Mechanical Analogs For One Dimensional Viscoelastic Models.

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Figure 3: one of five arrangements of springs and dashpots that leads to a Burgers' fluid model. Includes the Maxwell and Oldroyd-B model as special cases. (see Karra and Rajagopal (2009)).

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One dimensional Maxwell fluid model derivation  
based on the mechanical analog.  

$$f$$
  $K$   $F$   
Let  $\Delta$  denote the elongation  
Let  $\cdot$  denote the elongation  
Let  $\cdot$  denote time derivative.  
Then total elongation  $\Delta = \Delta_s + \Delta_d$  (1)  
elongation in the spring the dealoght.  
Equilibrium  $\Rightarrow$   $F_s = F_d = F$  (2).  
 $F_s -$  force in the spring  
 $F_d -$  force in the deslopht.

Constitutive Equations:  
For the spring: 
$$F_5 = K\Delta_5$$
. (3)  
For the dashpot:  $F_4 = \overline{\chi} A d$ . (4)  
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(i)  $-(L_1) \Rightarrow \Delta_5 + \Delta_4 = \frac{F_5}{K} + \frac{F_4}{\overline{\chi}} = -\frac{F}{K} + \frac{F_7}{\overline{\chi}} = \Delta$   
In terms of stress-strain this leads to  
 $\overline{(\tau + \lambda \hat{\sigma} = \gamma \hat{\epsilon})}$   
Units of  $\overline{\chi}$  and  $\gamma$  are different.  
 $dim(\lambda) = T$  (time)  
 $dim(\gamma) = ME^{1}T^{-1}$ 

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## • Helmholtz Potential Approach: Will not always work.

# • Pick rate of entropy production

- due to conduction
- Output to work being converted into energy in thermal form
- I due to mixing
- due to phase transition
- I due to growth

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- Part of the energy supplied can
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  - change the "strain" or "stored" energy
  - 🕘 etc.
- Need to know how the energy that is stored can be recovered
  - In a purely mechanical process
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  - letc.

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#### $\psi$ – Helmholtz potential

• If the material is incompressible, then we need:

$$\operatorname{div}(\mathbf{v}) = 0. \tag{21}$$

- Maximize ξ subject to (20) and (21) as constraints.
- Auxiliary function to be maximized

$$\Phi = \xi + \lambda_1 \left( \xi - \mathbf{T} \cdot \mathbf{D} - \rho \dot{\psi} \right) + \lambda_2 (\operatorname{div}(\mathbf{v})).$$
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$$\mathbf{\bar{A}} := \frac{d\mathbf{A}}{dt} - \mathbf{L}\mathbf{A} - \mathbf{A}\mathbf{L}^{\mathsf{T}}$$
  
Upper convected Oldroyd derivative  
•  $\mathbf{\bar{B}}_{\kappa_{p(t)}}^{\nabla} = -2\mathbf{F}_{\kappa_{p(t)}}\mathbf{D}_{\kappa_{p(t)}}\mathbf{F}_{\kappa_{p(t)}}^{\mathsf{T}}$   
•  $\mathbf{L}_{\kappa_{p(t)}} := \dot{\mathbf{G}}\mathbf{G}^{-1}$   
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- Get a generalization of the Oldroyd-B model.
- Can develop generalization of Burgers' fluid, etc.
- Can develop generalization of any system of springs and dashpots.
- Can obtain models which use ideas such as "confirmation tensors": FENE-P model.
- Can develop models for anisotropic fluids.
- However, cannot be used to develop certain classes of viscoelastic fluid.
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- Pick  $\xi = \eta \mathbf{D}_{\kappa_{p(t)}} \cdot \mathbf{B}_{\kappa_{p(t)}} \mathbf{D}_{\kappa_{p(t)}} + \eta_1 \mathbf{D} \cdot \mathbf{D}_{\kappa_{p(t)}}$
- Get a generalization of the Oldroyd-B model.
- Can develop generalization of Burgers' fluid, etc.
- Can develop generalization of any system of springs and dashpots.
- Can obtain models which use ideas such as "confirmation tensors": FENE-P model.
- Can develop models for anisotropic fluids.
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 Ideas developed by D'Alembert and Bernoulli and used by Lagrange that appeal to constraint forces doing no work are WRONG. This is at the heart of how one shows that

$$\mathbf{T} = -p\mathbf{I} + \mathbf{S}$$

in an incompressible fluid.

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$$G = G(\mathbf{S}), \quad \mathbf{S} = \frac{\mathbf{T}}{\rho}$$
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- $\xi = \xi (\mathbf{S}, \rho, \mathbf{D})$
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- A. N. Whitehead

- Most people would rather die than think. Most do. - B. Russell
- He who criticizes must first read.

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