Robust adaptive *hp* discontinuous Galerkin finite element methods for the Helmholtz equation

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Overview



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Introduction

- Helmholtz equation
- Discontinuous Galerkin FEM
- Error estimation and mesh adaptation
- Problems with FEM for Helmholtz

2 hp-robust error estimation

- Helmholtz as shifted Poisson
- A posteriori error estimator
- Reliability
- Reliability
- Suitable localized reconstructions
- Efficiency

Interpretended in the second secon

Section 1

Introduction



Let $\Omega \subset \mathbb{R}^d$, d = 2 be a bounded polygonal domain. We seek $u : \Omega \mapsto \mathbb{C}$ such that

 $-\Delta u - k^2 u = f \qquad \text{in } \Omega,$ $\nabla u \cdot \boldsymbol{n} - iku = g \qquad \text{on } \partial\Omega, \qquad (\text{Robin/Impedence BC})$

where

$$k = \frac{\omega L}{c}$$

is the wavenumber (ω is the frequency of the wave, L is the measure of the domain, and c is the speed of sound in the material). Wavenumber is related to the wave length

$$\lambda = \frac{2\pi}{k}.$$

Discontinuous Galerkin FEM



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Multiplying by a test function and integrating by parts gives the weak formulation: Find $u \in H^1(\Omega)$ such that

$$\int_{\Omega} (\nabla u \cdot \nabla \bar{v} - k^2 u \bar{v}) \, d\mathbf{x} + \int_{\Gamma_R} i k u \bar{v} \, ds = \int_{\Omega} f \bar{v} \, d\mathbf{x} + \int_{\Gamma_R} g \cdot \mathbf{n} \bar{v} \, ds$$

for all $v \in H^1(\Omega)$.

Well-posedness: [Melenk, 1995]

Discontinuous Galerkin FEM

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for all $v \in H^1(\Omega)$. We want to search for a solution in a finite dimensional subspace of $H^1(\Omega)$. To that end we.

- subdivide the domain Ω into a mesh \mathcal{T}_h of non-overlapping triangles \mathcal{T} , where each element has a size h_T and denote by \mathcal{E}'_h the union of all interior edges in the mesh
- multiple by test functions v and integrate by parts elementwise

$$\sum_{T\in\mathcal{T}_h}\left(\int_T (\nabla u\cdot\nabla\overline{v}-k^2u\overline{v})\,d\mathbf{x}-\int_{\partial T}\nabla u\cdot\boldsymbol{n}_T\overline{v}\,ds\right)=\sum_{T\in\mathcal{T}_h}\int_T f\overline{v}\,d\mathbf{x}.$$



Replace continuous functions u, v by discrete functions u_{hp}, v_{hp} in the finite element space

$$V_{hp} = \left\{ v_{hp} \in L^2(\Omega) : \left. v_{hp} \right|_{T} \in \mathbb{P}_{p_T}(T) \text{ for all } T \in \mathcal{T}_h \right\}$$

$$\mathbb{P}_{p_{\mathcal{T}}}(\mathcal{T}) = \mathsf{polynomials} ext{ of degree} \leq p_{\mathcal{T}} ext{ in } \mathcal{T}$$

symmetrize and add jump-penalty terms on \mathcal{E}_h^I and $\partial \Omega$, without losing consistency

$$\{\!\!\{\cdot\}\!\!\} = mean value \qquad [\![\cdot]\!]_N = jump in normal direction$$



DG method

Find
$$u_{hp} \in V_{hp}$$
 such that $a_{hp}(u_{hp}, v_{hp}) = \ell_{hp}(v_{hp})$ for all $v_{hp} \in V_{hp}$.

$$\begin{split} a_{h}(u,v) &:= \int_{\Omega} (\nabla_{h} u \cdot \nabla_{h} \overline{v} - k^{2} u \overline{v}) \, d\mathbf{x} - \int_{\mathcal{E}_{h}^{l}} (\llbracket u \rrbracket_{N} \{\!\!\{\nabla_{h} \overline{v}\}\!\!\} + \{\!\!\{\nabla_{h} u\}\!\}_{I\!\!I} \llbracket \overline{v} \rrbracket_{N}) \, ds \\ &- \mathrm{i}\beta \int_{\mathcal{E}_{h}^{l}} \frac{\mathrm{h}}{\mathrm{p}} \llbracket \nabla_{h} u \rrbracket_{N} \llbracket \nabla_{h} \overline{v} \rrbracket_{N} \, ds - \mathrm{i}\alpha \int_{\mathcal{E}_{h}^{l}} \frac{\mathrm{p}^{2}}{\mathrm{h}} \llbracket u \rrbracket_{N} \cdot \llbracket \overline{v} \rrbracket_{N} \, ds \\ &- \gamma \int_{\partial \Omega} k \frac{\mathrm{h}}{\mathrm{p}} (u \nabla_{h} \overline{v} \cdot \mathbf{n} + \nabla_{h} u \cdot \mathbf{n} \overline{v}) \, ds \\ &- \mathrm{i}\gamma \int_{\partial \Omega} \frac{\mathrm{h}}{\mathrm{p}} \nabla_{h} u \cdot \mathbf{n} \nabla_{h} \overline{v} \cdot \mathbf{n} \, ds - \mathrm{i} \int_{\partial \Omega} k \left(1 - \gamma k \frac{\mathrm{h}}{\mathrm{p}}\right) u \overline{v} \, ds \\ &\ell_{hp}(v) := \int_{\Omega} f \overline{v} \, d\mathbf{x} + i\gamma \int_{\partial \Omega} \frac{\mathrm{h}}{\mathrm{p}} g \, \nabla_{h} \overline{v} \cdot \mathbf{n} \, ds + \int_{\partial \Omega} \left(1 - \gamma k \frac{\mathrm{h}}{\mathrm{p}}\right) g \overline{v} \, ds \end{split}$$

We use $\alpha = 10$, $\beta = 1$ and $\gamma = 1/4$. a priori analysis proves well-posedness and quasi-optimal error estimates [Melenk, Parsania, Sauter 2013]



Adaptive mesh refinement:

- Refine elements using either:
 - element subdivision (h-refinement)
 - increasing polynomial degree of the element (*p*-refinement)
- Need to estimate which elements need refining; therefore, need computable (a posteriori) local error indicators η_T for each element

A global *a posteriori* error estimate can be computed by summing local error indicators, which can be used to estimate when the actual error reaches a desired accuracy.

error \lesssim estimate	\implies	reliability
estimate \lesssim error	\implies	efficiency

If the constants in these inequalities are independent of h and p we have a robust estimate.



Problems with FEM:

- Number of degrees of freedom required to obtain given accuracy increases with wave number k.
- Error: best approximation + phase lag:

$$\|
abla_h(u-u_h)\|_{L^2(\Omega)} \lesssim (kh)^p + k(kh)^{2p};$$



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convergence like the best approximation when $k(kh)^{2p} \leq (kh)^{p}$, i.e.



Section 2

hp-robust error estimation



[C., Gedicke, Perugia, SISC 2019]

The observation that the pollution effect is related to the phase lead of the numerical solution leads to the definition of the shifted solution as the key auxiliary construction for the analysis of local a posteriori error estimation.

— Babuška, Ihlenburg, Strouboulis, Gangaraj 1997



[C., Gedicke, Perugia, SISC 2019]

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— Babuška, Ihlenburg, Strouboulis, Gangaraj 1997

We consider the Helmholtz problem as a shifted Poisson problem, and use methods based on equilibrated fluxes and potential reconstructions; cf. [Braess, Pillwein, Schöberl 2009], [Ern, Vohralík 2015], [Dolejší, Ern, Vohralík 2016].

The idea of using a shifted Poisson is related to the *a posteriori* error analysis for eigenvalue problems.

[Cancès, Dusson, Maday, Stamm, Vohralík 2017 & 2018]



Helmholtz:

$$\begin{aligned} -\Delta u - k^2 u &= f & \text{in } \Omega, \\ \nabla u \cdot \boldsymbol{n} - \mathrm{i} k u &= g & \text{on } \partial \Omega. \end{aligned}$$

Shifted Poisson (Neumann):

$$\begin{aligned} -\Delta w &= f + k^2 u_{hp} & \text{in } \Omega, \\ \nabla w \cdot \boldsymbol{n} &= g + i k u_{hp} - \gamma k (g - \nabla u_{hp} \cdot \boldsymbol{n} + i k u_{hp}) & \text{on } \partial \Omega. \end{aligned}$$

The extra term on the right hand side is required for compatibility, which is necessary to prove that the flux reconstruction we define below exists.



Definition (Flux reconstruction)

Given $u_{hp} \in V_{hp}$, a equilibrated flux reconstruction for u_{hp} is any function $\sigma_{hp} \in H(\text{div}; \Omega)$ that satisfies

$$\int_{T} \operatorname{div} \boldsymbol{\sigma}_{hp} \, d\mathbf{x} = \int_{T} (f + k^{2} u_{hp}) \, d\mathbf{x} \qquad \forall T \in \mathcal{T}_{h},$$
$$\int_{E} \boldsymbol{\sigma}_{hp} \cdot \boldsymbol{n} = -\int_{E} (g + \mathrm{i}k u_{hp} - \gamma k (g - \nabla u_{hp} \cdot \boldsymbol{n} + \mathrm{i}k u_{hp})) \, ds \quad \forall E \subset \partial \Omega$$

Definition (Potential (reconstruction))

We define a potential as any function

$$s_{hp} \in H^1_*(\Omega) \coloneqq \left\{ v \in H^1(\Omega) : rac{1}{|\Omega|} \int_\Omega v \ d{f x} = 0
ight\}.$$

Note σ_{hp} and s_{hp} are not necessarily piecewise polynomial (yet).

A posteriori error estimator



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For $u_{hp} \in V_{hp}$ we denote by $\mathcal{G}(u_{hp})$ its DG gradient:

$$\mathcal{G}(u_{hp}) \coloneqq \nabla_h u_{hp} - \sum_{E \in \mathcal{E}_h^{\prime}} \mathcal{L}_E^0(\llbracket u_{hp} \rrbracket) - \sum_{E \in \mathcal{E}_h^{\prime}} \mathcal{L}_E^1(\llbracket \nabla u_{hp} \rrbracket)$$

Error Estimator

$$\begin{split} \eta_{hp} &\coloneqq \sum_{T \in \mathcal{T}_h} \left(\|\mathcal{G}(u_{hp}) + \boldsymbol{\sigma}_{hp}\|_{0,T} + \frac{h_T}{j_{1,1}} \|f + k^2 u_{hp} - \operatorname{div} \boldsymbol{\sigma}_{hp}\|_{0,T} \right. \\ &+ C_{tr} \sum_{E \subset \partial T \cap \partial \Omega} h_E^{1/2} \|\boldsymbol{\sigma}_{hp} \cdot \boldsymbol{n} + g + \operatorname{i} k u_{hp} - \gamma k \frac{h}{p} (g - \nabla_h u_{hp} \cdot \boldsymbol{n} + \operatorname{i} k u_{hp})\|_{0,E} \right)^2 \\ &+ \sum_{T \in \mathcal{T}_h} \|\mathcal{G}(u_{hp}) - \nabla s_{hp}\|_{0,T}^2 \end{split}$$

Any admissable flux reconstructions and potentials

 \implies reliability

(error \leq estimator)

Suitable localized flux and potential reconstructions

 \implies efficiency & robustness

Scott Congreve (Charles University)

Robust Adaptive hp-DG for Helmholtz

(estimator \leq error) KNM Seminar 14/34



Theorem (Reliability)

$$\begin{split} \|\nabla u - \mathcal{G}(u_{hp})\|_{0,\Omega} \lesssim \eta_{hp} + k^2 \|u - u_{hp}\|_{0,\Omega} + k \|u - u_{hp}\|_{0,\partial\Omega} \\ + \|\gamma k \frac{\mathsf{h}}{\mathsf{p}} (g - \nabla_h u_{hp} + \mathsf{i} k u_{hp})\|_{0,\partial\Omega} \end{split}$$

The additional terms are higher order compared to the left-hand side providing that the resolution condition is met. [Sauter, Zech 2015]

Proof is similar to [Ern, Vohralík 2015].

Reliability



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Proof.

Let $s \in H^1_*(\Omega)$ be defined by the projection $\int_{\Omega} \nabla s \cdot \nabla \overline{v} \, d\mathbf{x} = \int_{\Omega} \mathcal{G}(u_{hp}) \cdot \nabla \overline{v} \, d\mathbf{x} \qquad \forall v \in H^1(\Omega).$ Then by orthogonality, for any $s_{hp} \in H^1_*(\Omega)$, $\|\nabla u - \mathcal{G}(u_{hp})\|_{0,\Omega}^2 = \|\nabla (u - s)\|_{0,\Omega}^2 + \|\nabla s - \mathcal{G}(u_{hp})\|_{0,\Omega}^2$ $= \|\nabla(u-s)\|_{0,\Omega}^2 + \min_{v \in H^1(\Omega)} \|\nabla v - \mathcal{G}(u_{hp})\|_{0,\Omega}^2$ $\leq \|\nabla(u-s)\|_{0,\Omega}^2 + \|\nabla s_{hp} - \mathcal{G}(u_{hp})\|_{0,\Omega}^2$ $= \sup_{v \in H^{1}_{\omega}(\Omega), \|\nabla v\| = 1} \int_{\Omega} \nabla (u - s) \cdot \nabla \overline{v} \, d\mathbf{x} + \|\nabla s_{hp} - \mathcal{G}(u_{hp})\|^{2}_{0,\Omega}$ $= \sup_{\mathbf{y}\in H^1(\Omega), \|\nabla \mathbf{y}\|=1} \int_{\Omega} \nabla (u - \mathcal{G}(u_{hp})) \cdot \nabla \overline{\mathbf{y}} \, d\mathbf{x} + \|\nabla s_{hp} - \mathcal{G}(u_{hp})\|_{0,\Omega}^2$

The first term is then boundable by definitions of weak formulation and σ_{hp} , Poincaré and trace inequalities, interpolation results, etc.



We have triangular meshes with no hanging nodes, so we can define for each node $z \in \mathcal{N}$:

- nodal patch ω_z
- nodal hat functions \u03c6_z (forms partition of unity)





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- nodal patch ω_z
- nodal hat functions \u03c6_z (forms partition of unity)



For proof of efficiency we define specific flux and potential reconstructions locally by mesh nodes:

$$egin{aligned} &\sigma_{hp} \coloneqq \sum_{z \in \mathcal{N}} oldsymbol{\zeta}_{hp}^z \ &s_{hp} \coloneqq \widetilde{s}_{hp} - rac{1}{|\Omega|} \int_\Omega \widetilde{s}_{hp} \, d\mathbf{x} & \widetilde{s}_{hp} \coloneqq \sum_{z \in \mathcal{N}} s_{hp}^z \end{aligned}$$

For the local flux reconstruction we solve local patch problems using Raviart–Thomas (RT_p) elements:



For all any domain $D \subset \mathbb{R}^2$ Raviart–Thomas elements have the properties that:

Local flux reconstruction



We solve the following local problem in mixed form for every node $z \in \mathcal{N}$: Find $(\zeta_{hp}^z, r_{hp}^z) \in \Sigma_{g^z, hp}^z \times Q_{hp}^z$ such that

$$\int_{\omega_z} \left(\zeta_{hp}^z \cdot \overline{\tau}_{hp} - r_{hp}^z \operatorname{div} \overline{\tau}_{hp} \right) \, d\mathbf{x} = - \int_{\omega_z} \psi_z \mathcal{G}(u_{hp}) \cdot \overline{\tau}_{hp} \, d\mathbf{x} \quad \forall \boldsymbol{\tau}_{hp} \in \Sigma_{0,hp}^z$$
$$\int_{\omega_z} \operatorname{div} \zeta_{hp}^z \overline{q}_{hp} \, d\mathbf{x} = \int_{\omega_z} f_z \overline{q}_{hp} \, d\mathbf{x} \qquad \forall q_{hp} \in Q_{hp}^z$$

where, for $p_z \ge 1$,

$$\begin{split} \Sigma_{g^z,hp}^z &\coloneqq \left\{ \boldsymbol{\tau}_{hp} \in \boldsymbol{H}(\operatorname{div},\omega_z) : \boldsymbol{\tau}_{hp} |_{\mathcal{T}} \in \operatorname{RT}_{p_z}(\mathcal{T}) \text{ for all } \mathcal{T} \in \mathcal{T}_h(z), \\ \boldsymbol{\tau}_{hp} \cdot \boldsymbol{n} &= 0 \text{ on } \partial \omega_z \setminus \partial \Omega, \\ \boldsymbol{\tau}_{hp} \cdot \boldsymbol{n} |_E &= \Pi_E^{p_z} g^z \text{ for all } E \subset \partial \omega_z \cap \partial \Omega \right\} \\ Q_{hp}^z &\coloneqq \left\{ q_{hp} \in Q_{hp}(\omega_z) : |\omega_z|^{-1} \int_{\omega_z} q_{hp} \, d\mathbf{x} = 0 \right\} \\ Q_{hp}(\omega_z) &\coloneqq \left\{ q_{hp} \in L^2(\omega_z) : q_{hp} |_{\mathcal{T}} \in \mathbb{P}_{p_z}(\mathcal{T}) \text{ for all } \mathcal{T} \in \mathcal{T}_h(z) \right\} \\ f^z &\coloneqq (f + k^2 u_{hp}) \psi_z - \mathcal{G}(u_{hp}) \cdot \nabla \psi_z, \qquad g_z \coloneqq \ldots \end{split}$$



Lemma (Flux reconstruction)

 $\sigma_{hp} = \sum_{z \in \mathcal{N}} \zeta_{hp}^{z}$ is an equilibrated flux reconstruction in $H(\operatorname{div}, \Omega)$, which satisfies for any $T \in \mathcal{T}_{h}$

$$\int_{T} \left(f + k^2 u_{hp} - \operatorname{div} \boldsymbol{\sigma}_{hp} \right) \overline{q}_{hp} \, d\mathbf{x} = 0,$$

for all $q_{hp} \in igcap_{z \in \mathcal{N}(T)} Q_{hp}(\omega_z)|_T$, and for any $E \subset \partial \Omega$

$$\int_{E} \left(\boldsymbol{\sigma}_{hp} \cdot \boldsymbol{n} + g + \mathrm{i} k u_{hp} - \gamma k \frac{\mathrm{h}}{\mathrm{p}} (g - \nabla_{h} u_{hp} \cdot \boldsymbol{n} + \mathrm{i} k u_{hp}) \right) \overline{q}_{hp} \, ds = 0,$$

for all $q_{hp} \in \bigcap_{z \in \mathcal{N}(E)} Q_{hp}(\omega_z)|_E$.



Proof.

- $\zeta_{hp}^z \in H(\operatorname{div}, \Omega)$ by zero extension $\implies \sigma_{hp} \in H(\operatorname{div}, \Omega)$.
- Prove just the first statement (second follows similarly):

$$\int_{T} \left(f + k^{2} u_{hp} - \operatorname{div} \boldsymbol{\sigma}_{hp} \right) \overline{q}_{hp} \, d\mathbf{x}$$

$$= \sum_{z \in \mathcal{N}(T)} \int_{T} \left(\psi_{z} (f + k^{2} u_{hp}) - \operatorname{div} \boldsymbol{\zeta}_{hp}^{z} \right) \overline{q}_{hp} \, d\mathbf{x},$$

$$= \sum_{z \in \mathcal{N}(T)} \int_{T} \left(f^{z} + \mathcal{G}(u_{hp}) \cdot \nabla \psi_{z} - \operatorname{div} \boldsymbol{\zeta}_{hp}^{z} \right) \overline{q}_{hp} \, d\mathbf{x},$$

$$= \sum_{z \in \mathcal{N}(T)} \int_{T} \mathcal{G}(u_{hp}) \cdot \nabla \psi_{z} \overline{q}_{hp} \, d\mathbf{x}, \qquad \left(\sum_{z \in \mathcal{N}(T)} \nabla \psi_{z} = 0 \right)$$

$$= 0.$$



Theorem (Flux reconstruction efficiency)

$$\begin{split} \|\mathcal{G}(u_{hp}) + \boldsymbol{\sigma}_{hp}\|_{0,\Omega} &\lesssim \|\nabla u - \mathcal{G}(u_{hp})\|_{0,\Omega} + k^2 \|u - u_{hp}\|_{0,\Omega} \\ &+ k \|u - u_{hp}\|_{0,\partial\Omega} + \operatorname{osc}(f) + \operatorname{osc}(g) \\ &+ \|\gamma k \frac{\mathsf{h}}{\mathsf{p}} (g - \nabla_h u_{hp} \cdot \boldsymbol{n} + \mathsf{i} k u_{hp})\|_{0,\partial\Omega} \\ &+ \|\mathsf{i} \gamma \frac{\sqrt{\mathsf{h}}}{\mathsf{p}} (g - \nabla_h u_{hp} \cdot \boldsymbol{n} + \mathsf{i} k u_{hp})\|_{0,\partial\Omega} \end{split}$$



Proof.

Define $r^z \in H^1_*(\omega_z)$ as the solution of the continuous problem

$$\begin{split} \int_{\omega_z} \nabla r^z \cdot \nabla \overline{v} \, d\mathbf{x} &= -\int_{\omega_z} \psi_z \mathcal{G}(u_{hp}) \cdot \nabla \overline{v} \, d\mathbf{x} + \sum_{T \in \mathcal{T}_h(z)} \int_T \Pi_T^{p_z} f^z \overline{v} \, d\mathbf{x} \\ &- \sum_{E \subset \omega_z \cap \partial \Omega} \int_E \Pi_E^{p_z} g^z \overline{v} \, ds \qquad \quad \forall v \in H^1(\omega_z). \end{split}$$

$$\begin{split} \|\mathcal{G}(u_{hp}) + \sigma_{hp}\|_{0,\Omega} &\leq \sum_{z \in \mathcal{N}} \|\psi_{z}\mathcal{G}(u_{hp}) + \zeta_{hp}^{z}\|_{0,\omega_{z}} \\ &\lesssim \sum_{z \in \mathcal{N}} \|\nabla r^{z}\|_{0,\omega_{z}} \quad [\text{Braess, Pillwein, Schöberl 2009}] \\ &\lesssim \|\nabla w - \mathcal{G}(u_{hp})\|_{0,\Omega} + (\text{terms on RHS}) \\ &\leq \|\nabla u - \mathcal{G}(u_{hp})\|_{0,\Omega} + \sup_{v \in H^{1}_{*}(\Omega), \|v\|_{0,\Omega} = 1} \int_{\Omega} \nabla(w - u) \cdot \nabla \overline{v} \, d\mathbf{x} + \dots \end{split}$$



For simplicity we consider only interior nodes. For $p_z \ge 1$

 $V_{hp}^{z} := \{ v_{hp} \in C^{0}(\overline{\omega_{z}}) : v_{hp} |_{T} \in \mathbb{P}_{p_{z}+1}(T) \, \forall T \in \mathcal{T}_{h}(z), v_{hp} = 0 \text{ on } \partial \omega_{z} \}.$ We then define

$$s_{hp}^{z} \coloneqq \underset{v_{hp} \in V_{hp}^{z}}{\operatorname{arg\,min}} \| \nabla_{h}(\psi_{z} u_{hp}) - \nabla v_{hp} \|_{0,\omega_{z}},$$

which is equivalent to finding $s_{hp}^z \in V_{hp}^z$ such that

$$\int_{\omega_z} \nabla s_{hp}^z \cdot \nabla \overline{v}_{hp} \, d\mathbf{x} = \int_{\omega_z} \nabla_h (\psi_z u_{hp}) \cdot \nabla \overline{v}_{hp} \, d\mathbf{x}, \qquad \text{for all } v_{hp} \in V_{hp}^z.$$



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 $V_{hp}^{z} := \{ v_{hp} \in C^{0}(\overline{\omega_{z}}) : v_{hp} |_{T} \in \mathbb{P}_{p_{z}+1}(T) \, \forall T \in \mathcal{T}_{h}(z), v_{hp} = 0 \text{ on } \partial \omega_{z} \}.$ We then define

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which is equivalent to finding $s_{hp}^z \in V_{hp}^z$ such that

$$\int_{\omega_z} \nabla s_{hp}^z \cdot \nabla \overline{v}_{hp} \, d\mathbf{x} = \int_{\omega_z} \nabla_h (\psi_z u_{hp}) \cdot \nabla \overline{v}_{hp} \, d\mathbf{x}, \qquad \text{for all } v_{hp} \in V_{hp}^z.$$

Lemma (Potential reconstruction)

$$s_{hp} \coloneqq \widetilde{s}_{hp} - rac{1}{|\Omega|} \int_{\Omega} \widetilde{s}_{hp} \, d\mathbf{x}, \qquad ext{where } \widetilde{s}_{hp} \coloneqq \sum_{z \in \mathcal{N}} s_{hp}^z$$

is a potential.



Theorem (Potential reconstruction efficiency)

$$\begin{split} \|\mathcal{G}(u_{hp}) + \nabla s_{hp}\|_{0,\Omega}^2 \lesssim \|\nabla(u - u_{hp})\|_{0,\Omega}^2 + \sum_{E \subset \partial \Omega} h_E^{-1} \|\Pi_E^0 \llbracket u_{hp} \rrbracket \|_{0,E}^2 \\ + \sum_{E \subset \partial \Omega} \beta^2 h_E \|\mathbf{p}^{-1} \Pi_E^0 \llbracket \nabla u_{hp} \rrbracket \|_{0,E}^2 \end{split}$$

Proof.

We can write the above minimization in mixed form and show it has an underlying continuous problem: Find $r^z \in H^1_*(\omega_z)$ such that:

$$\int_{\omega_z} \nabla r^z \cdot \nabla \overline{v} \, d\mathbf{x} = - \int_{\omega_z} \operatorname{rot}_h(\psi_z \, u_{hp}) \cdot \nabla \overline{v} \, d\mathbf{x} \qquad \forall v \in H^1(\omega_z).$$

Following [Ern & Vohralík, 2015] we can show that

$$\begin{aligned} \|\nabla(u_{hp}-s_{hp})\|_{0,\mathcal{T}} &\lesssim \sum_{z\in\mathcal{N}(\mathcal{T})} \|\nabla r^{z}\|_{0,\omega_{z}} \\ \|\nabla r^{z}\|_{0,\omega_{z}}^{2} &\lesssim \|\nabla(u-u_{hp})\|_{0,\omega_{z}}^{2} + \sum_{E\subset\partial\omega_{z}\cap\partial\Omega} h_{E}^{-1} \|\Pi_{E}^{0}\llbracket u-u_{hp}]\!]\|_{0,E}^{2} \end{aligned}$$

Efficiency



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Theorem (Efficiency)

$$\begin{split} \eta_{hp} &\lesssim \|\nabla(u - u_{hp})\|_{0,\Omega} + k^2 \|u - u_{hp}\|_{0,\Omega} + k \|u - u_{hp}\|_{0,\partial\Omega} \\ &+ \operatorname{osc}(f) + \operatorname{osc}(g) + \|\gamma k \frac{\mathsf{h}}{\mathsf{p}} (g - \nabla_h u_{hp} \cdot \boldsymbol{n} + \mathrm{i} k u_{hp})\|_{0,\partial\Omega} \\ &+ \|\mathrm{i} \gamma \frac{\sqrt{\mathsf{h}}}{\mathsf{p}} (g - \nabla_h u_{hp} \cdot \boldsymbol{n} + \mathrm{i} k u_{hp})\|_{0,\partial\Omega} \\ &+ \left(\sum_{E \subset \partial\Omega} h_E^{-1} \|\Pi_E^0 \llbracket u_{hp} \rrbracket \|_{0,E}^2 \right)^{1/2} + \left(\sum_{E \subset \partial\Omega} \beta^2 h_E \|\mathsf{p}^{-1} \Pi_E^0 \llbracket \nabla u_{hp} \rrbracket \|_{0,E}^2 \right)^{1/2}. \end{split}$$

Proof.

The flux reconstruction efficiency and potential reconstruction efficiency bounds the first and last terms of η_{hp} respectively. The other terms are bound by the oscillation terms osc(f) and osc(g) respectively.

Section 3

hp-adaptive mesh refinement

hp-adaptive mesh refinement



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We can re-write the error indicator in local form as
$$\eta_{hp} = \sum_{T \in \mathcal{T}_h} \eta_T$$
.

hp-adaptive mesh refinement

Construct mesh \mathcal{T}_h , with uniform polynomial degree, and FE space V_{hp} .

for j = 0, 1, 2, ... do Solve $a_{hp}(u_{hp}, v_{hp}) = \ell_{hp}(v_{hp})$ on $V_{hp}^{(0)}$ Compute ζ_{hp}^z and s_{hp}^z (in parallel) and η_T for $T \in T_h$ do if $\eta_T \ge \theta \max_{T \in T_h} \eta_T$ then Perform *h*- or *p*-refinement on *T* [Melenk & Wohlmuth 2001] end if end for Perform mesh smoothing (remove hanging nodes) end for

To compute ζ_{hp}^z we let $p_z = \max_{T \subset \omega_z} p_T + 1$ and to compute s_{hp}^z we let $p_z = \max_{T \subset \omega_z} p_T$. We use $\theta = 0.75$ (maximum marking strategy).



hp-refinement strategy [Melenk & Wohlmuth 2001]

Decision on h- or p-refinement as performed as follows:

if T marked for refinement then if $\eta_T > \eta_T^{\text{pred}}$ then *h*-refinement: Divide T into 2 (T_{\pm}) using newest vertex bisection $(\eta_{T_{\perp}}^{\text{pred}})^2 \leftarrow \frac{1}{2} \gamma_h \left(\frac{1}{2}\right)^{p_T} \eta_T^2$ else *p*-refinement: $p_T \leftarrow p_T + 1$ $(\eta_{\tau}^{\text{pred}})^2 \leftarrow \gamma_{\rho} \eta_{T}^2$ end if else $(\eta_{\tau}^{\text{pred}})^2 \leftarrow \gamma_n (\eta_{\tau}^{\text{pred}})^2$ end if

We use
$$\gamma_h=$$
 4, $\gamma_{p}=$ 0.4, $\gamma_n=$ 1, and set $\eta_{T}^{
m pred}=\infty$ initially.

Square domain example



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Let $\Omega = (0,1)^2$, f = 0, and select g such that the analytical solution is

$$u(\mathbf{x}) = \mathcal{H}_0^{(1)}\left(k\sqrt{(x_1+1/4)^2+x_2^2}\right),$$

where $\mathcal{H}_0^{(1)} =$ Hankel function of the first kind. We consider k = 20, 50.





k = 20

Square domain example



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Scott Congreve (Charles University)

L-shaped domain example



Let $\Omega = (-1,1)^2 \setminus ((0,1) \times (-1,0))$, f = 0, and select g such that

$$u(r,\varphi) = \mathcal{J}_{2/3}(kr)\sin(2\varphi/3),$$

where $\mathcal{J}_{2/3}$ denotes the Bessel function of first kind. We consider k = 20, 50.



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Internal reflection/refraction



BLES UNIVERSITY of mathematics and physics

We now consider a wavenumber k given by the piecewise constant function

$$k(x, y) = \begin{cases} k_1 := \omega n_1 & \text{if } y \leq 0, \\ k_2 := \omega n_2 & \text{if } y > 0, \end{cases}$$

where, we let $\omega = 20$, $n_1 = 2$, and $n_2 = 1$, with appropriate boundary conditions, such that , for a constant $0 \le \theta_i \le \pi/2$,

$$u(x,y) = \begin{cases} T e^{i(K_1 x + K_2 y)} & \text{if } y > 0, \\ e^{ik_1(x\cos(\theta_i) + y\sin(\theta_i))} + R e^{ik_1(x\cos(\theta_i) - y\sin(\theta_i))} & \text{if } y < 0, \end{cases}$$

where $K_1 = k_1 \cos(\theta_i)$, $K_2 = \sqrt{k_2^2 - k_1^2 \cos^2(\theta_i)}$,

$$R = -\frac{K_2 - k_1 \sin(\theta_i)}{K_2 + k_1 \sin(\theta_i)},$$

and T = 1 + R.



















Summary:

- a posteriori error estimator based for Helmholtz
- Shown reliability and efficiency providing resolution condition met
- Demonstrated robust in polynomial degree

Further work:

- The analysis of the potential reconstruction is 2D only. [Ern, Vohralík 2017] has argument for 3D.
- Trefftz discontinuous Galerkin FEM