Adaptive Refinement for hp-version Trefftz Discontinuous Galerkin Methods for the Homogeneous Helmholtz Problem

Scott Congreve

Fakultät für Mathematik. Universität Wien

loint work with Ilaria Perugia (Universität Wien) Paul Houston (University of Nottingham)

16th European Finite Element Fair (Heidelberg)



Trefftz DG for Helmholtz

- Helmholtz Equation
- Trefftz DG

2 Adaptive Refinement

- Plane Wave Direction Refinement
- A posteriori Error Estimates
- *hp*-adaptive Refinement



Let $\Omega \subset \mathbb{R}^d$, d = 2, 3 be a bounded polygonal/polyhedral domain.

$-\Delta u - k^2 u = 0$	in Ω ,
<i>u</i> = 0	on Γ_D ,
$\nabla u \cdot \mathbf{n} + i k \vartheta u = g_R$	on Γ_R .

Trefftz FEM Spaces



Polynomial DG Finite Element Spaces: DGFEM uses polynomial basis functions defined on a reference element \widehat{K} :

$$V_q^{DG}(\mathcal{T}_h) \coloneqq \{ v \in L^2(\Omega) : v |_{\mathcal{K}} \circ F_{\mathcal{K}} \in \mathcal{S}_{q_{\mathcal{K}}}(\widehat{\mathcal{K}}), \mathcal{K} \in \mathcal{T}_h \}.$$

Trefftz FEM Spaces



Polynomial DG Finite Element Spaces: DGFEM uses polynomial basis functions defined on a reference element \widehat{K} :

$$V_q^{DG}(\mathcal{T}_h) \coloneqq \{ v \in L^2(\Omega) : v |_{\mathcal{K}} \circ \mathcal{F}_{\mathcal{K}} \in \mathcal{S}_{q_{\mathcal{K}}}(\widehat{\mathcal{K}}), \mathcal{K} \in \mathcal{T}_h \}.$$

Trefftz Finite Element Space: Use basis functions defined element-wise based on general solutions to the PDE. First define the local Trefftz spaces

$$T(K) \coloneqq \{v|_K : -\Delta u - k^2 u = 0\}$$

and let

$$T(\mathcal{T}_h) \coloneqq \{ v \in L^2(\Omega) : v |_K \in T(K), K \in \mathcal{T}_h \}.$$

Trefftz FEM Spaces



Polynomial DG Finite Element Spaces: DGFEM uses polynomial basis functions defined on a reference element \widehat{K} :

$$V_q^{DG}(\mathcal{T}_h) \coloneqq \{ v \in L^2(\Omega) : v |_{\mathcal{K}} \circ F_{\mathcal{K}} \in \mathcal{S}_{q_{\mathcal{K}}}(\widehat{\mathcal{K}}), \mathcal{K} \in \mathcal{T}_h \}.$$

Trefftz Finite Element Space: Use basis functions defined element-wise based on general solutions to the PDE. First define the local Trefftz spaces

$$T(K) \coloneqq \{v|_K : -\Delta u - k^2 u = 0\}$$

and let

$$T(\mathcal{T}_h) \coloneqq \{ v \in L^2(\Omega) : v |_{\mathcal{K}} \in T(\mathcal{K}), \mathcal{K} \in \mathcal{T}_h \}.$$

We let $V_p(K) \subset T(K)$ be a finite dimensional local space; then, the Trefftz FE Space is given by

$$V_p(\mathcal{T}_h) \coloneqq \{ v \in T(\mathcal{T}_h) : v_K \in V_p(K), K \in \mathcal{T}_h \}.$$

Plane Waves



$$V_{p}(K) = \left\{ v : v(\boldsymbol{x}) = \sum_{\ell=1}^{p_{K}} \alpha_{\ell} e^{ik\boldsymbol{d}_{\ell} \cdot (\boldsymbol{x} - \boldsymbol{x}_{K})}, \alpha_{\ell} \in \mathbb{C} \right\}$$

where p_K is the number of *degrees of freedom* for the element K, d_I , $I = 1, \dots, N_K$ are p_K (roughly) evenly spaced unit direction vectors, and x_K is the centre of the element.

Plane Waves



$$V_{p}(K) = \left\{ v : v(\boldsymbol{x}) = \sum_{\ell=1}^{p_{K}} \alpha_{\ell} e^{ik\boldsymbol{d}_{\ell} \cdot (\boldsymbol{x} - \boldsymbol{x}_{K})}, \alpha_{\ell} \in \mathbb{C} \right\}$$

where p_K is the number of *degrees of freedom* for the element K, d_I , $I = 1, \dots, N_K$ are p_K (roughly) evenly spaced unit direction vectors, and x_K is the centre of the element.

Trefftz DG has less degrees of freedom than high-order polynomials for the same accuracy.

Number of Degrees of Freedom

Plane Waves

$$V_{p}(K) = \left\{ v : v(\boldsymbol{x}) = \sum_{\ell=1}^{p_{K}} \alpha_{\ell} e^{ik\boldsymbol{d}_{\ell} \cdot (\boldsymbol{x} - \boldsymbol{x}_{K})}, \alpha_{\ell} \in \mathbb{C} \right\}$$

where p_K is the number of *degrees of freedom* for the element K, d_I , $I = 1, \dots, N_K$ are p_K (roughly) evenly spaced unit direction vectors, and x_K is the centre of the element.

Trefftz DG has less degrees of freedom than high-order polynomials for the same accuracy.



Number of Degrees of Freedom



Direction Vectors



[Sloan & Womersley, 2004]

TDGFEM for Helmholtz



Trefftz Discontinuous Galerkin FEM for Helmholtz

Find $u_{hp} \in V_p(\mathcal{T}_h)$ such that,

$$\mathcal{A}_h(u_{hp}, v_{hp}) = \ell_h(v_{hp}),$$

for all $v_{hp} \in V_p(\mathcal{T}_h)$, where

$$\begin{split} \mathcal{A}_{h}(u,v) &= \int_{\mathcal{F}_{h}^{l}} \{\!\!\{u\}\!\} [\![\nabla_{h}\bar{v}]\!] \, ds - \int_{\mathcal{F}_{h}^{l}} \beta(ik)^{-1} [\![\nabla_{h}u]\!] [\![\nabla_{h}\bar{v}]\!] \, ds \\ &- \int_{\mathcal{F}_{h}^{l} \cup \mathcal{F}_{h}^{D}} \{\!\!\{\nabla_{h}u\}\!\} \cdot [\![\bar{v}]\!] \, ds + \int_{\mathcal{F}_{h}^{l} \cup \mathcal{F}_{h}^{D}} \alpha ik [\![u]\!] \cdot [\![\bar{v}]\!] \, ds \\ &+ \int_{\mathcal{F}_{h}^{R}} (1-\delta) u \nabla_{h} \bar{v} \cdot \mathbf{n} \, ds - \int_{\mathcal{F}_{h}^{R}} \delta(ik\vartheta)^{-1} (\nabla_{h}u \cdot \mathbf{n}) (\nabla_{h} \bar{v} \cdot \mathbf{n}) \, ds \\ &- \int_{\mathcal{F}_{h}^{R}} \delta \nabla_{h} u \cdot \mathbf{n} \bar{v} \, ds + \int_{\mathcal{F}_{h}^{R}} (1-\delta) ik \vartheta u \bar{v} \, ds, \\ \ell_{h}(v) &= - \int_{\mathcal{F}_{h}^{R}} \delta(ik\vartheta)^{-1} g_{R} \nabla_{h} \bar{v} \cdot \mathbf{n} \, ds + \int_{\mathcal{F}_{h}^{R}} (1-\delta) g_{R} \bar{v} \, ds. \end{split}$$

Scott Congreve (Universität Wien)



Selecting plane wave directions which align with the wave direction of the analytical solution can reduce the error.

Several existing approaches exist for selecting pane wave directions:

- Ray-tracing requires a source term. [Betcke & Phillips, 2012]
- Approximate

$$\frac{\nabla e(\boldsymbol{x}_0)}{ike(\boldsymbol{x}_0)},$$

where *e* is the error. [Gittelson, 2008 (Master's Thesis)]

 Adding an extra unknown (the optimal angle of rotation) to the basis functions. [Amara, Chaudhry, Diaz, Djellouli & Fiedler, 2014]



Selecting plane wave directions which align with the wave direction of the analytical solution can reduce the error.

Several existing approaches exist for selecting pane wave directions:

- Ray-tracing requires a source term. [Betcke & Phillips, 2012]
- Approximate

$$\frac{\nabla e(\boldsymbol{x}_0)}{ike(\boldsymbol{x}_0)},$$

where *e* is the error. [Gittelson, 2008 (Master's Thesis)]

 Adding an extra unknown (the optimal angle of rotation) to the basis functions. [Amara, Chaudhry, Diaz, Djellouli & Fiedler, 2014]

We propose using the Hessian of the numerical solution, based on work on anisotropic meshes for standard FE [Formaggia & Perotto, 2001, 2003].



Plane Wave Refinement Algorithm (2D)

Let $(\lambda_1, \mathbf{v}_1), (\lambda_2, \mathbf{v}_2)$ be the eigenpairs of $\mathbf{H}(\operatorname{Re}(u_h(\mathbf{x}_K)))$, and $(\mu_1, \mathbf{w}_1), (\mu_2, \mathbf{w}_2)$ the eigenpairs of $\mathbf{H}(\operatorname{Im}(u_h(\mathbf{x}_K)))$ s.t. $|\lambda_1| \ge |\lambda_2|$, $|\mu_1| \ge |\mu_2|$; then, for constant C > 1, we can select the first plane wave direction as follows:

 $|\lambda_1| \ge C|\lambda_2| \mid |\mu_1| \ge C|\mu_2| \mid |\lambda_1| \ge C|\mu_1| \mid |\mu_1| \ge C|\lambda_1| \parallel \text{First PW}$

1	1	1	X	v ₁
1	1	×	1	w_1
✓	1	X	X	$\frac{(\boldsymbol{v}_1 + \boldsymbol{w}_1)}{\ \boldsymbol{v}_1 + \boldsymbol{w}_1\ }$
1	×	1	X	v ₁
1	×	×	-	-
×	1	×	1	w_1
×	1	-	X	-
×	×	_	-	-

[C., Houston, Perugia (Submitted)]





9 / 16





Evaluating at $\mathbf{x}_{K} + \delta \mathbf{v}$ we note that the normal is \mathbf{v} , so we can calculate

$$\frac{\nabla u_h(\boldsymbol{x}_K + \delta \boldsymbol{v}) \cdot \boldsymbol{v} + iku_h(\boldsymbol{x}_K + \delta \boldsymbol{v})}{iku_h(\boldsymbol{x}_K + \delta \boldsymbol{v})}.$$





Evaluating at $\mathbf{x}_{\mathcal{K}} + \delta \mathbf{v}$ we note that the normal is \mathbf{v} , so we can calculate

$$\frac{\nabla u_h(\boldsymbol{x}_K + \delta \boldsymbol{v}) \cdot \boldsymbol{v} + iku_h(\boldsymbol{x}_K + \delta \boldsymbol{v})}{iku_h(\boldsymbol{x}_K + \delta \boldsymbol{v})}.$$

We can compare this to the impedance for u:

$$\frac{\nabla u(\boldsymbol{x}_{K} + \delta \boldsymbol{v}) \cdot \boldsymbol{v}}{iku(\boldsymbol{x}_{K} + \delta \boldsymbol{v})} + 1 = \begin{cases} 2, & \text{if } \boldsymbol{d} = \boldsymbol{v}, \end{cases}$$





Evaluating at $\mathbf{x}_{\mathcal{K}} + \delta \mathbf{v}$ we note that the normal is \mathbf{v} , so we can calculate

$$\frac{\nabla u_h(\boldsymbol{x}_K + \delta \boldsymbol{v}) \cdot \boldsymbol{v} + iku_h(\boldsymbol{x}_K + \delta \boldsymbol{v})}{iku_h(\boldsymbol{x}_K + \delta \boldsymbol{v})}.$$

We can compare this to the impedance for u:

$$\frac{\nabla u(\boldsymbol{x}_{K} + \delta \boldsymbol{v}) \cdot \boldsymbol{v}}{iku(\boldsymbol{x}_{K} + \delta \boldsymbol{v})} + 1 = \begin{cases} 2, & \text{if } \boldsymbol{d} = \boldsymbol{v}, \\ 0, & \text{if } \boldsymbol{d} = -\boldsymbol{v}. \end{cases}$$



An *a posteriori* error bounds exists for the *h*-version of the method in \mathbb{R}_2 .

A posteriori Error Bound — h-version Only

For the TDGFEM, with the constant flux parameters, the following error bound holds:

$$\begin{split} \|u - u_{h}\|_{L^{2}(\Omega)}^{2} &\leq C(k, d_{\Omega}) \left\{ \left\| \alpha^{1/2} h_{F}^{s} \llbracket u_{h} \rrbracket \right\|_{L^{2}(\mathcal{F}_{h}^{I} \cup \mathcal{F}_{h}^{D})}^{2} + \frac{1}{k^{2}} \|\beta^{\frac{1}{2}} h_{F}^{s} \llbracket \nabla u_{h} \rrbracket \|_{L^{2}(\mathcal{F}_{h}^{I})}^{2} \\ &+ \frac{1}{k^{2}} \left\| \delta^{1/2} h_{F}^{s} \left(g_{R} - \nabla u_{h} \cdot \boldsymbol{n}_{F} + ik\vartheta u_{h} \right) \right\|_{L^{2}(\mathcal{F}_{h}^{R})}^{2} \right\} \end{split}$$

where s depends on the regularity of the solution to the adjoint problem $(z \in H^{3/2+s}(\Omega))$.

[Kapita, Monk & Warburton, 2015]



A posteriori Error Bound — hp-version

We propose the following potential *a posteriori* error bound with constants derived numerical to ensure the bound is efficient:

$$\begin{aligned} \|u - u_{hp}\|_{L^{2}(\Omega)}^{2} &\leq C \Biggl\{ k \Biggl\| \alpha^{1/2} h_{F}^{1/2} q_{F}^{-1/2} \llbracket u_{hp} \rrbracket \Biggr\|_{L^{2}(\mathcal{F}_{h}^{I} \cup \mathcal{F}_{h}^{D})}^{2} \\ &+ \|\beta^{\frac{1}{2}} h_{F}^{3/2} q_{F}^{-3/2} \llbracket \nabla u_{hp} \rrbracket \|_{L^{2}(\mathcal{F}_{h}^{I})}^{2} \\ &+ \left\| \delta^{1/2} h_{F}^{3/2} q_{F}^{-3/2} \left(g_{R} - \nabla u_{hp} \cdot \boldsymbol{n}_{F} + i k u_{hp} \right) \right\|_{L^{2}(\mathcal{F}_{h}^{R})}^{2} \Biggr\} \end{aligned}$$

for smooth solution of the adjoint.

[C., Houston, Perugia (Submitted)]



Modified hp-refinement Strategy [Melenk & Wohlmuth, 2001]

Let $\mathcal{T}_{h,0}$ be the initial mesh, $\mathcal{T}_{h,i}$ the mesh after *i* refinements, $\eta_{K,i}$ the error indicator for $K \in \mathcal{T}_{h,i}$, and $\eta_{K,i}^{\text{pred}}$ the predicted error for $K \in \mathcal{T}_{h,i}$.

for
$$K \in \mathcal{T}_{h,i}$$
 do
if K is marked for refinement then
if $\eta_{K,i}^2 > (\eta_{K,i}^{\text{pred}})^2$ then
h-refinement: Subdivide K into N sons $K_s, s \in 0, ..., N$
 $(\eta_{K_s,i+1}^{\text{pred}})^2 \leftarrow \frac{1}{N} \gamma_h \left(\frac{1}{2}\right)^{2q_K} \eta_{K,i}^2, i \leq s \leq N$
else
p-refinement: $q_K \leftarrow q_K + 1$
 $(\eta_{K,i+1}^{\text{pred}})^2 \leftarrow \gamma_p \eta_{K,i}^2$

end if

else $(\eta_{K,i+1}^{\mathrm{pred}})^2 \leftarrow \gamma_n (\eta_{K,i}^{\mathrm{pred}})^2$ end if

end for



Consider the smooth (analytic) solution (for Acoustic Wave Propagation)

$$u(x,y) = \mathcal{H}_0^{(1)}(k\sqrt{(x+1/4)^2 + y^2}),$$

on the domain $\Omega = (0, 1)^2$ with suitable Robin BCs. Consider *h*- and *hp*-refinement for k = 20.





Consider the smooth (analytic) solution (for Acoustic Wave Propagation)

$$u(x,y) = \mathcal{H}_0^{(1)}(k\sqrt{(x+1/4)^2 + y^2}),$$

on the domain $\Omega = (0, 1)^2$ with suitable Robin BCs. Consider *h*- and *hp*-refinement for k = 20.





Consider the smooth (analytic) solution (for Acoustic Wave Propagation)

$$u(x,y) = \mathcal{H}_0^{(1)}(k\sqrt{(x+1/4)^2+y^2}),$$

on the domain $\Omega = (0, 1)^2$ with suitable Robin BCs. Consider *h*- and *hp*-refinement for k = 20.



Scott Congreve (Universität Wien)

hp-TDGFEM Adaptive Refinement

13 / 16



Consider the smooth (analytic) solution (for Acoustic Wave Propagation)

$$u(x,y) = \mathcal{H}_0^{(1)}(k\sqrt{(x+1/4)^2+y^2}),$$

on the domain $\Omega = (0, 1)^2$ with suitable Robin BCs. Consider *h*- and *hp*-refinement for k = 50.





Consider the smooth (analytic) solution (for Acoustic Wave Propagation)

$$u(x,y) = \mathcal{H}_0^{(1)}(k\sqrt{(x+1/4)^2 + y^2}),$$

on the domain $\Omega = (0, 1)^2$ with suitable Robin BCs. Consider *h*- and *hp*-refinement for k = 50.



Scott Congreve (Universität Wien)



Consider the smooth (analytic) solution (for Acoustic Wave Propagation)

$$u(x,y) = \mathcal{H}_0^{(1)}(k\sqrt{(x+1/4)^2+y^2}),$$

on the domain $\Omega = (0, 1)^2$ with suitable Robin BCs. Consider *h*- and *hp*-refinement for k = 50.



Scott Congreve (Universität Wien)



Consider the 3D smooth (analytic) solution (for Acoustic Wave Propagation)

$$u(\mathbf{x}) = \mathrm{e}^{ik\mathbf{d}\cdot\mathbf{x}},$$

on the domain $\Omega = (0, 1)^3$, where $d_i = 1/\sqrt{3}$ for i = 1, 2, 3, with suitable Robin BCs.

Consider *h*- and *hp*-refinement for k = 20.





Consider the 3D smooth (analytic) solution (for Acoustic Wave Propagation)

$$u(\mathbf{x}) = \mathrm{e}^{ik\mathbf{d}\cdot\mathbf{x}},$$

on the domain $\Omega = (0, 1)^3$, where $d_i = 1/\sqrt{3}$ for i = 1, 2, 3, with suitable Robin BCs.

Consider *h*- and *hp*-refinement for k = 20.





Consider the 3D smooth (analytic) solution (for Acoustic Wave Propagation)

$$u(\mathbf{x}) = \mathrm{e}^{ik\mathbf{d}\cdot\mathbf{x}},$$

on the domain $\Omega = (0, 1)^3$, where $d_i = 1/\sqrt{3}$ for i = 1, 2, 3, with suitable Robin BCs.

Consider *h*- and *hp*-refinement for k = 50.





Consider the 3D smooth (analytic) solution (for Acoustic Wave Propagation)

$$u(\mathbf{x}) = \mathrm{e}^{ik\mathbf{d}\cdot\mathbf{x}},$$

on the domain $\Omega = (0, 1)^3$, where $d_i = 1/\sqrt{3}$ for i = 1, 2, 3, with suitable Robin BCs.

Consider *h*- and *hp*-refinement for k = 50.





Consider the non-smooth solution (for Acoustic Wave Propagation)

$$u(r, \theta) = \mathcal{J}_{2/3}(kr)\sin(2\theta/3),$$

on the domain L-shaped domain $\Omega = (-1,1)^2 \setminus (0,1) \times (-1,1)$, with suitable Robin BCs.

Consider *h*- and *hp*-refinement for k = 20.



Scott Congreve (Universität Wien)



Consider the non-smooth solution (for Acoustic Wave Propagation)

$$u(r, \theta) = \mathcal{J}_{2/3}(kr)\sin(2\theta/3),$$

on the domain L-shaped domain $\Omega = (-1,1)^2 \setminus (0,1) \times (-1,1)$, with suitable Robin BCs.

Consider *h*- and *hp*-refinement for k = 20.





Consider the non-smooth solution (for Acoustic Wave Propagation)

$$u(r, \theta) = \mathcal{J}_{2/3}(kr)\sin(2\theta/3),$$

on the domain L-shaped domain $\Omega=(-1,1)^2\setminus(0,1)\times(-1,1),$ with suitable Robin BCs.

Consider *h*- and *hp*-refinement for k = 20.



Scott Congreve (Universität Wien)

EFEF 2018 (Heidelberg)

15 / 16



Consider the non-smooth solution (for Acoustic Wave Propagation)

$$u(r, \theta) = \mathcal{J}_{2/3}(kr)\sin(2\theta/3),$$

on the domain L-shaped domain $\Omega=(-1,1)^2\setminus(0,1)\times(-1,1),$ with suitable Robin BCs.

Consider *h*- and *hp*-refinement for k = 50.



Scott Congreve (Universität Wien)



Consider the non-smooth solution (for Acoustic Wave Propagation)

$$u(r, \theta) = \mathcal{J}_{2/3}(kr)\sin(2\theta/3),$$

on the domain L-shaped domain $\Omega = (-1,1)^2 \setminus (0,1) \times (-1,1)$, with suitable Robin BCs.

Consider *h*- and *hp*-refinement for k = 50.



Scott Congreve (Universität Wien)



Consider the non-smooth solution (for Acoustic Wave Propagation)

$$u(r, \theta) = \mathcal{J}_{2/3}(kr)\sin(2\theta/3),$$

on the domain L-shaped domain $\Omega=(-1,1)^2\setminus(0,1)\times(-1,1),$ with suitable Robin BCs.

Consider *h*- and *hp*-refinement for k = 50.



Scott Congreve (Universität Wien)

EFEF 2018 (Heidelberg)

15 / 16



Summary:

- With plane wave basis functions it is possible to refine the wave directions.
- *hp*-adaptive refinement results in exponential convergence.
- Combining plane wave direction adaptivity with *hp*-adaptive refinement often leads to reduced error compared to standard refinement.

Future Aims:

- Develop robust *hp*-version *a posteriori* error bounds..
- Use the eigenvalues/eigenvectors to develop anisotropic *p*-refinement (unevenly spaced plane waves).