#### Incompressible piezoviscous fluids: first steps, a long ways to go.

M. Lanzendörfer

Mathematical Institute, Charles University in Prague

Institute of Computer Science, AS CR

M. Lanzendörfer (Charles University; ICS CAS)

Incompressible piezoviscous fluids

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#### Motivation,

- hydrodynamic lubrication
- Reynolds equation

#### incompressible piezoviscous fluids,

- viscosity depending on pressure and shear rate
- implicitely constituted models

#### challenges and results,

- inf-sup inequality and uniqueness of the pressure
- subclass of models: restricted growth of extra stress with pressure
- well-posed PDEs, convergence of FEM, pictures

#### and their limitations.

- subclass of models? limited range of pressures and shear rates
- relevance to hydrodynamic lubrication

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#### Motivation,

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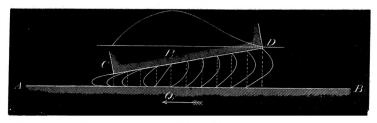
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IV. On the Theory of Lubrication and its Application to Mr. BEAUCHAMP TOWER'S Experiments, including an Experimental Determination of the Viscosity of Olive Oil.

By Professor Osborne Reynolds, LL.D., F.R.S.

Received December 29, 1885,-Read February 11, 1886.

Fig. 9.

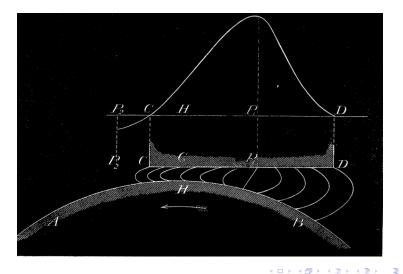


#### This is the explanation of continuous lubrication.

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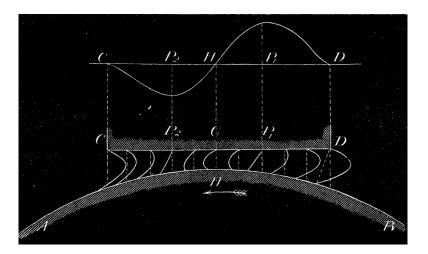
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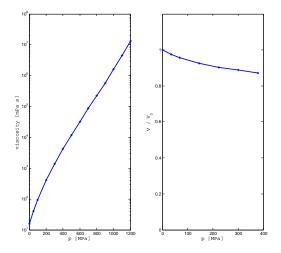
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#### incompressible piezoviscous fluids,

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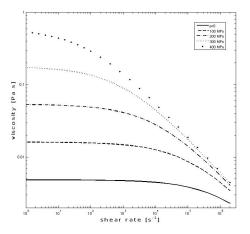
## Viscosity at large pressure and shear rate Viscosity and volume variation with pressure for *squalane* ("representing a low viscosity paraffinic mineral oil", S. Bair, *Tribology Letters*, 2006).



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### Viscosity at large pressure and shear rate

Viscosity for *SAE 10W/40 reference oil RL 88/1*, (partly) by Hutton, Jones, Bates, *SAE*, 1983.



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Mathematical formulation inside  $(0, T) \times \Omega$ :

$$\begin{aligned} \operatorname{div} \boldsymbol{v} &= 0\\ \partial_{\tau} \boldsymbol{v} + \operatorname{div}(\boldsymbol{v} \otimes \boldsymbol{v}) - \operatorname{div} \boldsymbol{S} &= -\nabla \pi + \boldsymbol{f},\\ \boldsymbol{S} &= 2 \nu(\pi, |\boldsymbol{D}(\boldsymbol{v})|^2) \boldsymbol{D}(\boldsymbol{v}) \end{aligned}$$

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- Cauchy stress tensor  $\boldsymbol{T} = -\pi \boldsymbol{I} + 2\nu(\pi, |\boldsymbol{D}|^2)\boldsymbol{D}$ , tr  $\boldsymbol{D} = 0$
- $\pi$  is the mean normal stress,  $\pi = -\frac{1}{3} \operatorname{tr} \boldsymbol{T}$ ,
- implicitely constituted model

$$m{T}-rac{1}{3}( ext{tr}~m{T})m{I}-2
u(-rac{1}{3} ext{tr}~m{T},|m{D}|)m{D}=0$$

see Rajagopal, J. Fluid Mech., 2006 (and Málek, Rajagopal, 2006, 2007)

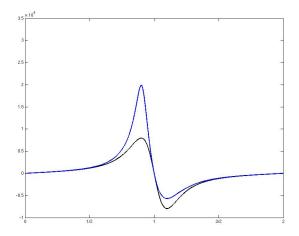
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Viscosity formulas used in applications

$$u = 
u(\pi, |\boldsymbol{D}(\boldsymbol{v})|^2) = \begin{cases} \sim \exp(lpha \pi), \\ \sim (1 + |\boldsymbol{D}(\boldsymbol{v})|^2)^{\frac{p-2}{2}}, & 1$$

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#### challenges and results,

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- subclass of models: restricted growth of extra stress with pressure
- well-posed PDEs, convergence of FEM, pictures

- > the lubrication works and is used since before the invention of wheel
- the viscosity-pressure relation is present in the very basis of the theory of elastohydrodynamic lubrication
- so why should we bother with mathematical theory now?

- > the lubrication works and is used since before the invention of wheel
- the viscosity-pressure relation is present in the very basis of the theory of elastohydrodynamic lubrication
- so why should we bother with mathematical theory now?
- Iubrication is used everywhere (transportation, electricity production)
- any optimization can save energy consumption and prolongate the lifespan
- quantitative predictions are needed !

- > the lubrication works and is used since before the invention of wheel
- the viscosity-pressure relation is present in the very basis of the theory of elastohydrodynamic lubrication
- so why should we bother with mathematical theory now?

## [Bair, Gordon, 2006]:

"  $\ldots$  there has been relatively little progress since the classic Newtonian solutions  $\ldots$  toward relating film thickness and traction to the properties of individual liquid lubricants

and it not clear at this time that full numerical solutions can even be obtained for heavily loaded contacts using accurate models.

One central issue is the validity of Reynolds equation, derived under the isoviscous assumption..."

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One central issue is the validity of Reynolds equation, derived under the isoviscous assumption. . .  $"\,$ 

Rajagopal, Szeri, Proc. R. Soc. Lond. A, 2003

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- so why should we bother with mathematical theory now?

#### [Babuška, Strouboulis, 2007 book]:

- It is necessary to realize that the FEM is a numerical method for constructing approximate solutions of a well defined mathematical problem. A necessary prerequisite is that the mathematical problem has 'good' properties: existence of solutions, possibly uniqueness, continuous dependence on input data.
- Mathematical formulation is a simplification of reality, and the existence of a physical solution which can be observed in an experiment, does not guarantee that the solution of the mathematical problem exists and has the expected properties following from physical intuition.

#### challenges and results,

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Mathematical formulation inside  $(0, T) \times \Omega$ :

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Problem well-posedness—first observations

$$\nu = \nu(\pi)$$

- M. Renardy, Comm. Part. Diff. Eq., 1986.
- ▶ F. Gazzola, Z. Angew. Math. Phys., 1997.
- ▶ F. Gazzola, P. Secchi, Navier-Stokes eq.: th. and num. meth. 1998.

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# Mathematical formulation inside $(0, T) \times \Omega$ :

$$\partial_{\tau} \boldsymbol{\nu} + \operatorname{div}(\boldsymbol{\nu} \otimes \boldsymbol{\nu}) - \operatorname{div} \boldsymbol{S} = -\nabla \pi + \boldsymbol{f},$$
  
$$\boldsymbol{S} = 2 \nu (\pi, |\boldsymbol{D}(\boldsymbol{\nu})|^2) \boldsymbol{D}(\boldsymbol{\nu})$$

Problem well-posedness-first positive results

$$\left| rac{\partial oldsymbol{\mathcal{S}}}{\partial oldsymbol{D}} \sim (1 + |oldsymbol{D}|^2)^{rac{p-2}{2}} \qquad \left| rac{\partial oldsymbol{\mathcal{S}}}{\partial \pi} 
ight| \leq \gamma_0 \left(1 + |oldsymbol{D}|^2\right)^{rac{p-2}{4}} \qquad 1$$

- Málek, Nečas, Rajagopal, Arch. Rational Mech. Anal., 2002.
- Hron, Málek, Nečas, Rajagopal, Math. Comput. Simulation, 2003.
- Málek, Rajagopal, Handbook of mathematical fluid dynamics, 2007.

# Mathematical formulation inside $(0, T) \times \Omega$ : $div \mathbf{v} = 0$ $\partial_{\tau} \mathbf{v} + div(\mathbf{v} \otimes \mathbf{v}) - div \mathbf{S} = -\nabla \pi + \mathbf{f},$ $\mathbf{S} = 2\nu(\pi, |\mathbf{D}(\mathbf{v})|^{2}) \mathbf{D}(\mathbf{v})$ on the boundary $(0, T) \times \partial \Omega = \Gamma_{D} \cup \Gamma_{N} \cup \Gamma_{P}$ : $\mathbf{v} \cdot \mathbf{n} = 0 \text{ and } - \mathbf{T}\mathbf{n} = \sigma \mathbf{v} \qquad \text{on } \Gamma_{N}$

Bulíček, Málek, Rajagopal, Indiana Univ. Math. J., 2007

Bulíček, Málek, Rajagopal, SIAM J. Math. Anal., 2009

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## Mathematical formulation inside $\Omega$ :

$$\begin{aligned} &\operatorname{div} \boldsymbol{v} &= 0\\ &\operatorname{div}(\boldsymbol{v} \otimes \boldsymbol{v}) - \operatorname{div} \boldsymbol{S} &= -\nabla \pi + \boldsymbol{f},\\ &\boldsymbol{S} &= 2 \nu(\pi, |\boldsymbol{D}(\boldsymbol{v})|^2) \boldsymbol{D}(\boldsymbol{v}) \end{aligned}$$

on the boundary 
$$\partial \Omega = \Gamma_D \cup \Gamma_N \cup \Gamma_P$$
:  
 $\mathbf{v} \cdot \mathbf{n} = 0 \text{ and } - \mathbf{T}\mathbf{n} = \sigma \mathbf{v} \quad \text{on } \Gamma_N$   
 $\mathbf{v} = \mathbf{v}_D \quad \text{on } \Gamma_D \quad \text{if } \Gamma_P = \emptyset,$   
 $-\mathbf{T}\mathbf{n} = \mathbf{b}(\mathbf{v}) \quad \text{on } \Gamma_P \quad \text{then } \int_{\Omega_0} \pi \, \mathrm{d}\mathbf{x} = 0$ 

- Franta, Málek, Rajagopal, Proc. Royal Soc. A, 2005
- Lanzendörfer, Nonlin. Anal.: Real World App., 2009

## Mathematical formulation inside $\Omega$ :

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- Stebel & Lanzendörfer, Appl. Mat.-Czech., 2011
- Stebel & Lanzendörfer, Math. Comput. Simulat., 2011
- Hirn, Stebel & Lanzendörfer, IMA J. Num. Anal., 2012 (el.)

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$$(q, \operatorname{div} \boldsymbol{w})_{\Omega} = 0$$
$$([\nabla \boldsymbol{v}] \boldsymbol{v}, \boldsymbol{w})_{\Omega} + (\boldsymbol{S}(\pi, \boldsymbol{D}(\boldsymbol{v})), \boldsymbol{D}(\boldsymbol{w}))_{\Omega} - (\pi, \operatorname{div} \boldsymbol{w})_{\Omega} = (\boldsymbol{f}, \boldsymbol{w})_{\Omega} - (\boldsymbol{b}(\boldsymbol{v}), \boldsymbol{w})_{\Gamma_{P}}$$

$$(q, \operatorname{div} \boldsymbol{w})_{\Omega} = 0$$
  
 $(\boldsymbol{S}(\pi, \boldsymbol{D}(\boldsymbol{v})), \boldsymbol{D}(\boldsymbol{w}))_{\Omega} - (\pi, \operatorname{div} \boldsymbol{w})_{\Omega} = (\boldsymbol{f}, \boldsymbol{w})_{\Omega} - (\boldsymbol{b} , \boldsymbol{w})_{\Gamma_{P}}$ 

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$$(q, \operatorname{div} \boldsymbol{w})_{\Omega} = 0$$
  
 $(\boldsymbol{S}(\pi, \boldsymbol{D}(\boldsymbol{v})), \boldsymbol{D}(\boldsymbol{w}))_{\Omega} - (\pi, \operatorname{div} \boldsymbol{w})_{\Omega} = (\boldsymbol{f}, \boldsymbol{w})_{\Omega} - (\boldsymbol{b} , \boldsymbol{w})_{\Gamma_{P}}$ 

Test eq. by solution

$$egin{aligned} & (m{S}(\pi,m{D}(m{v})),m{D}(m{v}))_\Omega \sim |m{D}(m{v})|^p \pm 1 \ & \|m{D}(m{v})\|_p \leq \mathcal{K} \implies \|m{v}\|_{1,p} + \|m{S}\|_{p'} \leq \mathcal{K} \end{aligned}$$

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$$(q, \operatorname{div} \boldsymbol{w})_{\Omega} = 0$$
  
 $(\boldsymbol{S}(\pi, \boldsymbol{D}(\boldsymbol{v})), \boldsymbol{D}(\boldsymbol{w}))_{\Omega} - (\pi, \operatorname{div} \boldsymbol{w})_{\Omega} = (\boldsymbol{f}, \boldsymbol{w})_{\Omega} - (\boldsymbol{b} \quad , \boldsymbol{w})_{\Gamma_{P}}$ 

Inf–sup inequality and the boundedness of  $\partial_{\pi} \boldsymbol{S}$ 

$$0 < \beta \leq \inf_{q \in \mathcal{L}^{p'}_{b.c.}(\Omega)} \sup_{\boldsymbol{w} \in \mathbf{W}^{1,p}_{b.c.}(\Omega)} \frac{(q, \operatorname{div} \boldsymbol{w})_{\Omega}}{\|q\|_{p'} \|\boldsymbol{w}\|_{1,p}}$$

$$\beta \|\pi\|_{p'} \le \|\boldsymbol{S}(\pi, \boldsymbol{D}(\boldsymbol{v}))\|_{p'} + \|\boldsymbol{f} + \boldsymbol{b}\| \le K$$

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$$egin{aligned} & (q, \operatorname{div} oldsymbol{w})_\Omega = 0 \ & (oldsymbol{S}(\pi, oldsymbol{D}(oldsymbol{v})), oldsymbol{D}(oldsymbol{w}))_\Omega - (\pi, \operatorname{div} oldsymbol{w})_\Omega = (oldsymbol{f}, oldsymbol{w})_\Omega - (oldsymbol{b} \ \ \ , oldsymbol{w})_{\Gamma_P} \end{aligned}$$

Inf–sup inequality and the boundedness of  $\partial_{\pi} \boldsymbol{S}$ 

$$0 < \beta \leq \inf_{q \in \mathcal{L}_{b,c.}^{p'}(\Omega)} \sup_{\boldsymbol{w} \in \boldsymbol{W}_{b,c.}^{1,p}(\Omega)} \frac{(q, \operatorname{div} \boldsymbol{w})_{\Omega}}{\|q\|_{p'} \|\boldsymbol{w}\|_{1,p}}$$

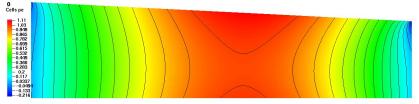
Pressure uniquely determined by velocity?

$$\begin{split} \beta \|\pi^1 - \pi^2\|_{p'} &\leq \|\boldsymbol{S}(\pi^1, \boldsymbol{D}(\boldsymbol{v})) - \boldsymbol{S}(\pi^2, \boldsymbol{D}(\boldsymbol{v}))\|_{p'} \leq \left\| \int_{\pi^1}^{\pi^2} \frac{\partial \boldsymbol{S}(\pi, \boldsymbol{D}(\boldsymbol{v}))}{\partial \pi} \mathrm{d}\pi \right\|_{p'} \\ &\leq \gamma_0 \|\pi^1 - \pi^2\|_{p'} \quad \text{where} \quad \left| \frac{\partial \boldsymbol{S}}{\partial \pi} \right| \leq \gamma_0 \end{split}$$

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#### Flow in a converging channel

Newtonian model  $\nu = const$ 



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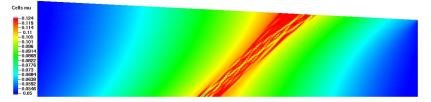
#### Flow in a converging channel Barus model $\nu = \exp(\alpha \pi)$ , $\alpha = 0.306$ Cells pe 2.72 2.54 2.36 2.18 2 1.82 1.64 1.46 1.28 1.1 -0.923 -0.743 -0.564 0.205 0.025 Cells mu -0.115 -0.102 0 103 0.086 0.082 0.078 0.0744 0.0703 0.0662 0.0622 0.0581 0.054

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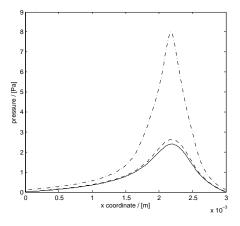
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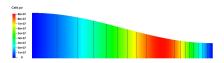
#### Flow in a converging channel Barus model $\nu = \exp(\alpha \pi)$ , $\alpha = 0.3061$



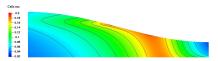
## Sensitivity on boundary data



Pressure for  $\pi(0) = \pi(L) = 0$  (full), 1 MPa (dashed) a 10 MPa (dash dotted).



pressure  $\pi$ 



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viskosity  $\nu(\pi)$ 

#### Conclusion

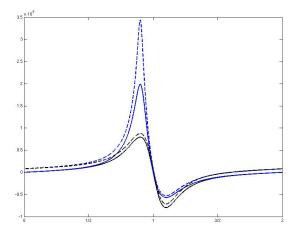
mathematical theory being build succesfully for

$$rac{\partial oldsymbol{\mathcal{S}}}{\partial oldsymbol{D}} \sim \left(1 + |oldsymbol{D}|^2
ight)^{rac{p-2}{2}} \qquad \left|rac{\partial oldsymbol{\mathcal{S}}}{\partial \pi}
ight| \leq \gamma_0 \left(1 + |oldsymbol{D}|^2
ight)^{rac{p-2}{4}} \qquad 1$$

- ▶ for  $|\partial \pmb{S} / \partial \pi| > 1$  both the theory and the standard numerical approach fails
- is the model well-posed?

#### Thank you for your attention

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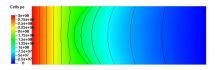
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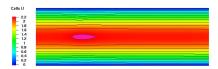
## Poiseuille flow



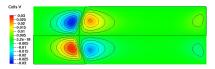
Cols and 1990

the pressure  $\boldsymbol{\pi}$ 

the viscosity  $\nu(p)$ 



the velocity component  $\boldsymbol{v} \cdot \boldsymbol{e}_{x}$ 



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the velocity component  $\boldsymbol{v} \cdot \boldsymbol{e}_{y}$