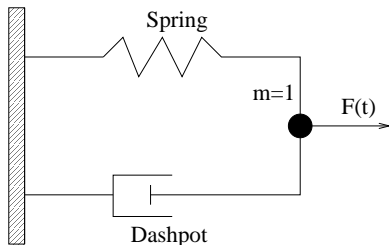
$$x'' + F_d + F_s = F(t)$$

$x$  ..... displacement

$F_d$  ..... dashpot force

$F_s$  ..... spring force

$F(t)$  ..... external force



$$x'' + F_d + F_s = F(t)$$

“common” approach:

$$F_s = f(x) \quad (\text{spring})$$

$$F_d = g(x') \quad (\text{dashpot})$$

$$x'' + g(x') + f(x) = F(t)$$

apply the standard  
ODE theory

# “Reversed” constitutive relations

IDEA:

what if we assume

$$x = f(F_s) \quad (\text{spring})$$

$$x' = g(F_d) \quad (\text{dashpot})$$

PHILOSOPHICALLY: kinematics ( $x$  and  $x'$ ) are a consequence, and hence a function of the forces ( $F_s$  and  $F_d$ ).

$$x'' + F_d + F_s = F(t)$$

$$x = f(F_s)$$

$$x' = g(F_d)$$

differential-algebraic  
system of equations

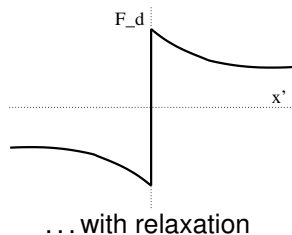
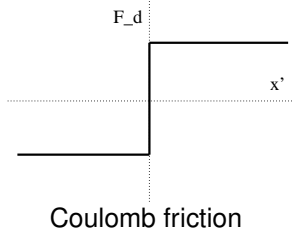
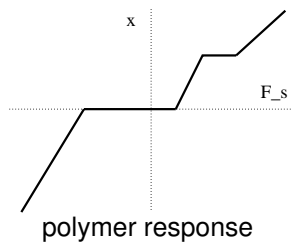
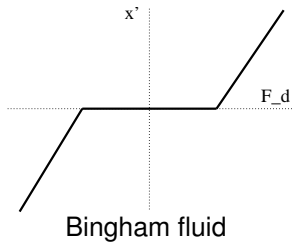
For some materials, it is even reasonable to assume:

$$f(x, F_s) = 0 \quad (\text{spring})$$

$$g(x', F_d) = 0 \quad (\text{dashpot})$$

That is to say, fully implicit constitutive relations.

# Examples



- ① oscillators with reversed (monotone) constitutive relations
- ② oscillator with (generalized) Coulomb friction
- ③ problem: uniqueness for 2nd order ODE's

# Oscillators with reversed constitutive relations

$$\begin{aligned}x'' + F_d + F_s &= F(t) \\x &= f(F_s) \\x' &= g(F_d)\end{aligned}$$

- $f, g$  continuous, **non-decreasing**
- $|f(u)|, |g(u)| \sim |u|$  for  $|u| \rightarrow \infty$
- $F(t) \in L^2(0, T)$



**THEOREM 1.** There is at least one global solution.

*Proof.*

① approximation: 
$$\begin{aligned}x &= f_k(F_s) & f_k &= f + k^{-1} \text{Id} \\x' &= g_k(F_d) & g_k &= g + k^{-1} \text{Id}\end{aligned}$$

②  $f_k, g_k$  invertible  $\rightsquigarrow$  
$$x'' + \underbrace{\{g_k\}_{-1}(x')}_{F_d} + \underbrace{\{f_k\}_{-1}(x)}_{F_s} = F(t)$$

③ coercivity of  $f, g \implies k$ -independent estimates

④ limit  $k \rightarrow \infty$  (use monotonicity of  $f, g$ ).

## ... uniqueness ... ?

$x_1, x_2 \dots$  solutions;

$F_d^i, F_s^i, i = 1, 2 \dots$  the corresponding forces.

$$(x_1 - x_2)'' + (F_d^1 - F_d^2) + (F_s^1 - F_s^2) = 0 \quad / \cdot (x^1 - x^2)'$$

$$\frac{1}{2} \frac{d}{dt} \{ (x_1 - x_2)' \}^2 + \underbrace{(F_d^1 - F_d^2)(x^1 - x^2)'}_{\geq 0} + \underbrace{(F_s^1 - F_s^2)(x^1 - x^2)'}_{???$$

assume in addition:

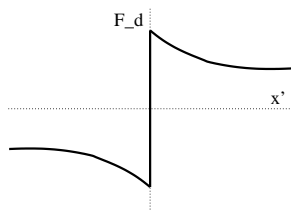
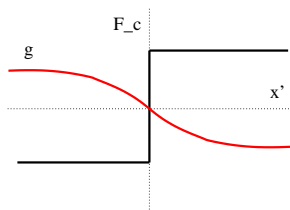
- structural properties of  $f$
- $F(t) \equiv F_0 \dots$  **autonomous case**

$\implies$  **THEOREM 2.** Global (forward) uniqueness

# Coulomb friction with relaxation

$$x'' + F_d + kx = F(t)$$
$$F_d = F_c + g(x')$$
$$(F_c, x') \in \mathcal{A}$$

$F_c$  ..... Coulomb-like friction force  
 $\mathcal{A}$  ..... monotone graph  
 $g(\cdot)$  ..... relaxation function



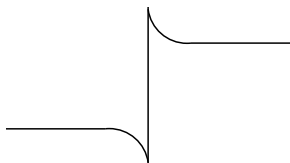
- $g$  continuous,  $|g(u)| \leq c(1 + |u|)$
- $\mathcal{A}$  maximal monotone, coercive

$\implies$  **THEOREM 1.** Global existence of solutions.

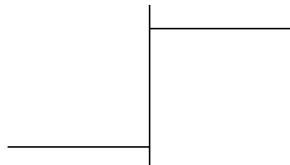
- moreover:  $g$  locally lipschitz

$\implies$  **THEOREM 2.** Global (forward) uniqueness.

examples of nonuniqueness:



(steep relaxation)



(non-monotone graph)

# Simplification: uniqueness for ODE

**motivation:**

$$x'' + F_d + F_s = F(t)$$

$\uparrow$              $\uparrow$   
 $x'$              $x$

① neglect  $F_s$  and  $x$        $\rightsquigarrow y' + f(y, t) = 0$       ( $y = x'$ )

② neglect  $F_d$  and  $x'$        $\rightsquigarrow x'' + f(x, t) = 0$

# Uniqueness for 1st order ODE ?

$$y' + f(y, t) = 0$$

- $f(\cdot, t)$  locally lipschitz: YES
- $f(\cdot, t)$  only Hölder: NO
- $f(\cdot, t)$  non-decreasing: YES (forward)

# Uniqueness for 2nd order ODE ?

$$x'' + f(x, t) = 0$$

- $f(\cdot, t)$  locally lipschitz: YES
- $f(\cdot, t)$  only Hölder: NO
- $f(\cdot, t)$  non-decreasing: **NO** in general
  - linear counterexample:  $x'' + Q(t)x = 0, \quad Q(t) \geq 0.$
  - autonomous problem:  $\implies$  **uniqueness**
  - “quasi-autonomous” case:  $x'' + h(x) = f(t)$  **????**

Thank you.